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# Spectral analysis of the zigzag process 

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#### Abstract

The zigzag process is a variant of the telegraph process with position dependent switching intensities. A characterization of the $L^{2}$-spectrum for the generator of the one-dimensional zigzag process is obtained in the case where the marginal stationary distribution on $\mathbb{R}$ is unimodal and the refreshment intensity is zero. Sufficient conditions are obtained for a spectral mapping theorem, mapping the spectrum of the generator to the spectrum of the corresponding Markov semigroup. Furthermore results are obtained for symmetric stationary distributions and for perturbations of the spectrum, in particular for the case of a non-zero refreshment intensity. In the examples we consider (including a Gaussian target distribution) a slight increase of the refreshment intensity above zero results in a larger $L^{2}$-spectral gap, corresponding to an improved convergence in $L^{2}$.


Résumé. Le processus de ZigZag est une variante du processus du télégraphe avec des intensités de retournement dépendant de la position. Nous obtenons une caractérisation du spectre en norme $L^{2}$ du générateur du processus de zigzag unidimensionnel, dans le cas où la distribution marginale stationnaire sur $\mathbb{R}$ est unimodale, avec une fréquence de rééchantillonnage nulle. Nous obtenons des conditions suffisantes pour un théorème d'isomorphisme spectral, identifiant le spectre du générateur à celui du semi-groupe de Markov correspondant. Par ailleurs, nous obtenons des résultats pour les distributions stationnaires symétriques ainsi que pour les perturbations du spectre dans le cas particulier d'une fréquence de rééchantillonnage non nulle. Enfin, nous considérons dans les exemples (avec une distribution cible gaussienne) un faible taux de rafraichissement positif par rapport aux résultats du taux nul, ce qui induit une plus grande bande dans le spectre $L^{2}$, correspondant à une convergence améliorée en norme $L^{2}$.

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## 1. Introduction

In recent years piecewise deterministic Markov processes have emerged as a useful computational tool in stochastic simulation, for example for Bayesian statistics or statistical physics. A particular instance of such a process is the zigzag process, described already in e.g. [12,21] and given its name in [5]. The zigzag process is a variant of, and extends the telegraph process [16] and is intimately related to the Bouncy Particle Sampler [6,22].

The zigzag process can be defined in multiple dimensions [4] but here we consider only one spatial dimension. In this setting the zigzag process is a Markov process $(X(t), \Theta(t))$ in the state space $E:=\mathbb{R} \times\{-1,+1\}$. Conditional on the velocity process $(\Theta(t))_{t \geq 0}$, the position process $(X(t))_{t \geq 0}$ in $\mathbb{R}$ is completely determined by the relation $X(t)=X(0)+$ $\int_{0}^{t} \Theta(s) d s$. The velocity process $(\Theta(t))_{t \geq 0}$ in $\{-1,+1\}$ switches sign at inhomogeneous Poisson rate $\lambda(X(t), \Theta(t))$, where $\lambda: E \rightarrow[0, \infty)$ is a continuous function. Thus the switching intensity function $\lambda$ has a crucial impact on the dynamics exhibited by the zigzag process.

For a given absolutely continuous potential function $U: \mathbb{R} \rightarrow \mathbb{R}$, the switching intensity $\lambda$ can be chosen as $\lambda(x, \theta)=\left(\theta U^{\prime}(x) \vee 0\right)+\lambda_{\text {refr }}(x)$ for some non-negative refreshment intensity function $\lambda_{\text {refr }}: \mathbb{R} \rightarrow[0, \infty)$. This condition results in the zigzag process having stationary distribution $\mu$ on $E$ with marginal density on $\mathbb{R}$ proportional to $e^{-U}$. This observation makes the zigzag process into a useful computational device for the stochastic simulation of a given target density proportional to $e^{-U}$.

In the design of stochastic simulation algorithms it is vital to understand the convergence to stationarity, i.e., the quality of the approximation of the probability density proportional to $e^{-U}$ by the law of $X(t)$. For the zigzag process this question has already been addressed in several ways. In [5,13] a Lyapunov function argument has been used to show exponentially fast convergence to stationarity in total variation norm. In [3] a central limit theorem for $\frac{1}{\sqrt{T}} \int_{0}^{T}(f(X(s), \Theta(s))-\pi(f)) d s$, for $T \rightarrow \infty$, was obtained under suitable conditions. Here $\pi$ denotes the normalization of $\mu$ into a probability distribution. The hypocoercivity approach of [9] is applied in [1] to the zigzag process, amongst other piecewise deterministic processes, in order to obtain exponential convergence in $L^{2}(\mu)$. Recently [19] provided estimates on the exponential convergence of PDMPs using a time-dependent Poincaré inequality, that seem to result in a good understanding of dependence on, e.g., dimensionality, anisotropy and convexity. However, typically these results either do not obtain a tractable quantitative bound on the rate of convergence, and/or provide only a one-sided estimate on the rate of convergence.

In order to obtain quantitative result on the rate of convergence to stationarity, it is natural to investigate the spectral properties of the Markov semigroup $(P(t))_{t \geq 0}$ connected to the zigzag process, and this is the aim of this work. The probabilistic interpretation of $P(t)$ is that

$$
P(t) f(x, \theta)=\mathbb{E}_{x, \theta} f(X(t), \Theta(t)), \quad t \geq 0
$$

where $\mathbb{E}_{x, \theta}$ denotes the expectation over trajectories of the Markov process $(X(t), \Theta(t))$ with initial position $(X(0), \Theta(0))=(x, \theta)$. In this paper we take a fully analytical approach, and thus construct and analyse the semigroup $(P(t))_{t \geq 0}$ using functional analytical arguments with no reliance on probabilistic results.

In a similar vein in [20] a rather complete understanding of the spectral properties in $L^{2}(\mu)$ is obtained for a periodic variant of the telegraph process, with uniform marginal stationary distribution on the position space. As in $[1,20]$ we consider the spectrum of the zigzag semigroup in the Hilbert space $L^{2}(\mu)$. Although the zigzag semigroup is not selfadjoint, it still has remarkable symmetry properties as a semigroup in $L^{2}(\mu)$. In particular we obtain an elegant characterization of the adjoint of the generator (see Lemma 3.4), and a decomposition of the spectrum in the case of a symmetric stationary distribution (Section 5).

Our results rely for a significant part on the explicit solution of the eigenvalue and resolvent equations, which is possible if $\lambda_{\text {refr }}=0$. We anticipate that the results can be extended to $\lambda_{\text {refr }} \neq 0$ and to a wider variety of function spaces, using an approach based upon characteristic matrices [15], which will be the focus of further research.

The structure of this work is as follows. First, in Section 2 we introduce notation and our main assumptions which are used throughout this document. In Section 3 we construct, imposing minimal assumptions and taking a purely analytical approach, the zigzag semigroup $(P(t))_{t \geq 0}$ on $L^{2}(\mu)$, where we characterize the domain of the infinitesemal generator $L$ explicitly (Theorem 3.6). Under mild conditions it is established that the resolvent of $(P(t))_{t \geq 0}$ is compact with immediate strong implications for the spectrum. Then, in Section 4 the spectrum of the zigzag semigroup is studied under the assumption of unimodality of $e^{-U}$ and vanishing refreshment, $\lambda_{\text {refr }}(x)=0$ for all $x$. It turns out that in this case we can find a holomorphic function $Z: \mathbb{C} \rightarrow \mathbb{C}$ such that the spectrum of $L$ is identical to the roots of $Z$ (Theorem 4.4). Under a further suitable polynomial growth condition on the potential function, a spectral mapping theorem is shown to hold, so that the spectrum of $(P(t))_{t \geq 0}$ can expressed in terms of the roots of $Z$ (Theorem 4.13), yielding in particular a characterization of the spectral gap of $(P(t))_{t \geq 0}$ in $L^{2}(\mu)$ (Theorem 4.14).

Interestingly, if $U$ is symmetric, as discussed in Section 5, then the state space $L^{2}(\mu)$ can be separated into two subspaces that are invariant with respect to the zigzag semigroup. In particular this implies a decomposition of the spectrum (Theorem 5.3). In Section 6 we determine the effect of small perturbations of the generator on the spectrum of the zigzag semigroup, allowing us to study the effect of a small non-zero refreshment intensity.

In Section 7 some fundamental examples are discussed along with a numerical illustration of the spectrum. Interestingly, in the examples we consider (including a Gaussian target distribution) a slight increase above zero in the zigzag refreshment intensity results in a larger spectral gap and thus an improved convergence in $L^{2}$. In contrast, the asymptotic variance increases when the refreshment rate increases [3]. This apparent contradiction between convergence in law and convergence of empirical averages is also observed in reversible, nearly-periodic Markov chains [24] and non-reversible Markov chains [25].

Finally we include an appendix containing several technical lemmas.

## 2. Notation and assumptions

Let $E=\mathbb{R} \times\{-1,+1\}$ and equip $E$ with the product topology generated by the usual topologies on $\mathbb{R}$ and $\{-1,+1\}$. We also equip $E$ with the product $\sigma$-algebra generated by products of Lebesgue sets in $\mathbb{R}$ and all subsets of $\{-1,+1\}$.

Associated to a potential function $U: \mathbb{R} \rightarrow \mathbb{R}$, which we assume to be continuously differentiable, we define the switching intensity function

$$
\begin{equation*}
\lambda(x, \theta):=\left(\theta U^{\prime}(x) \vee 0\right)+\lambda_{\text {refr }}(x), \quad(x, \theta) \in E . \tag{2.1}
\end{equation*}
$$

Here $(a \vee b):=\max (a, b)$ for any $a, b \in \mathbb{R}$, and $\lambda_{\text {refr }}: \mathbb{R} \rightarrow[0, \infty)$ is a continuous function. The function $\lambda_{\text {refr }}$ is referred to as the refreshment rate or excess switching rate. In order to obtain the main results in this paper we will often assume that $\lambda_{\text {refr }}(x)=0$ for all $x \in \mathbb{R}$; this assumption will always be made explicit.

The switching intensities $\lambda$ define (under suitable conditions) a Markov process in $E$, with associated Markov semigroup $(P(t))_{t \geq 0}$ on $B(E)$, the Banach space of bounded measurable functions on $E$. We call $(P(t))$ the zigzag semigroup. For a probabilistic construction see, e.g., $[4,5]$.

It is the aim of this paper to carry out a detailed investigation of properties of this semigroup as a strongly continuous semigroup in $L^{2}(\mu)$, where $\mu$ denotes a stationary distribution of the process (to be made precise later). We will be particularly interested in spectral properties of the semigroup.

### 2.1. Notation

For a measure $\mu$ on a measurable space $E$ we write $\mu(f)=\int_{E} f d \mu$, for $f \in L^{1}(\mu)$. Let $E=\mathbb{R} \times\{-1,+1\}$. For $f: E \rightarrow$ $\mathbb{C}$ we write $f^{\theta}=f(\cdot, \theta), \theta= \pm 1$. Recall the Sobolev spaces $W^{k, p}(\Omega, v)$ for $\Omega$ an open subset of $\mathbb{R}$ and $v$ a Borel measure on $\Omega$ and write $W^{k, p}(\Omega)$ when the underlying measure is Lebesgue measure. Let $\mu$ denote a Borel measure on $E$. We say that $f \in W^{k, p}(\mu)$ if $f(\cdot, \theta) \in W^{k, p}(\mathbb{R}, \mu)$ for $\theta \in\{-1,+1\}$. Similarly we say that $f \in W_{\text {loc }}^{k, p}(E)$ if, for any $\theta \in\{-1,+1\}$, $f(\cdot, \theta)$ restricted to $F$ is in $W^{k, p}(F)$ for any compact set $F \subset \mathbb{R}$. Mappings $M$ of functions $f: \mathbb{R} \rightarrow \mathbb{C}$ yielding a function $M f: \mathbb{R} \rightarrow \mathbb{C}$ may be extended to mappings of functions $f: E \rightarrow \mathbb{C}$ by defining $(M f)^{\theta}=M f(\cdot, \theta)=M f^{\theta}$ for $\theta= \pm 1$. For example, for $f \in W^{1, p}(E)$ we write $\partial_{x} f: E \rightarrow \mathbb{R}$ for the function such that $\left(\partial_{x} f\right)^{\theta}=\partial_{x}\left(f^{\theta}\right)$ for $\theta= \pm 1$. Similarly a function space of functions $f: \mathbb{R} \rightarrow \mathbb{R}$ is readily extended to functions $f: E \rightarrow \mathbb{R}$. For example, $C_{c}^{\infty}(E)$ denotes the space of all functions $f: E \rightarrow \mathbb{C}$ such that $f^{\theta} \in C_{c}^{\infty}(\mathbb{R})$ for $\theta= \pm 1$.

### 2.2. Assumptions on the potential function

We introduce several different assumptions on the potential function $U$ for later reference. See Section 7 for the verification of these conditions for a specific family of potential functions.

Assumption A1. $U \in C^{2}(\mathbb{R}), e^{-U} \in L^{1}(\mathbb{R})$ and, for some $\delta \in(0,1), m \in \mathbb{R}$, and $M>0$,

$$
m \leq U^{\prime \prime}(x) \leq \delta\left|U^{\prime}(x)\right|^{2}+M, \quad x \in \mathbb{R}
$$

Furthermore there are constants $c_{1}, c_{2} \geq 0$ such that $\lambda_{\text {refr }}(x) \leq c_{1}+c_{2}\left|U^{\prime}(x)\right|$ for all $x \in \mathbb{R}$.

Assumption A2. $\lim _{x \rightarrow \infty} U^{\prime}(x)=+\infty$ and $\lim _{x \rightarrow-\infty} U^{\prime}(x)=-\infty$.
Assumption A3 (Unimodality). $U(0)=0, U^{\prime}(x) \leq 0$ for $x \leq 0$ and $U^{\prime}(x) \geq 0$ for $x \geq 0$.

Assumption 44 (Growth condition). There are constants $C \leq 0, M \geq 0, m>0$ and $p>1$ such that

$$
\begin{equation*}
U(y) \geq U(x)+C+m|x-y|^{p} \quad \text { for all } x, y \text { for which } M \leq x \leq y \text { or } y \leq x \leq-M \tag{2.2}
\end{equation*}
$$

Note that for continuously differentiable $U$, Assumption A4 cannot hold for $C>0$. Indeed, if (2.2) would hold with $C>0$ then for all $x \geq M$

$$
U^{\prime}(x)=\lim _{h \downarrow 0} \frac{U(x+h)-U(x)}{h} \geq \lim _{h \downarrow 0} \frac{C}{h}+m h^{p-1}=+\infty
$$

Remark 2.1. If Assumption (2.2) holds for some $M>0$, then it also holds for $M=0$ (for a different value of $m$ ), which simplies proofs later on. This result is delegated to the appendix, see Lemma A.2.

Examples of potential functions for which Assumption A4 is satisfied are given by $U(x)=|x|^{p}$ for $p>1$, with $C=0$ and $m=1$. This can be seen by defining, for fixed $x \geq 0, f(y):=|y|^{p}-|x|^{p}-|y-x|^{p}$ and noting that $f^{\prime}(y) \geq 0$ for $y \geq x$.

Remark 2.2. For convenience of the reader we provide the following lemma which gives an easily verifiable condition for Assumption A4 to hold. The reverse implication does not hold: see Lemma 7.1 for an example where the conditions of Lemma 2.3 are not satisfied whereas Assumptions A2 and A4 are satisfied.

Lemma 2.3. Suppose $U \in C^{2}(\mathbb{R})$ and $U^{\prime \prime}(x) \geq m$ for all $|x| \geq M$, where $m>0$ and $M \geq 0$. Then Assumptions $A 2$ and A4 hold.

Proof. The implication of Assumption A2 is immediate. If the assumptions of the lemma are satisfied then $U^{\prime}(x) \geq 0$ for $x$ sufficiently large (say for $x \geq \widetilde{M} \geq M$ ). By integrating it follows that for $y \geq x \geq \widetilde{M}$ and $\widetilde{m}=m / 2$ we have that

$$
U(y) \geq U(x)+U^{\prime}(x)(y-x)+\widetilde{m}(y-x)^{2} \geq U(x)+\widetilde{m}(y-x)^{2} .
$$

We may use an analogous argument for $y \leq x \leq-M$ to obtain Assumption A4 for $C=0$ and $p=2$.
In Section 5 we will investigate the implications of the following symmetry assumption.
Assumption A5 (Symmetry). $U(x)=U(-x)$ for $x \in \mathbb{R}$.
Remark 2.4. For convenience of the reader we remark that the following list of assumptions implies all conditions mentioned so far, but is somewhat stronger than necessary. In particular (iii) and (v) are stronger than necessary for most of the results.
(i) $U \in C^{2}(\mathbb{R})$.
(ii) For some $\delta \in(0,1)$, and $M>0$,

$$
U^{\prime \prime}(x) \leq \delta\left|U^{\prime}(x)\right|^{2}+M, \quad x \in \mathbb{R} .
$$

(iii) Symmetry: $U(x)=U(-x)$ for $x \in \mathbb{R}$.
(iv) Unimodality: $U(0)=0$ and $U^{\prime}(x) \geq 0$ for $x \geq 0$.
(v) $m$-strong convexity: $U^{\prime \prime}(x) \geq m$ for all $|x| \geq M$, where $m>0$ and $M \geq 0$.

In particular note that $e^{-U} \in L^{1}(\mathbb{R})$ is implied by (v).
Example 2.5 (Gaussian distribution). Suppose $U(x)=\frac{x^{2}}{2 \sigma^{2}}$ with $\sigma>0$. Then Assumptions A1, A2, A3, A4 and A5 are satisfied. Here Assumption A4 follows, e.g., by Lemma 2.3. We will study a more general family of distributions, including the Gaussian distribution, in Section 7.

Example 2.6 ( $t$-distribution). Suppose $U(x)=\left(\frac{\alpha+1}{2}\right) \log \left(1+\frac{x^{2}}{\alpha}\right)$, with $\alpha>0$. In this case the probability measure on $\mathbb{R}$ with density proportional to $e^{-U}$ is the $t$-distribution with $\alpha$ degrees of freedom. In particular if $\alpha=1$ this is the Cauchy distribution. Then $U^{\prime}(x)=\frac{(\alpha+1) x}{\alpha+x^{2}}$ and $U^{\prime \prime}(x)=\frac{(\alpha+1)\left(\alpha-x^{2}\right)}{\left(x^{2}+\alpha\right)^{2}}$. Since $U^{\prime \prime}$ is a bounded function for any $\alpha>0$, Assumption A1 is satisfied. However, $\left|U^{\prime}(x)\right| \rightarrow 0$ as $|x| \rightarrow \infty$ for any $\alpha>0$, so that Assumption A2 is not satisfied. Assumptions A3 and A5 are trivially satisfied; Assumption A4 is not satisfied.

## 3. The zigzag semigroup in $L^{2}(\mu)$

In this section we take a purely functional analytic approach to the construction of the zigzag semigroup in $L^{2}(\mu)$. That is, we do not refer to the probabilistic construction of the process as carried out, e.g., in [5]. This construction is discussed in Section 3.1. Next, we establish that under reasonable conditions (i.e. Assumptions A1, A2) the semigroup has a compact resolvent, which is important in the characterization of the spectrum from Section 4 onwards.

### 3.1. Construction of the zigzag semigroup

Define measures $v$ on $\mathbb{R}$ and $\mu$ on $E$, respectively by

$$
\nu(A)=\int_{A} e^{-U(x)} d x, \quad \text { and } \quad \mu(A \times\{\theta\})=v(A), \quad A \in \mathcal{B}(\mathbb{R}), \theta \in\{-1,+1\}
$$

At this point there is no requirement that $e^{-U}$ be integrable. Therefore the measure $\mu$ is not necessarily finite, but as a consequence of the continuity of $U$ it is $\sigma$-finite. We consider the complex Hilbert space $L^{2}(\mu)$ with inner product $\langle\cdot, \cdot \cdot\rangle$ and norm $\|\cdot\|$. The space $L^{2}(\mu)$ is separable by [7, p. 36]. Define a linear operator $(L, \mathcal{D}(L))$ on the domain

$$
\begin{equation*}
\mathcal{D}(L)=\left\{f \in L^{2}(\mu) \cap W_{\mathrm{loc}}^{1,2}(E):(x, \theta) \mapsto \theta \partial_{x} f(x, \theta)+\lambda(x, \theta)(f(x,-\theta)-f(x, \theta)) \in L^{2}(\mu)\right\} \tag{3.1}
\end{equation*}
$$

by

$$
\begin{equation*}
L f(x, \theta)=\theta \partial_{x} f(x, \theta)+\lambda(x, \theta)(f(x,-\theta)-f(x, \theta)) \tag{3.2}
\end{equation*}
$$

We introduce the flipping operator $F f(x, \theta)=f(x,-\theta)$ and we will often use the shorthand notation $L f=\theta \partial_{x} f+$ $\lambda(F f-f)$.

The aim of this section is to establish that $L$ is the generator of a semigroup $(P(t))_{t \geq 0}$ in $L^{2}(\mu)$, which we will refer to as the zigzag semigroup. In order to do so, we will first show that $L$ is a closed dissipative operator (Lemmas 3.1, 3.3), determine a core for $L$ (Lemma 3.2), and next characterize the adjoint of $L$, finding essentially that $L^{\star}=F L F$ (Lemmas 3.4, 3.5). These results then immediately imply the wellposedness of the Cauchy problem, i.e., the construction of the semigroup $P(t)$, see Theorem 3.6.

Lemma 3.1. The operator $(L, \mathcal{D}(L))$ is closed.
Proof. Let $\left(f_{n}\right) \subset \mathcal{D}(L)$ be a converging sequence in $L^{2}(\mu)$ with $\lim _{n \rightarrow \infty} f_{n}=f \in L^{2}(\mu)$, and suppose $\left(g_{n}\right):=\left(L f_{n}\right)$ converges in $L^{2}(\mu)$ to $g$. Note that

$$
\begin{equation*}
\partial_{x} f_{n}(x, \theta)=\theta\left(g_{n}(x, \theta)-\lambda(x, \theta)\left(f_{n}(x,-\theta)-f_{n}(x, \theta)\right)\right) \tag{3.3}
\end{equation*}
$$

and define in similar spirit

$$
h(x, \theta):=\theta(g(x, \theta)-\lambda(x, \theta)(f(x,-\theta)-f(x, \theta)))
$$

We will first show that $f \in W_{\mathrm{loc}}^{1,2}(E)$ and that $\partial_{x} f=h$. To this end let $\phi \in C_{c}^{\infty}(E)$. Then

$$
\begin{aligned}
-\int_{E} f \partial_{x} \phi d x & =-\int_{E}\left(\lim _{n \rightarrow \infty} f_{n}\right) \partial_{x} \phi d x=-\lim _{n \rightarrow \infty} \int_{E} f_{n} \partial_{x} \phi d x \\
& =\lim _{n \rightarrow \infty} \int_{E}\left(\partial_{x} f_{n}\right) \phi d x=\int_{E}\left(\lim _{n \rightarrow \infty} \partial_{x} f_{n}\right) \phi d x=\int_{E} h \phi d x
\end{aligned}
$$

In he above derivation we were allowed to exchange the limit and integral twice, since (for the first interchange), $f_{n} \rightarrow$ $f$ in $L^{2}(\mu)$, and using (for the second interchange) that $f_{n} \rightarrow f$ in $L^{2}(\mu), g_{n} \rightarrow g$ in $L^{2}(\mu)$, and $\lambda$ is bounded on compact intervals, and thus $\partial_{x} f_{n}$ converges in $L_{\text {loc }}^{2}(\mathbb{R})$ by (3.3). This establishes that $\partial_{x} f=h$. Furthermore, taking a subsequence, $\theta \partial_{x} f_{n_{k}}=g_{n_{k}}-\lambda\left(F f_{n_{k}}-f_{n_{k}}\right) \rightarrow g-\lambda(F f-f)=\theta \partial_{x} f$ almost everywhere (by $L^{2}$-convergence of $f_{n}$ and $g_{n}$ ). Therefore $g_{n_{k}}=L f_{n_{k}} \rightarrow \theta \partial_{x} f+\lambda(F f-f)=g$ almost everywhere. Since also $g_{n_{k}} \rightarrow g$ in $L^{2}(\mu)$ we have $f \in \mathcal{D}(L)$ and $L f=g$.

Lemma 3.2. $C_{c}^{\infty}(E)$ is a core for $L$.
Proof. We use ideas from the proof of [18, Proposition 8.7.3]. Let $f \in \mathcal{D}(L)$, and let $\xi:[0, \infty) \rightarrow[0,1]$ be a smooth function such that $\xi=1$ in $[0,1 / 2]$ and $\xi=0$ in $[1, \infty)$, and define $\xi_{n}(x)=\xi(|x| / n)$. Then define $f_{n}(x, \theta)=$ $f(x, \theta) \xi_{n}(x)$. Then every $f_{n}$ has compact support, $f_{n} \in W^{1,2}(\mu)$, and $f_{n} \rightarrow f$ in $L^{2}(\mu)$ by dominated convergence. Furthermore we have

$$
L f_{n}(x, \theta)=\theta \xi_{n}^{\prime}(x) f(x, \theta)+\xi_{n}(x) L f(x, \theta)
$$

Note that $\sup _{x \in \mathbb{R}}\left|\xi_{n}^{\prime}(x)\right| \leq C / n$ for some $C>0$. It follows that, as $n \rightarrow \infty$,

$$
\left\|L f_{n}(x, \theta)-L f(x, \theta)\right\|_{L^{2}(\mu)} \leq C / n\|f\|_{L^{2}(\mu)}+\left\|\left(1-\xi_{n}\right) L f\right\|_{L^{2}(\mu)} \rightarrow 0
$$

where the second term converges to zero by dominated converence. Thus, functions with compact support are dense in $\mathcal{D}(L)$ with respect to the graph norm of $L$. We can now limit ourselves to the situation in which $f$ has compact support.

Let $\varphi \in C_{c}^{\infty}(\mathbb{R})$ be a smooth function, compactly supported in $[-1,1]$ such that $\varphi=1$ in $[-1 / 2,1 / 2], \varphi(x) \in[0,1]$ for all $x$, and $\|\varphi\|_{L^{1}(\mathbb{R})}=1$. Set $\varphi_{n}(x)=n \varphi(x / n)$ for $n \in \mathbb{N}$, and for compactly supported $f$, define

$$
f_{n}(x, \theta)=\left(f \star \varphi_{n}\right)(x, \theta):=\int_{\mathbb{R}} \varphi_{n}(x-y) f(y, \theta) d y, \quad x \in \mathbb{R}, n \in \mathbb{N} .
$$

Then $f_{n} \in C_{c}^{\infty}(E)$, uniformly supported in $\operatorname{supp}(f)+B(1)$, and $f_{n} \rightarrow f$ in $L^{2}(\mu)$. Furthermore, since $\partial_{x} f_{n}=\left(\partial_{x} f\right) \star \varphi_{n}$, and $\lambda$ is bounded uniformly in $n$ on the support of ( $f_{n}$ ), it follows that $\left\|L f_{n}-L f\right\|_{L^{2}(\mu)} \rightarrow 0$.

Lemma 3.3 (Dissipativity). $\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}-f^{-}\right|^{2} \in L^{1}(\nu)$ for any $f \in \mathcal{D}(L)$, and

$$
\begin{equation*}
\operatorname{Re} \int_{E}(L f) \bar{f} d \mu \leq-\frac{1}{2} \int_{\mathbb{R}}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}-f^{-}\right|^{2} d \nu, \quad f \in \mathcal{D}(L) . \tag{3.4}
\end{equation*}
$$

Proof. First we establish the result for $f \in C_{c}^{\infty}(E)$. By partial integration, using compact support of $f$, and noting that (2.1) leads to the equality $\lambda^{+}(x)-\lambda^{-}(x)=U^{\prime}(x)$, we find that

$$
\begin{aligned}
\operatorname{Re} & \sum_{\theta= \pm 1} \int_{-\infty}^{\infty}(L f)^{\theta} \overline{f^{\theta}} d v(x) \\
= & \operatorname{Re} \int_{-\infty}^{\infty}\left\{\left(\partial_{x} f^{+}+\lambda^{+}\left(f^{-}-f^{+}\right)\right) \overline{f^{+}}+\left(-\partial_{x} f^{-}+\lambda^{-}\left(f^{+}-f^{-}\right)\right) \overline{f^{-}}\right\} e^{-U} d x \\
= & \int_{-\infty}^{\infty} \frac{1}{2} \partial_{x}\left(\left|f^{+}\right|^{2} e^{-U}\right)+\operatorname{Re}\left(\frac{1}{2} U^{\prime} f^{+}+\lambda^{+}\left(f^{-}-f^{+}\right)\right) \overline{f^{+}} e^{-U} d x \\
& +\int_{-\infty}^{\infty}\left(-\frac{1}{2} \partial_{x}\left(\left|f^{-}\right|^{2} e^{-U}\right)+\operatorname{Re}\left(-\frac{1}{2} U^{\prime} f^{-}+\lambda^{-}\left(f^{+}-f^{-}\right)\right) \overline{f^{-}} e^{-U} d x\right. \\
= & \operatorname{Re} \int_{-\infty}^{\infty}\left(\lambda^{+} f^{-} \overline{f^{+}}+\lambda^{-} f^{+} \overline{f^{-}}-\frac{1}{2}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}\right|^{2}-\frac{1}{2}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{-}\right|^{2}\right) e^{-U} d x \\
= & -\frac{1}{2} \int_{-\infty}^{\infty}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}-f^{-}\right|^{2} e^{-U} d x .
\end{aligned}
$$

For $f \in \mathcal{D}(L)$, using that $C_{c}^{\infty}(E)$ is a core for $L$ (Lemma 3.2), take a sequence $f_{n} \rightarrow f$, with $f_{n} \in C_{c}^{\infty}(E)$, and with convergence in the graph norm of $L$. By Fatou,

$$
\frac{1}{2} \int_{\mathbb{R}}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}-f^{-}\right|^{2} d \nu \leq-\operatorname{Re} \int_{E}(L f) \bar{f} d \mu<\infty
$$

Recall that the adjoint operator of a densely defined operator $(A, \mathcal{D}(A))$ on $L^{2}(\mu)$ is defined on the domain

$$
\mathcal{D}\left(A^{\star}\right)=\left\{g \in L^{2}(\mu): \text { there is an } h \in L^{2}(\mu) \text { such that }\langle A f, g\rangle=\langle f, h\rangle \text { for all } f \in \mathcal{D}(A)\right\},
$$

by $A^{\star} g=h$, with $h$ the function as described in the domain of the adjoint.
Define an operator $\left(L^{\times}, \mathcal{D}\left(L^{\times}\right)\right.$) by

$$
\mathcal{D}\left(L^{\times}\right)=\left\{g \in L^{2}(\mu) \cap W_{\mathrm{loc}}^{1,2}(E):(x, \theta) \mapsto-\theta \partial_{x} g(x, \theta)+\lambda(x,-\theta)(g(x,-\theta)-g(x, \theta)) \in L^{2}(\mu)\right\}
$$

and

$$
L^{\times} g(x, \theta)=-\theta \partial_{x} g(x, \theta)+\lambda(x,-\theta)(g(x,-\theta)-g(x, \theta)), \quad g \in \mathcal{D}\left(L^{\times}\right),(x, \theta) \in E .
$$

This operator is going to be our candidate for the adjoint operator $L^{\star}$.
Lemma 3.4. $\left(L^{\star}, \mathcal{D}\left(L^{\star}\right)\right)=\left(L^{\times}, \mathcal{D}\left(L^{\star}\right)\right)$.

Proof. We first show that $\mathcal{D}\left(L^{\star}\right) \subset \mathcal{D}\left(L^{\star}\right)$, and that $L^{\star}$ conincides with $L^{\star}$ on its domain. Suppose $g \in \mathcal{D}\left(L^{\star}\right)$, and let $h \in L^{2}(\mu)$ be so that $\langle L f, g\rangle=\langle f, h\rangle$ for every $f \in \mathcal{D}(L)$. Note that $C_{c}^{\infty}(E) \subset \mathcal{D}(L)$. For $f \in C_{c}^{\infty}(E)$ and $\theta= \pm 1$, we have

$$
\begin{aligned}
\sum_{\theta= \pm 1} \int_{\mathbb{R}} \theta \partial_{x} f(x, \theta) \overline{g(x, \theta)} e^{-U(x)} d x & =\left\langle\theta \partial_{x} f, g\right\rangle \\
& =\langle(L f-\lambda(F f-f)), g\rangle=\langle f, h\rangle-\langle\lambda(F f-f), g\rangle \\
& =\sum_{\theta= \pm 1} \int_{\mathbb{R}} f \overline{(h-F(\lambda g)+\lambda g)} e^{-U} d x .
\end{aligned}
$$

Since this holds in particular for any $f$ such that either $f(\cdot,+1)=0$ or $f(\cdot,-1)=0$, it follows that $\theta g e^{-U}$ has a weak derivative $-(h-F(\lambda g)+\lambda g) e^{-U}$. From this we may conclude (by Lemma A.1) that $g=\left(\theta g e^{-U}\right)\left(\theta e^{U}\right) \in W_{\mathrm{loc}}^{1,2}(E)$ with weak derivative

$$
\partial_{x} g=-\theta(h-F(\lambda g)+\lambda g)+U^{\prime} g .
$$

Isolating $h$ on the left hand side gives

$$
\begin{aligned}
h(x, \theta) & =-\theta \partial_{x} g(x, \theta)+\lambda(x,-\theta) g(x,-\theta)-\lambda(x, \theta) g(x, \theta)+\theta U^{\prime}(x) g(x, \theta) \\
& =-\theta \partial_{x} g(x, \theta)+\lambda(x,-\theta)(g(x,-\theta)-g(x, \theta)),
\end{aligned}
$$

using that $\theta U^{\prime}=\lambda-F \lambda$. Since $h \in L^{2}(\mu)$ the right-hand side is also in $L^{2}(\mu)$. So $g \in \mathcal{D}\left(L^{\times}\right)$and $h=L^{\times} g$.
Conversely, if $g \in \mathcal{D}\left(L^{\times}\right)$and $f \in \mathcal{D}(L)$, we may approximate $f$ by $\tilde{f} \in C_{c}^{\infty}(E)$. By partial integration it follows that $\langle L \tilde{f}, g\rangle=\left\langle\tilde{f}, L^{\times} g\right\rangle$. Since $\tilde{f}$ can be arbitrarily close to $f$ in the graph norm of $L$ by Lemma 3.2, it follows that $\langle L f, g\rangle=\left\langle f, L^{\times} g\right\rangle$. This establishes that $\mathcal{D}\left(L^{\times}\right) \subset \mathcal{D}\left(L^{\star}\right)$.

The following lemma establishes that $L^{\times}=F L F$.
Lemma 3.5. $f \in \mathcal{D}(L)$ if and only if $F f \in \mathcal{D}\left(L^{\times}\right)$, and $F L f=L^{\times}$Ff for $f \in \mathcal{D}(L)$. In particular $L^{\star}=L^{\times}$is dissipative.

Proof. If $f \in \mathcal{D}(L)$, then $f \in W_{\text {loc }}^{1,2}(E)$, and therefore $F f \in W_{\text {loc }}^{1,2}(E)$. Applying $L^{\times}$(formally) to $F f$ yields

$$
L^{\times} F f(x, \theta)=-\theta \partial_{x} f(x,-\theta)+\lambda(x, \theta)(f(x, \theta)-f(x,-\theta))=F L f(x, \theta) .
$$

Since $L f \in L^{2}(\mu)$, it follows that $F L f \in L^{2}(\mu)$. This establishes that $F f \in \mathcal{D}\left(L^{\times}\right)$and $F L f=L^{\times} F f$ for $f \in \mathcal{D}(L)$. The reverse inclusion is analogous. Since $L$ is dissipative by Lemma 3.3, this applies to $L^{\times}$.

We summarize our findings in the following theorem.
Theorem 3.6. $L$ is a closed dissipative operator which generates a strongly continuous contraction semigroup $(P(t))_{t \geq 0}$ in $L^{2}(\mu)$. The adjoint of $L$ is given by $L^{\star}=F L F$, i.e., $\mathcal{D}\left(L^{\star}\right)=\{F f: f \in \mathcal{D}(L)\}$ and $L^{\star} g=F L F g$ for $g \in \mathcal{D}\left(L^{\star}\right)$. The adjoint semigroup $\left(P^{\star}(t)\right)_{t \geq 0}$ is given by $P^{\star}(t)=F P(t) F$ for $t \geq 0$.

Proof. The fact that $L$ generates a strongly continuous semigroup follows by [10, Theorem II.3.17], using that $L$ and $L^{\star}$ are dissipative. In particular the range of $\gamma-L$ is equal to $L^{2}(\mu)$ ([10, Theorem II.3.15]) for any $\gamma>0$. The statement on the adjoints of $L^{\star}$ and $P^{\star}$ follows by Lemma 3.4.

### 3.2. Compactness of the resolvent of the zigzag semigroup generator

We will now investigate conditions under which the resolvent of the generator of the zigzag semigroup is compact. This is of fundamental importance in the characterization of the spectrum in subsequent sections.

Let $W^{1,2}(\nu)$ denote the usual Sobolev space of functions in $L^{2}(\nu)$ which have a weak derivative in $L^{2}(\nu)$. Since $W^{1,2}(\nu) \subset W_{\text {loc }}^{1,2}(\mathbb{R})$, all $f \in W^{1,2}(\nu)$ are absolutely continuous. Finally we denote by $W^{1,2}(\mu)$ the (equivalence classes of) measurable functions $f: E \rightarrow \mathbb{R}$ such that $f^{\theta} \in W^{1,2}(\nu)$ for $\theta= \pm 1$.

We may obtain a better understanding of $\mathcal{D}(L)$ if we suppose Assumption A1 and A2 are satisfied. In particular, $e^{-U}$ is assumed to belong to $L^{1}(\mathbb{R})$ so that $v$ and $\mu$ are finite measures. The relevance of these assumptions follows for a large part from the following lemmas.

Lemma 3.7 ([18, Lemma 8.5.2, Theorem 8.5.3]). Suppose that Assumption A1 holds. Then for any $p \in[2,+\infty)$ there exists a positive constant $C$ such that

$$
\int_{\mathbb{R}}|f|^{p}\left|U^{\prime}\right|^{2} d \nu \leq C\|f\|_{W^{1, p}(\nu)}^{p}, \quad f \in W^{1, p}(\nu)
$$

If additionally Assumption $A 2$ is satisfied, then the embedding $W^{1, p}(\nu) \subset L^{p}(\nu)$ is compact for any $p \in[2,+\infty)$.
Lemma 3.8. Suppose that $U$ satisfies Assumption A1. Suppose $f \in W_{\mathrm{loc}}^{1,2}(\mathbb{R}) \cap L^{2}(\nu)$ is such that $f^{\prime}-U^{\prime} f \in L^{2}(\nu)$. Then $f \in W^{1,2}(\nu)$ and

$$
\begin{equation*}
\int_{\mathbb{R}}\left|f^{\prime}\right|^{2} d \nu \leq-m \int_{\mathbb{R}}|f|^{2} d \nu+\int_{\mathbb{R}}\left|f^{\prime}-U^{\prime} f\right|^{2} d \nu<\infty \tag{3.5}
\end{equation*}
$$

Proof. First suppose $f \in C_{c}^{\infty}(\mathbb{R})$. We compute

$$
\int_{\mathbb{R}}\left|f^{\prime}-U^{\prime} f\right|^{2} d \nu=\int_{\mathbb{R}}\left\{\left|f^{\prime}\right|^{2}-f^{\prime} U^{\prime} \bar{f}-U^{\prime} f\left(\overline{f^{\prime}}\right)+\left|U^{\prime}\right|^{2}|f|^{2}\right\} d \nu .
$$

For the cross terms we find by partial integration that

$$
\begin{aligned}
\int_{\mathbb{R}}\left\{-\left(f^{\prime}\right) U^{\prime} \bar{f}-U^{\prime} f\left(\overline{f^{\prime}}\right)\right\} d \nu & =-\int_{\mathbb{R}}\left(\partial_{x}|f|^{2}\right) U^{\prime} e^{-U} d x \\
& =\int_{\mathbb{R}}\left\{|f|^{2} U^{\prime \prime}-\left|U^{\prime}\right|^{2}|f|^{2}\right\} e^{-U} d x
\end{aligned}
$$

Thus,

$$
\int_{\mathbb{R}}\left|f^{\prime}-U^{\prime} f\right|^{2} d \nu=\int_{\mathbb{R}}\left\{\left|f^{\prime}\right|^{2}+|f|^{2} U^{\prime \prime}\right\} d \nu
$$

Combining with the assumption that $U^{\prime \prime} \geq m$ and rearranging yields the stated conclusion for $f \in C_{c}^{\infty}(\mathbb{R})$.
Analogous to the proof of Lemma 3.2, it may be established that $C_{c}^{\infty}(\mathbb{R})$ is a core for the closed operator $(C, \mathcal{D}(C))$ on $L^{2}(\nu)$ with $C f=f^{\prime}-U^{\prime} f$ and $\mathcal{D}(C)=\left\{f \in L^{2}(\nu) \cap W_{\text {loc }}^{1,2}: C f \in L^{2}(\nu)\right\} ;$ see Lemma A. 3 for details. Letting $f_{n} \rightarrow f$ in the graph norm of $C$ with $\left(f_{n}\right) \subset C_{c}^{\infty}(\mathbb{R})$, we find that, using Fatou's lemma, and after taking a subsequence for which $f_{n}^{\prime}$ converges almost everywhere,

$$
\begin{aligned}
\left\|f^{\prime}\right\|_{L^{2}(\nu)}^{2} & =\left\|\lim _{n \rightarrow \infty} f_{n}^{\prime}\right\|_{L^{2}(\nu)}^{2} \\
& \leq \liminf _{n \rightarrow \infty}\left\|f_{n}^{\prime}\right\|_{L^{2}(\nu)}^{2} \\
& \leq \liminf _{n \rightarrow \infty}\left(-m\left\|f_{n}\right\|_{L^{2}(\nu)}^{2}+\left\|C f_{n}\right\|_{L^{2}(\nu)}^{2}\right) \\
& =-m\|f\|_{L^{2}(\nu)}^{2}+\|C f\|_{L^{2}(v)}^{2}<\infty
\end{aligned}
$$

We can now state sufficient conditions for the resolvent of the generator of the zigzag semigroup to be compact.
Theorem 3.9 (Compact resolvent). Suppose Assumption A1 is satisfied. Then $\mathcal{D}(L)=W^{1,2}(\mu)$. If also Assumption A2 holds then the embedding $W^{1,2}(\mu) \subset L^{2}(\mu)$ is compact and in particular, for any $\gamma>0$, the resolvent operator $(\gamma-$ $L)^{-1}: L^{2}(\mu) \rightarrow L^{2}(\mu)$ is compact.

Proof. First we prove that $\mathcal{D}(L) \subset W^{1,2}(\mu)$. Suppose $f \in \mathcal{D}(L)$ and write $L f=h$. Write $g_{1}=f^{+}-f^{-}$and $g_{2}=$ $f^{+}+f^{-}$. We have

$$
\begin{aligned}
& \partial_{x} g_{1}=h^{+}-\lambda^{+}\left(f^{-}-f^{+}\right)+h^{-}-\lambda^{-}\left(f^{+}-f^{-}\right)=h^{+}+h^{-}+U^{\prime} g_{1}, \\
& \partial_{x} g_{2}=h^{+}-\lambda^{+}\left(f^{-}-f^{+}\right)-h^{-}+\lambda^{-}\left(f^{+}-f^{-}\right)=h^{+}-h^{-}-\left(2 \lambda_{\text {refr }}+\left|U^{\prime}\right|\right) g_{1} .
\end{aligned}
$$

We see that $g_{1}$ satisfies $g_{1}^{\prime}-U^{\prime} g_{1}=h^{+}+h^{-} \in L^{2}(\nu)$ so that, by Lemma 3.8, $g_{1} \in W^{1,2}(\nu)$. Next by the estimate on $\lambda_{\text {refr }}$ of Assumption A1, Lemma 3.7 and the second equation, $g_{2} \in W^{1,2}(\nu)$. The inclusion $\mathcal{D}(L) \subset W^{1,2}(\mu)$ follows since $f^{ \pm}$ may be expressed as linear combinations of $g_{1}$ and $g_{2}$.

Conversely, we now prove that $W^{1,2}(\mu) \subset \mathcal{D}(L)$. Since $\lambda_{\text {refr }}(x) \leq c_{1}+c_{2}\left|U^{\prime}(x)\right|$ for some constants $c_{1}, c_{2}$, it follows $\stackrel{\sim}{\text { that }}|\lambda(x, \theta)| \leq C_{1}+C_{2}\left|U^{\prime}(x)\right|$ for some constants $C_{1}, C_{2}$. By Assumption A1 and Lemma 3.7, $\|f\|_{L^{2}(\mu)}+\|L f\|_{L^{2}(\mu)} \leq$ $\widetilde{C}\|f\|_{W^{1,2}(\mu)}$ for some $\widetilde{C}>0$, which establishes that $W^{1,2}(\mu) \subset \mathcal{D}(L)$.

Let $\gamma>0$. Recall that $\gamma-L: \mathcal{D}(L) \rightarrow L^{2}(\mu)$ is bijective, so the inverse map exists, and is bounded by dissipativity (see [10, Proposition II.3.14]. Since the embedding $W^{1,2} \hookrightarrow L^{2}(\mu)$ is compact by Lemma 3.7 it follows from [10, Proposition II.4.25] that the resolvent is compact.

Remark 3.10. If Assumption A 1 is satisfied, then (3.4) holds with equality, i.e., for all $f \in \mathcal{D}(L)$ we have

$$
\operatorname{Re} \int_{E}(L f) \bar{f} d \mu=-\frac{1}{2} \int_{\mathbb{R}}\left(\lambda^{+}+\lambda^{-}\right)\left|f^{+}-f^{-}\right|^{2} d \nu, \quad f \in \mathcal{D}(L) .
$$

This follows from the proof of Lemma 3.3. Indeed in the proof the above relation is established for $f \in C_{c}^{\infty}(E)$. Now for $f \in \mathcal{D}(L)=W^{1,2}(\mu)$, take a sequence $\left(f_{n}\right) \subset C_{c}^{\infty}(E)$ converging in $W^{1,2}(\mu)$ towards $f$, and note that both sides of the above equation converge, using Lemma 3.7.

The fact that the resolvent is compact under suitable conditions leads naturally to the question whether the semigroup is eventually compact. We will now establish that this is not the case. The following lemma first states in a technical way that the zigzag semigroup consists simply of translations in regions of the space where the switching intensity vanishes, a necessary ingredient for the conclusion of Proposition 3.12 that the zigzag semigroup is in general not eventually compact.

Lemma 3.11. Suppose $\lambda_{\text {refr }}(x)=0$ for all $x$. Suppose $\theta U^{\prime}(x) \leq 0$ for some $\theta \in\{-1,+1\}$ and all $x \in I$, where $I=[a, b]$ is a non-empty interval. Then $P(t) f(x, \theta)=f(x+\theta t, \theta)$ for all $t \geq 0$ and $\mu$-almost all $x \in \mathbb{R}$ for which $\{x+\theta s: 0 \leq$ $s \leq t\} \subset I$.

Proof. First suppose $f \in \mathcal{D}(L)$. Write $\varphi(t, x):=P(t) f(x, \theta)$ with $\theta$ as in the assumption of the lemma. Then, for $x \in I$

$$
\partial_{t} \varphi(t, x)=\partial_{t} P(t) f(x, \theta)=L P(t) f(x, \theta)=\theta \partial_{x} P(t) f(x, \theta)=\theta \partial_{x} \varphi(t, x) .
$$

Fix ( $x, t$ ) such that $x+\theta s \in[a, b]$ for $0 \leq s \leq t$. Consider the characteristic curve $X(s)=x+\theta s, T(s)=t-s$. We find that, since $X(s) \in I$ for all $0 \leq s \leq t$, then $\frac{d}{d s} \varphi(T(s), X(s))=0$, so that

$$
\varphi(t, x)=\varphi(T(0), X(0))=\varphi(T(t), X(t))=\varphi(0, x+\theta t)=f(x+\theta t) .
$$

Now if $f_{n} \rightarrow f$ in $L^{2}(\mu)$ with $\left(f_{n}\right) \subset \mathcal{D}(L)$, then $P(t) f_{n_{k}}(x, \theta) \rightarrow P(t) f(x, \theta), \mu$-almost all $x$, for a subsequence $f_{n_{k}}$, and $f_{n_{k_{l}}}(x+\theta t, \theta) \rightarrow f(x+\theta t, \theta), \mu$-almost all $x$, for a further subsequence, which yields the result for general $f$.

Proposition 3.12. Suppose Assumption A3 is satisfied. Furthermore assume that $\lambda_{\mathrm{refr}}(x)=0$ everywhere. Then, for any $t>0$, the operator $P(t)$ is not compact.

Proof. For $n \in \mathbb{N}$ define functions $f_{n}$ by

$$
\tilde{f}_{n}^{+}(x)=\mathbb{1}_{[-1 / n,-1 /(n+1)]}(x), \quad \widetilde{f}_{n}^{-}(x)=0, \quad \widetilde{f}_{n}(\cdot, \pm 1)=\tilde{f}_{n}^{ \pm}, \quad f_{n}=\widetilde{f}_{n} /\left\|\tilde{f}_{n}\right\|_{L^{2}(\mu)}
$$

Then $\left(f_{n}\right)$ is a bounded sequence of functions in $L^{2}(\mu)$. Therefore if $P\left(t_{0}\right)$ were compact for some $t_{0}>0$, it should be possible to find a convergent subsequence of $\left(P\left(t_{0}\right) f_{n}\right)_{n=1}^{\infty}$. However we will show that this is impossible.

By Assumption A3, $U^{\prime}(x) \leq 0$ for $x \leq 0$. By Lemma 3.11,

$$
P\left(t_{0}\right) f(x,+1)=f\left(x+t_{0},+1\right) \quad \text { if } x \leq-t_{0}, \text { for any } f \in L^{2}(\mu) .
$$

Suppose $m \neq n$. Then

$$
\begin{aligned}
\left\|P\left(t_{0}\right) f_{n}-P\left(t_{0}\right) f_{m}\right\|_{L^{2}(\mu)}^{2} & \geq \int_{-\infty}^{-t_{0}}\left|P\left(t_{0}\right) f_{n}(x,+1)-P\left(t_{0}\right) f_{m}(x,+1)\right|^{2} e^{-U(x)} d x \\
& =\int_{-\infty}^{-t_{0}}\left|f_{n}^{+}\left(x+t_{0}\right)-f_{m}^{+}\left(x+t_{0}\right)\right|^{2} e^{-U(x)} d x \\
& =\int_{-\infty}^{0}\left(\left|f_{n}^{+}(y)\right|^{2}+\left|f_{m}^{+}(y)\right|^{2}\right) e^{-U\left(y-t_{0}\right)} d y \\
& =\frac{\int_{-1 / n}^{-1 /(n+1)} e^{-U\left(y-t_{0}\right)} d y}{\int_{-1 / n}^{-1 /(n+1)} e^{-U(y)} d y}+\frac{\int_{-1 / m}^{-1 /(m+1)} e^{-U\left(y-t_{0}\right)} d y}{\int_{-1 / m}^{-1 /(m+1)} e^{-U(y)} d y} \\
& \geq 2 e^{-U\left(-1-t_{0}\right)+U(0)} .
\end{aligned}
$$

This establishes that no subsequence of $\left(P\left(t_{0}\right) f_{n}\right)_{n \in \mathbb{N}}$ is Cauchy, and therefore that $P\left(t_{0}\right)$ is not compact.

## 4. Spectral theory of the zigzag generator

This section will describe key results concerning the characterization of the spectrum of the zigzag semigroup under the assumptions of unimodularity and no refreshments. The main result is Theorem 4.4 which, in summary, gives a description of the spectrum of $L$ as consisting of the roots of a particular complex valued function.

As usual define the resolvent set $\rho(A)$ of a closed operator $A$ by

$$
\rho(A)=\left\{\gamma \in \mathbb{C}: \gamma-A \text { is invertible and }(\gamma-A)^{-1} \text { is bounded }\right\}
$$

and the spectrum $\sigma(A)$ by $\sigma(A)=\mathbb{C} \backslash \rho(A)$. The point spectrum of $A$ is defined as

$$
\sigma_{p}(A)=\{\gamma \in \mathbb{C}: \gamma-A \text { is not injective }\} .
$$

Elements $\gamma \in \sigma_{p}(A)$ are called eigenvalues. For a disjoint decomposition of the spectrum $\sigma(A)=\sigma_{c} \cup \sigma_{u}$ where $\sigma_{u}$ is closed and $\sigma_{c}$ is compact, the spectral projection $P_{c}$ is defined as

$$
P_{c}=\frac{1}{2 \pi \mathrm{i}} \int_{\Gamma}(\zeta-A)^{-1} d \zeta
$$

where $\Gamma$ is a Jordan curve in the complement of $\sigma_{u}$ and enclosing $\sigma_{c}$. See [10, Section IV.1] or [17, Section III.6.4] for details. If $\Gamma$ encloses only a single element $\gamma \in \sigma(A)$ then we write $P_{\gamma}=P_{c}$.

We start with a simple corollary of Theorem 3.9.
Corollary 4.1. Suppose Assumptions A1 and A2 are satisfied. Then the spectrum of $L$ consists entirely of isolated eigenvalues with finite multiplicities, and the resolvent operator $(\gamma-L)^{-1}$ is compact for every $\gamma \in \rho(L)$.

Proof. This follows from Theorem 3.9 and [17, Theorem III.6.29].
There are also some interesting immediate consequences of the fact that $L^{\star}=F L F$ which we state in some generality in Lemma 4.2. Let $(\mathcal{H},\langle\cdot, \cdot\rangle$,$) denote a Hilbert space, and let J$ be a unitary operator on $\mathcal{H}$. We call a closed operator $(A, \mathcal{D}(A))$ on $\mathcal{H} J$-selfadjoint if $A^{\star}=J A J^{\star}$. In the case of the zigzag process we have that $L$ is $J$-selfadjoint for $J=F$.

Lemma 4.2. Suppose $(A, \mathcal{D}(A))$ is a closed operator on $\mathcal{H}$ such that $A$ is $J$-selfadjoint.
(i) $\rho(A)=\rho\left(A^{\star}\right)$.
(ii) $\sigma(A)=\{\bar{\gamma}: \gamma \in \sigma(A)\}$ and $\sigma_{p}(A)=\left\{\bar{\gamma}: \gamma \in \sigma_{p}(A)\right\}$ (i.e. $\sigma(A)$ and $\sigma_{p}(A)$ are closed under taking complex conjugate).
(iii) If $\gamma_{1}, \gamma_{2}, \gamma_{1} \neq \gamma_{2}$, are eigenvalues of A with eigenvectors $\phi_{1}, \phi_{2}$, respectively, then either $\gamma_{1}=\bar{\gamma}_{2}$ or $\phi_{1}$ and $\phi_{2}$ are $J$-orthogonal, i.e., $\left\langle J \phi_{1}, \phi_{2}\right\rangle=0$.
(iv) If $\phi$ is an eigenvector of $A$ with eigenvalue $\gamma$, then $J \phi$ is an eigenvector of $A^{\star}$ with eigenvalue $\gamma$.
(v) The adjoint of the spectral projection $P_{\gamma}$ associated to an isolated eigenvalue $\gamma \in \sigma(A)$ satisfies

$$
\left(P_{\gamma}\right)^{\star}=J P_{\bar{\gamma}} J^{\star}
$$

and is equal to the spectral projection associated with $A^{\star}$ for the eigenvalue $\bar{\gamma}$.

## Proof.

(i) This follows from $\left(\gamma-A^{\star}\right)^{-1}=J(\gamma-A)^{-1} J^{\star}$ for $\gamma \in \rho(A)$.
(ii) For general closed operators $A$ on a Hilbert space, we have that $\sigma\left(A^{\star}\right)$ equals the complex conjugate of $\sigma(A)$, and similarly for $\sigma_{p}(A)$ (e.g., [17, Theorem III.6.22]). Combining with the above equality used in the proof of (i) gives the stated result.
(iii) This follows from

$$
\gamma_{1}\left\langle J \phi_{1}, \phi_{2}\right\rangle=\left\langle J A \phi_{1}, \phi_{2}\right\rangle=\left\langle A^{\star} J \phi_{1}, \phi_{2}\right\rangle=\left\langle J \phi_{1}, A \phi_{2}\right\rangle=\bar{\gamma}_{2}\left\langle J \phi_{1}, \phi_{2}\right\rangle .
$$

(iv) If $\phi$ is an eigenvector of $A$ with eigenvalue $\gamma$, then $A^{\star} J \phi=J A \phi=\gamma J \phi$.
(v) Let $\bar{\Gamma}$ denote a counterclockwise contour in the complex plane around (only) the eigenvalue $\bar{\gamma}$. By [17, Theorem III.6.22], the adjoint of the corresponding spectral projection satisfies

$$
\left(P_{\gamma}\right)^{\star}=\frac{1}{2 \pi \mathrm{i}} \int_{\bar{\Gamma}}\left(\zeta-A^{\star}\right)^{-1} d \zeta=\frac{1}{2 \pi \mathrm{i}} J \int_{\bar{\Gamma}}(\zeta-A)^{-1} d \zeta J^{\star}=J P_{\bar{\gamma}} J^{\star} .
$$

For later reference we note the following general observation for eigenfunctions of linear operators on function spaces.
Lemma 4.3. Suppose $(A, \mathcal{D}(A))$ is a closed operator on $L^{2}(\mu)$. Suppose

$$
\psi \in \mathcal{D}(A) \quad \Longrightarrow \quad \operatorname{Re} \psi \in \mathcal{D}(A) \quad \text { and } \quad \operatorname{Im} \psi \in \mathcal{D}(A)
$$

and suppose A maps (a.e.) real valued functions in its domain into (a.e.) real valued functions. Suppose $\phi$ is an eigenfunction of $A$ with eigenvalue $\gamma$. Then $\bar{\phi}$ is an eigenfunction of $A$ with eigenvalue $\bar{\gamma}$.

Proof. By linearity of $A$,

$$
\overline{A \bar{\phi}}=\overline{A \operatorname{Re} \phi-\mathrm{i} A \operatorname{Im} \phi}=A \operatorname{Re} \phi+\mathrm{i} A \operatorname{Im} \phi=A \phi=\gamma \phi,
$$

which establishes the result.
We are now ready to state the main result of this section, which characterizes $\sigma(L)$. Define functions $\psi^{ \pm}(\gamma):=$ $\int_{0}^{ \pm \infty} U^{\prime}(\xi) e^{\mp 2 \gamma \xi-U(\xi)} d \xi$, i.e.,

$$
\begin{equation*}
\psi^{+}(\gamma):=\int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi \quad \text { and } \quad \psi^{-}(\gamma):=-\int_{-\infty}^{0} U^{\prime}(\xi) e^{2 \gamma \xi-U(\xi)} d \xi, \quad \gamma \in \mathbb{C}, \tag{4.1}
\end{equation*}
$$

and

$$
\begin{equation*}
Z(\gamma):=1-\psi^{+}(\gamma) \psi^{-}(\gamma), \quad \gamma \in \mathbb{C} . \tag{4.2}
\end{equation*}
$$

Define the set

$$
\begin{equation*}
\Sigma:=\{\gamma \in \mathbb{C}: Z(\gamma)=0\} . \tag{4.3}
\end{equation*}
$$

Under Assumption A2 the functions $\psi^{ \pm}$are well defined, using Lemma A.4.

Theorem 4.4 (Spectrum, resolvent and eigenfunctions of the zigzag generator). Suppose Assumptions A1, A2 and A3 are satisfied. Furthermore assume $\lambda_{\text {refr }}(x)=0$ for all $x$. Then $\sigma(L)=\sigma_{p}(L)=\Sigma$.

If $\gamma \in \rho(L)$ and $h \in L^{2}(\mu)$ then $f=(\gamma-L)^{-1} h$ is given by

$$
\begin{align*}
& f^{+}(x)= \begin{cases}e^{\gamma x}\left(k^{+}+\int_{x}^{0} e^{-\gamma \xi} h^{+}(\xi) d \xi\right), & x \leq 0, \\
e^{\gamma x+U(x)}\left(k^{+}-\int_{0}^{x} e^{-\gamma \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) e^{-\gamma \xi}\left(k^{-}+\int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta\right)\right] d \xi\right), & x>0,\end{cases} \\
& f^{-}(x)= \begin{cases}e^{-\gamma x+U(x)}\left(k^{-}-\int_{x}^{0} e^{\gamma \xi-U(\xi)}\left[h^{-}(\xi)-U^{\prime}(\xi) e^{\gamma \xi}\left(k^{+}+\int_{\xi}^{0} e^{-\gamma \eta} h^{+}(\eta) d \eta\right)\right] d \xi\right), & x \leq 0, \\
e^{-\gamma x}\left(k^{-}+\int_{0}^{x} e^{\gamma \xi} h^{-}(\xi) d \xi\right), & x>0 .\end{cases} \tag{4.4}
\end{align*}
$$

The constants $k^{ \pm}=k^{ \pm}(\gamma ; h)$ are given by

$$
\binom{k^{+}(\gamma ; h)}{k^{-}(\gamma ; h)}=\frac{1}{Z(\gamma)}\left(\begin{array}{cc}
1 & \psi^{+}(\gamma)  \tag{4.5}\\
\psi^{-}(\gamma) & 1
\end{array}\right) K(\gamma) h,
$$

with $K(\gamma): L^{2}(\mu) \rightarrow \mathbb{C}^{2}$ the bounded linear mapping given by

$$
\begin{equation*}
K(\gamma) h=\binom{\int_{0}^{\infty} e^{-\gamma \xi-U(\xi)} h^{+}(\xi) d \xi+\int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} \int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta d \xi}{\int_{-\infty}^{0} e^{\gamma \xi-U(\xi)} h^{-}(\xi) d \xi-\int_{-\infty}^{0} U^{\prime}(\xi) e^{2 \gamma \xi-U(\xi)} \int_{\xi}^{0} e^{-\gamma \eta} h^{+}(\eta) d \eta d \xi} \tag{4.6}
\end{equation*}
$$

For every $\gamma \in \sigma_{p}(L)$, the corresponding eigenfunctions form a one-dimensional space spanned by the function $f_{\gamma} \in$ $\mathcal{D}(L)$ defined by

$$
\begin{aligned}
& f_{\gamma}^{+}(x)=\psi^{+}(\gamma) e^{\gamma x}, \quad(x \leq 0), \quad f_{\gamma}^{+}(x)=e^{\gamma x+U(x)} \int_{x}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi, \quad(x \geq 0), \\
& f_{\gamma}^{-}(x)=e^{-\gamma x}, \quad(x \geq 0), \quad f_{\gamma}^{-}(x)=-\psi^{+}(\gamma) e^{-\gamma x+U(x)} \int_{-\infty}^{x} U^{\prime}(\xi) e^{2 \gamma \xi-U(\xi)} d \xi, \quad(x \leq 0) .
\end{aligned}
$$

If $\gamma_{1}, \gamma_{2} \in \sigma_{p}(L), \gamma_{1} \neq \gamma_{2}$, with corresponding eigenfunctions $f_{1}, f_{2}$ then either $\gamma_{1}=\bar{\gamma}_{2}$ or $\left\langle F f_{1}, f_{2}\right\rangle=0$.
Proof. Claim: $\Sigma \subset \sigma(L)$.
Proof of claim. We will prove the equivalent statement that $\rho(L) \subset \mathbb{C} \backslash \Sigma$. Suppose $\gamma \in \rho(L)$. In particular, for all $h \in L^{2}(\mu)$, there is a solution $f \in \mathcal{D}(L)$ to the equation $(\gamma-L) f=h . f$ satisfies the coupled system of ordinary differential equations

$$
\begin{aligned}
& \gamma f^{+}-\partial_{x} f^{+}-\lambda^{+}\left(f^{-}-f^{+}\right)=h^{+} \\
& \gamma f^{-}+\partial_{x} f^{-}-\lambda^{-}\left(f^{+}-f^{-}\right)=h^{-}
\end{aligned}
$$

For $x \leq 0$, we have $\lambda^{+}(x)=0$ and $\lambda^{-}(x)=-U^{\prime}(x)$. We find that

$$
f^{+}(x)=e^{\gamma x}\left(k^{+}+\int_{x}^{0} e^{-\gamma \xi} h^{+}(\xi) d \xi\right), \quad x \leq 0,
$$

for some constant $k^{+}$. By variation of constants, for $x \leq 0$,

$$
\begin{aligned}
f^{-}(x) & =e^{-\gamma x+U(x)}\left(k^{-}-\int_{x}^{0} e^{\gamma \xi-U(\xi)}\left[h^{-}(\xi)+\lambda^{-}(\xi) f^{+}(\xi)\right] d \xi\right) \\
& =e^{-\gamma x+U(x)}\left(k^{-}-\int_{x}^{0} e^{\gamma \xi-U(\xi)}\left[h^{-}(\xi)-U^{\prime}(\xi) e^{\gamma \xi}\left(k^{+}+\int_{\xi}^{0} e^{-\gamma \eta} h^{+}(\eta) d \eta\right)\right] d \xi\right) .
\end{aligned}
$$

Similarly, for $x \geq 0$, we obtain

$$
\begin{aligned}
& f^{-}(x)=e^{-\gamma x}\left(k^{-}+\int_{0}^{x} e^{\gamma \xi} h^{-}(\xi) d \xi\right) \\
& f^{+}(x)=e^{\gamma x+U(x)}\left(k^{+}-\int_{0}^{x} e^{-\gamma \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) f^{-}(\xi)\right] d \xi\right)
\end{aligned}
$$

$$
=e^{\gamma x+U(x)}\left(k^{+}-\int_{0}^{x} e^{-\gamma \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) e^{-\gamma \xi}\left(k^{-}+\int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta\right)\right] d \xi\right)
$$

Here no new integration constants are introduced because we required continuity of $f$. Indeed by assumption $f \in \mathcal{D}(L)$ and thus $f$ is continuous.

By Lemma A.4, we have that all the integrals occuring in the expressions for $f^{+}(x)$ for $x \geq 0$ and (analogously) $f^{-}(x)$ for $x \leq 0$, converge as $|x| \rightarrow \infty$. In order for $f \in L^{2}(\mu)$, it follows that necessarily

$$
\begin{align*}
& k^{+}-\int_{0}^{\infty} e^{-\gamma \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) e^{-\gamma \xi}\left(k^{-}+\int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta\right)\right] d \xi=0 \quad \text { and } \\
& k^{-}-\int_{-\infty}^{0} e^{\gamma \xi-U(\xi)}\left[h^{-}(\xi)-U^{\prime}(\xi) e^{\gamma \xi}\left(k^{+}+\int_{\xi}^{0} e^{-\gamma \eta} h^{+}(\eta) d \eta\right)\right] d \xi=0 \tag{4.7}
\end{align*}
$$

for otherwise the multiplication by $e^{ \pm \gamma x+U(x)}$ in the expression for $f$ would lead to the norm of $f$ being equal to $\infty$. (Note that the exponential growth by $e^{ \pm \operatorname{Re} \gamma x}$ in $f^{+}$for $x \geq 0$ and $f^{-}$for $x \leq 0$ is not problematic since it is compensated by $e^{-U(x)}$ in the $L^{2}(\mu)$ inner product, which by Assumption A2 will decay faster than any exponential.) Condition (4.7) can be written in terms of a linear system for $k^{ \pm}$:

$$
\left(\begin{array}{cc}
1 & -\psi^{+}(\gamma)  \tag{4.8}\\
-\psi^{-}(\gamma) & 1
\end{array}\right)\binom{k+}{k^{-}}=K(\gamma) h
$$

where, for $\gamma \in \mathbb{C}, K(\gamma): L^{2}(\mu) \rightarrow \mathbb{C}^{2}$ is defined by (4.6). The boundedness of $K$ follows by the proof of Lemma A.4.
Now suppose $\gamma \in \Sigma$. In this case the system (4.8) is singular. However the linear map $K(\gamma): L^{2}(\mu) \rightarrow \mathbb{C}^{2}$ is surjective, even when we restrict the domain to the functions in $L^{2}(\mu)$ having compact support. Indeed, we can produce the vector $(10)^{T}$ on the right hand side by choosing

$$
h^{-}(x)=0 \quad \text { and } \quad h^{+}(x)=e^{\gamma x+U(x)} \mathbb{1}_{[0,1]}(x), \quad x \in \mathbb{R}
$$

and in an analogous way we can choose $h$ so that the vector $(01)^{T}$ is obtained on the right hand side. Thus the right hand side as a function of $h$ having compact support has range $\mathbb{C}^{2}$, whereas the left hand side can only span a one-dimensional subspace due to the assumption $\gamma \in \Sigma$. It follows that there are choices $h$ such that no solution to the resolvent system $(\gamma-L) f=h$ exist, so that $\gamma \notin \rho(L)$; a contradiction. The solution of (4.8) yields the stated expression for $k^{ \pm}$.

Claim: $\sigma_{p}(L) \subset \Sigma$.
Proof of claim. Suppose $\gamma \in \sigma_{p}(L)$. There exists a function $f \in \mathcal{D}(L)$ such that $L f=\gamma f$. This corresponds to the system

$$
\begin{aligned}
& \gamma f^{+}-\partial_{x} f^{+}-\lambda^{+}\left(f^{-}-f^{+}\right)=0 \\
& \gamma f^{-}+\partial_{x} f^{-}-\lambda^{-}\left(f^{+}-f^{-}\right)=0
\end{aligned}
$$

Solving for $f^{+}$for $x \leq 0$ yields $f^{+}(x)=c^{+} e^{\gamma x}$, and similarly $f^{-}(x)=c^{-} e^{-\gamma x}$ for $x \geq 0$ for some constants $c^{ \pm} \in \mathbb{C}$. Then using variation of constants and insisting upon continuity in $x=0$, yields

$$
\begin{aligned}
& f^{+}(x)=e^{\gamma x+U(x)}\left(c^{+}-c^{-} \int_{0}^{x} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi\right), \quad x \geq 0 \\
& f^{-}(x)=e^{-\gamma x+U(x)}\left(c^{-}+c^{+} \int_{x}^{0} U^{\prime}(\xi) e^{2 \gamma \xi-U(\xi)} d \xi\right), \quad x \leq 0
\end{aligned}
$$

Again the integrals converge as $|x| \rightarrow \infty$ by Lemma A.4. Now in order for $f \in L^{2}(\mu)$, we require that the terms within parantheses in the expressions for $f^{ \pm}$vanish as $|x| \rightarrow \infty$. This yields the conditions

$$
\begin{equation*}
c^{+}=c^{-} \psi^{+}(\gamma), \quad \text { and } \quad c^{-}=c^{+} \psi^{-}(\gamma) \tag{4.9}
\end{equation*}
$$

This system admits a non-trivial solution if and only if $\gamma \in \Sigma$. The choice for $f$ in the statement of the proposition is obtained by taking $c^{-}=1$ and $c^{+}=\psi^{+}(\gamma)$.

We have established that $\Sigma \subset \sigma(L)$ and $\sigma_{p}(L) \subset \Sigma$. Since $L$ has compact resolvent (Corollary 4.1), $\sigma_{p}(L)=\sigma(L)$ and the proof is complete.

The final statement of the theorem follows by an application of Lemma 4.2(iii).

Corollary 4.5. The spectral projection corresponding to $\gamma \in \sigma(L)$ is given by

$$
P_{\gamma}=\frac{1}{2 \pi \mathrm{i}} \int_{\Gamma} \frac{1}{Z(\zeta)} M(\zeta)\left(\begin{array}{cc}
1 & \psi^{+}(\zeta)  \tag{4.10}\\
\psi^{-}(\zeta) & 1
\end{array}\right) K(\zeta) d \zeta
$$

where $\Gamma$ is a Jordan contour in $\mathbb{C}$ enclosing the eigenvalue $\gamma$ and no other eigenvalues, and where the family of bounded linear operators $(M(\zeta))_{\zeta \in \rho(L)}$ is given for fixed $\zeta \in \rho(L)$ as a bounded operator $M(\zeta): \mathbb{C}^{2} \rightarrow L^{2}(\mu)$ by

$$
\begin{aligned}
& (M(\zeta) a)^{+}(x):=M^{+}(\zeta) a(x):= \begin{cases}e^{\zeta x} a^{+}, & x \leq 0, \\
e^{\zeta x+U(x)}\left(a^{+}-a^{-} \int_{0}^{x} U^{\prime}(\xi) e^{-2 \zeta \xi-U(\xi)} d \xi\right), & x>0,\end{cases} \\
& (M(\zeta) a)^{-}(x):=M^{-}(\zeta) a(x):= \begin{cases}e^{-\zeta x+U(x)}\left(a^{-}+a^{+} \int_{x}^{0} U^{\prime}(\xi) e^{2 \zeta x-U(\xi)} d \xi\right), & x \leq 0, \\
e^{-\zeta x} a^{-}, & x>0,\end{cases}
\end{aligned}
$$

for $a=\binom{a^{+}}{a^{-}} \in \mathbb{C}^{2}$.
Note that the integral in (4.10) is a concatenation of bounded linear operators and as such the integrand is a well defined mapping from $\rho(L)$ into the bounded operators on $L^{2}(\mu)$, which makes the integral over closed contours well defined as a Bochner integral.

Proof. By [17, proof of Theorem 6.17], the spectral projections are given for $\gamma \in \sigma(L)$ by

$$
P_{\gamma}=\frac{1}{2 \pi \mathrm{i}} \int_{\Gamma}(\zeta-L)^{-1} d \zeta
$$

where $\Gamma$ is a contour enclosing $\gamma$ as stated. Note that the terms in the resolvent expression (4.4) which do not involve $k^{ \pm}$ amount to integral operators by an analytic integral kernel. We may check that such terms are holomorphic as operator valued functions of the complex-valued argument. Indeed, as discussed in [17, Section VII.1.1] weak and strong notions of holomorphicity coincide, and we may check that for example the operator $T^{+}(\zeta)$, constituting the non- $k^{ \pm}$-dependent terms in $\left((\zeta-L)^{-1} h\right)^{+}$, given by

$$
T^{+}(\zeta) h(x)= \begin{cases}e^{\zeta x} \int_{x}^{0} e^{-\zeta \xi} h^{+}(\xi) d \xi \\ e^{\zeta x+U(x)}\left(-\int_{0}^{x} e^{-\zeta \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) e^{-\zeta \xi} \int_{0}^{\xi} e^{\zeta \eta} h^{-}(\eta) d \eta\right]\right), & x>0\end{cases}
$$

is holomorphic as a function of $\zeta$ by verifying that $\left\langle g, T^{+}(\zeta) h\right\rangle_{L^{2}(\nu)}$ is holomorphic for indicator functions $g, h$. Analogously this can be verified for $\left((\zeta-L)^{-1} h\right)^{-}$.

Thus, if we perform the complex contour integration all terms in (4.4) which do not involve $k^{ \pm}$vanish, since such terms are holomorphic in $\zeta$. The constants $k^{ \pm}$are given by (4.5). Combining these observations yields the stated expression.

Corollary 4.6 (Simple roots). Suppose the assumptions of Theorem 4.4 are satisfied and $\gamma$ is a simple root of $Z$, so that $\gamma \in \sigma(L)$ and $Z^{\prime}(\gamma) \neq 0$. The corresponding spectral projection has rank one and is given by

$$
P_{\gamma} h=\frac{\left(1 \psi^{+}(\gamma)\right) K(\gamma) h}{Z^{\prime}(\gamma) \psi^{+}(\gamma)} f_{\gamma},
$$

where $f_{\gamma}$ is the eigenfunction corresponding to $\gamma$ as introduced in Theorem 4.4. An alternative expression for $P_{\gamma}$ is given by

$$
P_{\gamma} h=\frac{\left\langle h, F \bar{f}_{\gamma}\right\rangle}{\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle} f_{\gamma}
$$

Proof. Note that $Z(\gamma)=0$ yields $\psi^{-}(\gamma)=1 / \psi^{+}(\gamma)$. The Cauchy residue theorem applied to the expression for $P_{\gamma}$ of Corollary 4.5 yields that

$$
P_{\gamma}=\frac{1}{Z^{\prime}(\gamma)} M(\gamma)\binom{1}{\psi^{-}(\gamma)}\left(\begin{array}{ll}
1 & \left.\psi^{+}(\gamma)\right) K(\gamma) .
\end{array}\right.
$$

Next we compute

$$
\begin{aligned}
M^{+}(\gamma)\binom{1}{\psi^{-}(\gamma)}(x) & = \begin{cases}e^{\gamma x}, & x \leq 0 \\
e^{\gamma x+U(x)}\left(1-\psi^{-} \int_{0}^{x} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi\right), & x>0\end{cases} \\
& = \begin{cases}e^{\gamma x}, & x \leq 0 \\
\psi^{-}(\gamma) e^{\gamma x+U(x)} \int_{x}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi, & x>0\end{cases}
\end{aligned}
$$

and similarly

$$
M^{-}(\gamma)\binom{1}{\psi^{-}(\gamma)}(x)= \begin{cases}-e^{-\gamma x+U(x)} \int_{-\infty}^{x} U^{\prime}(\xi) e^{2 \gamma x-U(\xi)} d \xi, & x \leq 0 \\ \psi^{-}(\gamma) e^{-\gamma x}, & x>0\end{cases}
$$

Multiplying by $\frac{\psi^{+}(\gamma)}{\psi^{+}(\gamma)}$ and using that $\psi^{+}(\gamma) \psi^{-}(\gamma)=1$ yields the stated result.
Furthermore since $P_{\gamma}$ is of rank one and has range $\operatorname{span}\left\{f_{\gamma}\right\}$, there exists a $g \in L^{2}(\mu)$ such that $P_{\gamma} h=\langle h, g\rangle f_{\gamma}$. We have $P_{\gamma}^{\star} m=\left\langle m, f_{\gamma}\right\rangle g$, for $m \in L^{2}(\mu)$. Since $P_{\gamma}^{\star}$ is the spectral projection corresponding to eigenvalue $\bar{\gamma}$ of $L^{\star}$ (see [17, Theorem III.6.22]), it follows that $g$ is an eigenfunction of $L^{\star}$ with eigenvalue $\bar{\gamma}$. By Lemma 4.3, $\bar{g}$ is an eigenfunction of $L^{\star}$ with eigenvalue $\gamma$, and by Lemma 4.2(iv), $F \bar{g}$ is an eigenfunction of $L$ with eigenvalue $\gamma$, i.e., $F \bar{g}$ is parallel to $f_{\gamma}$. It follows that $F \bar{g}=\alpha f_{\gamma}$ for some $\alpha \in \mathbb{C}$, i.e., $g=\bar{\alpha} F \overline{f_{\gamma}}$. From $P_{\gamma}^{2}=P_{\gamma}$ it follows that $\left\langle f_{\gamma}, g\right\rangle=1$ which yields the correct value for $\alpha$.

Remark 4.7. Corollary 4.6 yields an expression for $Z^{\prime}(\gamma)$ for $\gamma \in \sigma(L)$, for $Z^{\prime}(\gamma) \neq 0$ :

$$
Z^{\prime}(\gamma)=\frac{\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle}{\left\langle h, F \bar{f}_{\gamma}\right\rangle} \frac{\left(1 \psi^{+}(\gamma)\right) K(\gamma) h}{\psi^{+}(\gamma)}
$$

It may be computed directly that $\left(1 \psi^{+}(\gamma)\right) K(\gamma) h=\left\langle h, F \bar{f}_{\gamma}\right\rangle$, so that

$$
Z^{\prime}(\gamma)=\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle / \psi^{+}(\gamma)=\psi^{-}(\gamma)\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle
$$

In fact this expression for $Z^{\prime}(\gamma)$ remains valid if $Z^{\prime}(\gamma)=0$, which may be verified by direct computation, showing that $\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle=0$, and making using of the following lemma.

Lemma 4.8. Suppose the assumptions of Theorem 4.4 hold. Then the functions $\psi^{ \pm}$(as defined in (4.1)) are holomorphic, and satisfy

$$
\begin{equation*}
\psi^{ \pm}(\gamma)=1 \mp 2 \gamma \int_{0}^{ \pm \infty} e^{\mp 2 \gamma \xi-U(\xi)} d \xi \tag{4.11}
\end{equation*}
$$

## Furthermore

$$
\begin{equation*}
\frac{d}{d \gamma} \psi^{ \pm}(\gamma)=\mp 2 \int_{0}^{ \pm \infty} e^{\mp 2 \gamma \xi-U(\xi)} d \xi+4 \gamma \int_{0}^{ \pm \infty} \xi e^{\mp 2 \gamma \xi-U(\xi)} d \xi, \quad \gamma \in \mathbb{C} \tag{4.12}
\end{equation*}
$$

Proof. This follows by partial integration.
Proposition 4.9 (Dominant eigenvalue). Suppose the assumptions of Theorem 4.4 hold. Then 0 is a simple eigenvalue of $L$ (i.e. its spectral projection has rank one). The spectral projection maps onto the space of constant functions.

Proof. From (4.11), $\psi^{ \pm}(0)=1$, so that $Z(0)=1-\psi^{+}(0) \psi^{-}(0)=0$. Furthermore, using (4.12),

$$
\frac{d}{d \gamma} Z(0)=-\psi^{+}(0) \frac{d}{d \gamma} \psi^{-}(0)-\psi^{-}(0) \frac{d}{d \gamma} \psi^{+}(0)=2 \int_{-\infty}^{\infty} e^{-U(\xi)} d \xi \neq 0
$$

The range of the spectral projection coincides with the span of the eigenfunction corresponding to the eigenvalue 0 , which is constant by Theorem 4.4.

Define the closed subspace $L_{0}^{2}(\mu):=\left\{f \in L^{2}(\mu): \mu(f)=0\right\}=\mathbb{1}^{\perp}$. From Proposition 4.9 and Corollary 4.6 it follows that $L_{0}^{2}(\mu)=\left(I-P_{0}\right) L^{2}(\mu)$, where $P_{0}=\langle\cdot, \mathbb{1}\rangle /\langle\mathbb{1}, \mathbb{1}\rangle \mathbb{1}$ denotes the spectral projection corresponding to the eigenvalue 0 . Since spectral projections associated with $L$ commute with the semigroup generated by $L$, it follows that the semigroup $(P(t))_{t \geq 0}$ leaves $L_{0}^{2}(\mu)$ invariant. The restriction of $(P(t))_{t \geq 0}$ to $L_{0}^{2}(\mu)$ has generator $L_{0}$, defined to be the restriction of $L$ to the domain $\mathcal{D}\left(L_{0}\right)=\mathcal{D}(L) \cap L_{0}^{2}(\mu)$. This establishes the following result.

Proposition 4.10 (Poisson equation). Suppose the assumptions of Theorem 4.4 are satisfied. Then $\sigma\left(L_{0}\right)=\sigma_{p}\left(L_{0}\right)=$ $\sigma(L) \backslash\{0\}$. In particular $L_{0}$ admits a bounded inverse.

### 4.1. Spectral gap

The following theorem establishes that the non-trivial spectrum of $L$ is strictly separated from the imaginary axis. The proof depends upon Lemmas A. 5 and A. 6 , which may be found in the appendix.

Theorem 4.11 (Spectral gap). Suppose the assumptions of Theorem 4.4 are satisfied. Define

$$
\kappa:=\inf \{-\operatorname{Re} \gamma: \gamma \in \sigma(L), \gamma \neq 0\} .
$$

Then $\kappa>0$.
Proof. Suppose on the contrary there exists a sequence $\left(\gamma_{n}\right) \in \sigma(L) \backslash\{0\}$ such that $\operatorname{Re} \gamma_{n} \rightarrow 0$. By Lemma A.5, it follows that $\operatorname{Im} \gamma_{n}$ is unbounded (for otherwise there would be an accumulation point $\mathrm{i} \beta$ for $\beta \in \mathbb{R}$ ). By extracting a subsequence, we may assume that $\left|\operatorname{Im} \gamma_{n}\right| \rightarrow \infty$. However, by Lemma A.6, applied to a small interval $\alpha \in[-\varepsilon, 0]$ this is impossible.

### 4.2. A spectral mapping theorem

In this section we will establish a spectral mapping theorem, which maps the spectrum of the infinitesemal generator to the spectrum of the semigroup. If the semigroup were eventually compact, this would be immediate (see, e.g., [10, Corollary IV.3.12]); however this is not the case as we have established in Proposition 3.12.

By the Gearhart-Herbst theorem [14, Theorem 4.5] spectrum of the semigroup may also originate from vertical lines in the complex plane (i.e. $\operatorname{Re} \gamma=$ constant) along which the resolvent of the generator is unbounded. We verify in Proposition 4.12 that this does not happen if we include Assumption A4, which details a sufficiently fast decay of the stationary distribution in its tails. Using Proposition 4.12 we can then prove the spectral mapping theorem, Theorem 4.13.

Proposition 4.12. Suppose Assumptions A1, A2, A3 and A4 are satisfied. Furthermore assume $\lambda_{\mathrm{refr}}(x)=0$ for all $x$. There is a family of constants $C(\alpha)$ such that for all $\alpha \in \mathbb{R}$ we have

$$
\limsup _{|\beta| \rightarrow \infty}\left\|(\alpha+\mathrm{i} \beta-L)^{-1}\right\| \leq C(\alpha)
$$

Proof. Suppose $\gamma \in \rho(L)$ and write $\gamma=\alpha+\mathrm{i} \beta$, where $\alpha, \beta \in \mathbb{R}$. In this proof we will often write $C(\alpha)$ for a positive constant which depends only on $\alpha$, whose value may be different at different locations in the proof. Let $h \in L^{2}(\mu)$ and let $f=(\gamma-L)^{-1} h$. The proof of Theorem 4.4 gives an expression for $f$ in terms of $h$. It only remains to provide adequate bounds for the expressions.

We have by the proof of Theorem 4.4, for $x \leq 0$, that

$$
f^{+}(x)=k^{+} e^{\gamma x}+e^{\gamma x} \int_{x}^{0} e^{-\gamma \xi} h^{+}(\xi) d \xi .
$$

We will first the following claim.
Claim 1: $\left\|x \mapsto f^{+}(x) \mathbb{1}_{\{x \leq 0\}}\right\|_{L^{2}(v)} \leq C(\alpha)\left(\left|k^{+}\right|+\left\|h^{+}\right\|_{L^{2}(v)}\right)$.
Proof of Claim 1. For the first term of $f^{+}, \int_{-\infty}^{0}\left|e^{2 \gamma x}\right| e^{-U(x)} d x=\int_{-\infty}^{0} e^{2 \alpha x} e^{-U(x)} d x \leq C(\alpha)$ by Assumption A2. For the second term of $f^{+}$, by Lemma A.7, taking $\psi(y)=h^{+}(-y) \mathbb{1}_{\{y \geq 0\}}$ and $\phi(y)=e^{-\gamma y} \mathbb{1}_{\{y \geq 0\}}$, we find that (substituting $x=-y$ and $\xi=-\eta$, and noting that Assumption A4 is invariant under this reparametrization)

$$
\left\|x \mapsto e^{\gamma x} \int_{x}^{0} e^{-\gamma \xi} h^{+}(\xi) d \xi \mathbb{1}_{\{x \leq 0\}}\right\|_{L^{2}(\nu)}=\left\|y \mapsto \int_{0}^{y} e^{-\gamma(y-\eta)} h^{+}(-\eta) d \eta \mathbb{1}_{\{y \geq 0\}}\right\|_{L^{2}(\tilde{\nu})}
$$

$$
\begin{aligned}
& =\|\phi \star \psi\|_{L^{2}(\tilde{v})} \\
& \leq\left(\int_{0}^{\infty} e^{-\alpha x-C / 2-(m / 2)|x|^{p}} d x\right)\left\|h^{+}\right\|_{L^{2}(v)}
\end{aligned}
$$

where $\tilde{\nu}(d y)=e^{-U(-y)} d y$. This establishes Claim 1.
By an analogous argument, we have that $\left\|x \mapsto f^{-}(x) \mathbb{1}_{\{x \geq 0\}}\right\|_{L^{2}(\nu)} \leq C(\alpha)\left(\left|k^{-}\right|+\left\|h^{-}\right\|_{L^{2}(\nu)}\right)$.
Next we will obtain a similar estimate for

$$
f^{+}(x)=e^{\gamma x+U(x)}\left(k^{+}-\int_{0}^{x} e^{-\gamma \xi-U(\xi)}\left[h^{+}(\xi)+U^{\prime}(\xi) e^{-\gamma \xi}\left(k^{-}+\int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta\right)\right] d \xi\right), \quad x \geq 0
$$

Claim 2: $\left\|x \mapsto f^{+}(x) \mathbb{1}_{\{x \geq 0\}}\right\|_{L^{2}(\nu)} \leq C(\alpha)\left(\left|k^{-}\right|+\left\|h^{+}\right\|_{L^{2}(\nu)}+\left\|h^{-}\right\|_{L^{2}(\nu)}\right)$.
Proof of Claim 2. Recall from the proof of Theorem 4.4 that $k^{+}$and $k^{-}$satisfy the relations (4.7). By substituting the expression for $k^{+}$in the above expression for $f^{+}$we obtain

$$
f^{+}(x)=e^{\gamma x+U(x)} \int_{x}^{\infty} e^{-\gamma \xi-U(\xi)}\left[\widetilde{h}^{+}(\xi)+k^{-} U^{\prime}(\xi) e^{-\gamma \xi}\right] d \xi, \quad x \geq 0
$$

Here $\widetilde{h}^{+}(\xi)=h^{+}(\xi)+U^{\prime}(\xi) e^{-\gamma \xi} \int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta, \xi \geq 0$. First we will establish Claim $2 a: x \mapsto \widetilde{h}^{+}(x) \mathbb{1}_{\{x \geq 0\}} \in L^{2}(\nu)$, with norm depending only on $\left\|h^{+}\right\|_{L^{2}(\nu)},\left\|h^{-}\right\|_{L^{2}(\nu)}$ and $\operatorname{Re} \gamma$. We only need to establish this for the second term of $\widetilde{h}^{+}$. We may estimate

$$
\left|e^{-\gamma \xi} \int_{0}^{\xi} e^{\gamma \eta} h^{-}(\eta) d \eta\right| \leq e^{-\alpha \xi} \int_{0}^{\xi} e^{\alpha \eta}\left|h^{-}(\eta)\right| d \eta
$$

An analogous argument as given in the proof of Claim 1 establishes that

$$
\phi: x \mapsto e^{-\alpha x} \int_{0}^{x} e^{\alpha \eta}\left|h^{-}(\eta)\right| d \eta \mathbb{1}_{\{x \geq 0\}} \in L^{2}(\nu)
$$

with $\|\phi\|_{L^{2}(\nu)}$ depending only on $\alpha$ and $\left\|h^{-}\right\|_{L^{2}(\nu)}$. Then $\phi$ has weak derivative $\partial_{x} \phi(x)=\left(\left|h^{-}(x)\right|-\alpha \phi(x)\right) \mathbb{1}_{\{x \geq 0\}}$ in $L^{2}(v)$ which may be bounded by a constant depending only on $\left\|h^{-}\right\|_{L^{2}(\nu)}$ and $\alpha$. Thus $\|\phi\|_{W^{1,2}(\nu)} \leq C(\alpha)\left\|h^{-}\right\|_{L^{2}(\nu)}$. By Lemma 3.7, $U^{\prime}(x) \phi(x) \in L^{2}(v)$ with norm depending only on $\alpha$ and $h$. This establishes Claim 2a. Next we establish Claim $2 b: \xi \mapsto U^{\prime}(\xi) e^{-\gamma \xi} \in L^{2}(\nu)$ with norm depending only on $\alpha$. Indeed

$$
\int_{0}^{\infty}\left(U^{\prime}(\xi)\right)^{2} e^{-2 \alpha \xi-U(\xi)} d \xi<\infty
$$

by Lemma 3.7. Combining Claims 2a and 2byields that for $\psi(\xi)=\widetilde{h}^{+}(\xi)+k^{-} U^{\prime}(\xi) e^{-\gamma \xi}$,

$$
\|\psi\|_{L^{2}(\nu)} \leq C(\alpha)\left(\left\|h^{+}\right\|_{L^{2}(\nu)}+\left\|h^{-}\right\|_{L^{2}(\nu)}+\left|k^{-}\right|\right)
$$

Now Claim 2 is proven by applying the second statement of Lemma A.7, applied to $\psi$ as defined above and $\phi(x)=e^{-\gamma x}$.
We may repeat the proofs of Claims 1 and 2 to establish analogous results for $f^{-}$. Thus we obtain that

$$
\|f\|_{L^{2}(\mu)} \leq C(\alpha)\left(\left|k^{+}\right|+\left|k^{-}\right|+\|h\|_{L^{2}(\mu)}\right)
$$

We will now express $k^{ \pm}$in terms of $h$ and $\gamma$. From (4.8) and (4.6), we obtain

$$
\binom{k^{+}}{k^{-}}=\frac{1}{Z(\gamma)} A(\gamma) K(\gamma) h, \quad \text { with } A(\gamma):=\left(\begin{array}{cc}
1 & \psi^{+}(\gamma) \\
\psi^{-}(\gamma) & 1
\end{array}\right)
$$

The first row of $K(\gamma)$ contains the expression

$$
\int_{0}^{\infty} e^{-\gamma \xi-U(\xi)} \widetilde{h}^{+}(\xi) d \xi
$$

where $\widetilde{h}^{+} \in L^{2}(\nu)$ is as in the proof of Claim 2 above. In the proof of Claim 2 it is established that $\widetilde{h}^{+} \in L^{2}(\nu)$, with norm depending only on $\|h\|_{L^{2}(\mu)}$ and $\alpha=\operatorname{Re} \gamma$. The same holds for the second row and we conclude that $\|K(\gamma) h\| \leq$ $C(\alpha)\|h\|_{L^{2}(\mu)}$ for some constant $C(\alpha)$ depending only on $\alpha$.

Recall $\psi^{+}(\gamma)=\int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi$. Since under our assumptions $x \mapsto U^{\prime}(x) e^{-U(x)-2 \alpha x} \mathbb{1}_{\{x \geq 0\}} \in L^{1}(\mathbb{R})$ for any $\alpha \in \mathbb{R}$, by Riemann-Lebesgue [10, Theorem C.8], we find for any $\alpha \in \mathbb{R}$ that $\lim _{|\beta| \rightarrow \infty} \psi^{+}(\alpha+\mathrm{i} \beta)=0$. By an analogous argument, $\lim _{|\beta| \rightarrow \infty} \psi^{-}(\alpha+\mathrm{i} \beta)=0$. Thus for any $\alpha$ the matrix $\beta \mapsto A(\alpha+\mathrm{i} \beta)$ is bounded. Furthermore

$$
\lim _{|\beta| \rightarrow \infty} Z(\alpha+\mathrm{i} \beta)=\lim _{|\beta| \rightarrow \infty}\left(1-\psi^{+}(\alpha+\mathrm{i} \beta) \psi^{-}(\alpha+\mathrm{i} \beta)\right)=1,
$$

so that the roots of $Z$ are contained in a bounded interval along the line $\operatorname{Re} \gamma=\alpha$.
Combining all estimates yields the stated result.
Theorem 4.13 (Spectrum of the zigzag semigroup). Suppose Assumptions A1, A2, A3 and A4 are satisfied. Furthermore assume $\lambda_{\text {refr }}(x)=0$ for all $x$. Then for all $t>0$,

$$
\begin{equation*}
\sigma(P(t)) \backslash\{0\}=\sigma_{p}(P(t)) \backslash\{0\}=\left\{e^{\gamma t}: \gamma \in \sigma(L)\right\} . \tag{4.13}
\end{equation*}
$$

Proof. Fix $t>0$. Spectral mapping of the point spectrum is well known [10, Theorem IV.3.7]: $\left\{e^{\gamma t}: \gamma \in \sigma_{p}(L)\right\}=$ $\sigma_{p}(P(t)) \backslash\{0\}$. By Theorem 4.4 we have $\sigma_{p}(L)=\sigma(L)$. Since $\sigma_{p}(P(t)) \backslash\{0\} \subset \sigma(P(t)) \backslash\{0\}$ it only remains to establish that $\sigma(P(t)) \backslash\{0\} \subset\left\{e^{\gamma t}: \gamma \in \sigma(L)\right\}$. Suppose $\eta \in \mathbb{C}, \eta \neq 0$, and consider the set

$$
\Gamma_{\eta}:=\left\{\gamma \in \mathbb{C}: e^{\gamma t}=\eta\right\}=\{\alpha+\beta i: \alpha=(\ln |\eta|) / t, \beta=(\arg \eta) / t+2 k \pi / t, k \in \mathbb{Z}\} .
$$

Suppose $\Gamma_{\eta} \cap \sigma(L)=\varnothing$. By Proposition 4.12 and the above characterization of $\Gamma_{\eta}$, it follows that $\gamma \mapsto\left\|(\gamma-L)^{-1}\right\|$ is bounded on $\Gamma_{\eta}$. Therefore, by [14, Theorem 4.5], $\eta \notin \sigma(P(t))$. In other words, if $\eta \in \sigma(P(t)) \backslash\{0\}$, then there is a $\gamma \in \Gamma_{\eta} \cap \sigma(L)$. This establishes that $\sigma(P(t)) \backslash\{0\} \subset\left\{e^{\gamma t}: \gamma \in \sigma(L)\right\}$.

Recall the spectral gap $\kappa$ of $L$,

$$
\kappa:=\inf \left\{-\operatorname{Re} \gamma: \gamma \in \sigma_{p}(L), \gamma \neq 0\right\},
$$

which by Theorem 4.11 satisfies $\kappa>0$. By Assumption A1, we have that $\mu(E)<\infty$. Define a Borel probability distribution $\pi$ on $E$ by

$$
\begin{equation*}
\pi(A)=\mu(A) / \mu(E), \text { for all Borel measurable } A \subset E \tag{4.14}
\end{equation*}
$$

Theorem 4.14. Suppose the assumptions of Theorem 4.13 are satisfied. There is a constant $M>0$ such that for any $f \in L^{2}(\mu)$,

$$
\|P(t) f-\pi(f)\|_{L^{2}(\mu)} \leq M e^{-\kappa t}\|f-\pi(f)\|_{L^{2}(\mu)} .
$$

For the proof we will consider the spectrum of the semigroup $(P(t))_{t \geq 0}$ restricted to $L_{0}^{2}(\mu)$ as defined just before Proposition 4.10, with infinitesemal generator $\widetilde{L}$.

Proof. Write $\widetilde{P}(t)$ for $\left.P(t)\right|_{L_{0}^{2}(\mu)}$. The argument of Theorem 4.13 may be repeated to establish that, for $t>0$,

$$
\sigma(\widetilde{P}(t)) \backslash\{0\}=\sigma_{p}(\widetilde{P}(t)) \backslash\{0\}=\left\{e^{\gamma t}: \gamma \in \sigma(\widetilde{L})\right\}=\left\{e^{\gamma t}: \gamma \in \sigma(L), \gamma \neq 0\right\} .
$$

Here we used that 0 is a simple eigenvalue (Proposition 4.9), which is removed from $\sigma(\widetilde{L})$, since $L_{0}^{2}(\mu)$ is by definition orthogonal to the associated eigenspace of constant functions. It then follows from [10, Proposition IV.2.2] that there exists an $M>0$ such that for the operator norm we have $\left\|P_{0}(t)\right\|_{L_{0}^{2}(\mu)} \leq M e^{-\kappa t}$. Applying $\widetilde{P}(t)$ to the function $\bar{f}=f-\pi(f)$ yields the stated result.

## 5. The symmetric case

Let us now consider Assumption A5 (Symmetry): $U(x)=U(-x)$ for all $x$. In this case the zigzag semigroup decouples into two separate semigroups:
(i) a 'climb-fall' semigroup acting on functions $f(x, \theta)$ which satisfy $f(x, \theta)=f(-x,-\theta)$, and
(ii) a second semigroup acting on functions $f(x, \theta)$ which satisfy $f(x, \theta)=-f(-x,-\theta)$.

The climb-fall semigroup (i) can be interpreted as the semigroup corresponding to the process $Y(t)$, where $Y(t)=$ $\Theta(t) X(t)$ with $(X(t), \Theta(t))$ denoting the Zig-Zag process, as described in the introduction. This process $Y(t)$ can (under Assumption A5) be seen to be a (piecewise deterministic) Markov process itself, with generator (5.1) below. The other semigroup (ii) is not a Markov semigroup.

We will now discuss this decomposition in more detail. Define an operator $T \in L\left(L^{2}(\mu)\right)$ by

$$
T f(x, \theta)=\frac{1}{\sqrt{2}}\left(f^{+}(x)+\theta f^{-}(-x)\right) .
$$

Then $T$ is unitary: Indeed, if $f, g \in L^{2}(\mu)$, then

$$
\begin{aligned}
\langle T f, T g\rangle= & \int_{\mathbb{R}}(T f)^{+}(x)(T g)^{+}(x) e^{-U(x)} d x+\int_{\mathbb{R}}(T f)^{-}(x)(T g)^{-}(x) e^{-U(x)} d x \\
= & \frac{1}{2} \int_{\mathbb{R}}\left(f^{+}(x)+f^{-}(-x)\right)\left(g^{+}(x)+g^{-}(-x)\right) e^{-U(x)} d x \\
& +\frac{1}{2} \int_{\mathbb{R}}\left(f^{+}(x)-f^{-}(-x)\right)\left(g^{+}(x)-g^{-}(-x)\right) e^{-U(x)} d x \\
= & \int_{\mathbb{R}}\left[f^{+}(x) g^{+}(x)+f^{-}(-x) g^{-}(-x)\right] e^{-U(x)} d x=\langle f, g\rangle,
\end{aligned}
$$

where the last equality holds due to Assumption A5. It may be checked that

$$
T^{\star} f(x, \theta)=T^{-1} f(x, \theta)=\frac{1}{\sqrt{2}}(f(\theta x,+1)+\theta f(\theta x,-1)) .
$$

Remark 5.1. Define subspaces of $L^{2}(\mu)$ as follows:

$$
\begin{aligned}
& \mathcal{H}_{S}:=\left\{f \in L^{2}(\mu): f^{+}(x)=f^{-}(-x) \text { for } x \in \mathbb{R}\right\}, \quad \mathcal{H}^{+}:=\left\{f \in L^{2}(\mu): f^{-}(x)=0 \text { for } x \in \mathbb{R}\right\} \\
& \mathcal{H}_{A}:=\left\{f \in L^{2}(\mu): f^{+}(x)=-f^{-}(-x) \text { for } x \in \mathbb{R}\right\}, \quad \mathcal{H}^{-}:=\left\{f \in L^{2}(\mu): f^{+}(x)=0 \text { for } x \in \mathbb{R}\right\} .
\end{aligned}
$$

We can write $f \in L^{2}(\mu)$ as

$$
\binom{f^{+}(x)}{f^{-}(x)}=\frac{1}{2}\binom{f^{+}(x)+f^{-}(-x)}{f^{-}(x)+f^{+}(-x)}+\frac{1}{2}\binom{f^{+}(x)-f^{-}(-x)}{f^{-}(x)-f^{+}(-x)},
$$

coressponding with the decomposition $L^{2}(\mu) \cong \mathcal{H}_{S} \oplus \mathcal{H}_{A}$. We thus have that $L^{2}(\mu) \cong \mathcal{H}^{+} \oplus \mathcal{H}^{-}, L^{2}(\mu) \cong \mathcal{H}_{S} \oplus \mathcal{H}_{A}$, $\mathcal{H}^{+} \perp \mathcal{H}^{-}$and moreover, using Assumption A5, $\mathcal{H}_{S} \perp \mathcal{H}_{A}$. Furthermore

$$
T \mathcal{H}_{S}=\mathcal{H}^{+}, \quad T \mathcal{H}_{A}=\mathcal{H}^{-}, \quad T^{\star} \mathcal{H}^{+}=\mathcal{H}_{S}, \quad T^{\star} \mathcal{H}^{-}=\mathcal{H}_{A} .
$$

In other words $T: \mathcal{H}_{S} \oplus \mathcal{H}_{A} \rightarrow \mathcal{H}^{+} \oplus \mathcal{H}^{-}$with the mapping $T$ respecting the direct sum.
The motivation for this transformation and these subspaces stems from the following proposition.
Proposition 5.2. Suppose Assumptions $A 1$ and $A 5$ are satisfied, and $\lambda_{\mathrm{refr}}(x)=\lambda_{\mathrm{refr}}(-x)$ for all $x$. The transformation of $(L, \mathcal{D}(L))$ by $T$, i.e., $(\tilde{L}, \mathcal{D}(L))$ with $\widetilde{L}=T L T^{\star}$, is given by $\mathcal{D}(\widetilde{L})=W^{1,2}(\mu)$ and

$$
(\widetilde{L} f)^{ \pm}(x)=\partial_{x} f^{ \pm}(x)+\lambda^{+}(x)\left( \pm f^{ \pm}(-x)-f^{ \pm}(x)\right), \quad f \in \mathcal{D}(\widetilde{L}) .
$$

This proposition states that the subspaces $\mathcal{H}^{ \pm}$are invariant for $\widetilde{L}$. In terms of the original generator $L$, this implies that the decomposition $L^{2}(\mu)=\mathcal{H}_{S} \oplus \mathcal{H}_{A}$ is invariant under $L$. That is, $L$ maps $f \in \mathcal{H}_{S} \cap \mathcal{D}(L)$ into $\mathcal{H}_{S}$ and $f \in \mathcal{H}_{A} \cap \mathcal{D}(L)$ into $\mathcal{H}_{A}$ respectively.

Proof. By Theorem 3.9 we have $\mathcal{D}(L)=W^{1,2}(\mu)$ which is invariant under $T$ and its inverse. Write $g(\cdot, \theta)=T^{\star} f(\cdot, \theta)=$ $\frac{1}{\sqrt{2}}\left(f^{+}(\theta \cdot)+\theta f^{-}(\theta \cdot)\right)$. Then

$$
\theta \partial_{x} g(\cdot, \theta)=\frac{1}{\sqrt{2}}\left[\left(\partial_{x} f^{+}\right)(\theta \cdot)+\theta\left(\partial_{x} f^{-}\right)(\theta \cdot)\right] .
$$

Using the above expression, and $\lambda^{+}(x)=\lambda^{-}(-x)$,

$$
\begin{aligned}
&\left(T L T^{\star} f\right)^{+}(x) \\
&= \frac{1}{\sqrt{2}}\left(L T^{\star} f^{+}(x)+L T^{\star} f^{-}(-x)\right) \\
&= \frac{1}{\sqrt{2}}\left(\partial_{x} g^{+}(x)-\partial_{x} g^{-}(-x)\right) \\
&+\frac{1}{\sqrt{2}}\left(\lambda^{+}(x)\left(g^{-}(x)-g^{+}(x)\right)+\lambda^{-}(-x)\left(g^{+}(-x)-g^{-}(-x)\right)\right) \\
&= \frac{1}{2}\left(\partial_{x} f^{+}(x)+\partial_{x} f^{-}(x)+\partial_{x} f^{+}(x)-\partial_{x} f^{-}(x)\right) \\
&+\frac{1}{2} \lambda^{+}(x)\left(f^{+}(-x)-f^{-}(-x)-f^{+}(x)-f^{-}(x)+f^{+}(-x)+f^{-}(-x)-f^{+}(x)+f^{-}(x)\right) \\
&= \partial_{x} f^{+}(x)+\lambda^{+}(x)\left(f^{+}(-x)-f^{+}(x)\right) .
\end{aligned}
$$

An analogous computation yields

$$
\left(T L T^{\star} f\right)^{-}(x)=\partial_{x} f^{-}(x)-\lambda^{+}(x)\left(f^{-}(-x)+f^{-}(x)\right)
$$

Recall that $v$ denotes the measure on $\mathbb{R}$ with density $e^{-U}$ with respect to Lebesgue measure. We see that under the transformation $T$, the generator decouples, and we may thus consider the two generators $L^{ \pm}$on the decoupled spaces given by

$$
\begin{equation*}
L^{+} f(x)=f^{\prime}(x)+\lambda^{+}(x)(f(-x)-f(x)), \quad f \in W^{1,2}(\nu), \tag{5.1}
\end{equation*}
$$

and

$$
L^{-} f(x)=f^{\prime}(x)+\lambda^{+}(x)(-f(-x)-f(x)), \quad f \in W^{1,2}(\nu)
$$

independently.

### 5.1. Spectral theory of the reduced semigroups

Here we consider the two semigroups with generator

$$
L^{ \pm} f(x)=f^{\prime}(x)+\lambda^{+}(x)( \pm f(-x)-f(x)), \quad f \in W^{1,2}(\nu)
$$

in $L^{2}(\nu)$. Because from an analytic point of view the two semigroups $L^{+}$and $L^{-}$are very similar, we carry out the analysis for the semigroups at once, at least as much as possible. We carefully keep track of the effect of the sign $\pm$ in the expressions that follow. We assume throughout Assumption A5, i.e., $U(x)=U(-x), x \in \mathbb{R}$.

Define functions $Z^{ \pm}: \mathbb{C} \rightarrow \mathbb{C}$ by

$$
\begin{equation*}
Z^{ \pm}(\gamma)=1 \mp \int_{0}^{\infty} U^{\prime}(\eta) e^{-2 \gamma \eta-U(\eta)} d \eta=1 \mp \psi^{+}(\gamma) \tag{5.2}
\end{equation*}
$$

and sets $\Sigma^{+} \subset \mathbb{C}$ and $\Sigma^{-} \subset \mathbb{C}$ by

$$
\begin{equation*}
\Sigma^{ \pm}=\left\{\gamma \in \mathbb{C}: Z^{ \pm}(\gamma)=0\right\} . \tag{5.3}
\end{equation*}
$$

Theorem 5.3. Suppose Assumptions A1, A2, A3 and A5 are satisfied. Furthermore assume $\lambda_{\mathrm{refr}}(x)=0$ for all $x$. Then $\sigma_{p}\left(L^{ \pm}\right)=\sigma\left(L^{ \pm}\right)=\Sigma^{ \pm}$. If $\gamma \in \rho\left(L^{ \pm}\right)$then, for $h \in L^{2}(\nu), f^{ \pm}=\left(\gamma-L^{ \pm}\right)^{-1} h \in W^{1,2}(\nu)$ admits the expression

$$
f^{ \pm}(x)= \begin{cases}e^{\gamma x}\left(\frac{m_{\nu}^{ \pm}(h)}{Z^{ \pm}(\gamma)}+\int_{x}^{0} e^{-\gamma \xi} h(\xi) d \xi\right), & x \leq 0,  \tag{5.4}\\ e^{\gamma x+U(x)} \int_{x}^{\infty} e^{-\gamma \xi-U(\xi)}\left[h(\xi) \pm U^{\prime}(\xi) e^{-\gamma \xi}\left(\frac{m_{\gamma}^{ \pm}(h)}{Z^{ \pm}(\gamma)}+\int_{-\xi}^{0} e^{-\gamma \eta} h(\eta) d \eta\right)\right] d \xi, & x>0,\end{cases}
$$

where $m_{\gamma}^{ \pm}(h)$ is given by

$$
\begin{equation*}
m_{\gamma}^{ \pm}(h)=\int_{0}^{\infty} e^{-\gamma \xi-U(\xi)}\left[h(\xi) \pm U^{\prime}(\xi) e^{-\gamma \xi} \int_{0}^{\xi} e^{\gamma \eta} h(-\eta) d \eta\right] d \xi \tag{5.5}
\end{equation*}
$$

For every $\gamma \in \sigma\left(L^{ \pm}\right)$the associated space of eigenfunctions is one-dimensional and spanned by $f_{\gamma}^{ \pm} \in W^{1,2}(\nu)$ given by

$$
f_{\gamma}^{ \pm}(x)= \begin{cases}e^{\gamma x}, & x \leq 0  \tag{5.6}\\ \pm e^{\gamma x+U(x)} \int_{x}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi, & x \geq 0\end{cases}
$$

In this proposition, and in the remainder of this manuscript, statements involving $\pm$ hold for the + and - cases, respectively. For example, $\sigma_{p}\left(L^{+}\right)=\sigma\left(L^{+}\right)=\Sigma^{+}$, and $f_{\gamma}^{+}$is an eigenfunction corresponding to the eigenvalue $\gamma \in \Sigma^{+}$ for the operator $L^{+}$.

Proof. The details of this proof (e.g., bounds on integrals) are analogous to those given in the proof of Theorem 4.4 and will be omitted. We only carry out the computations yielding the stated expressions. First suppose $\gamma \in \rho\left(L^{ \pm}\right)$. The resolvent equation is $\gamma f^{ \pm}-L^{ \pm} f^{ \pm}=h$ where $h \in L^{2}(\nu)$. For $x \leq 0$, since $\lambda^{+}(x)=0$, this becomes $\gamma f^{ \pm}(x)-\left(f^{ \pm}\right)^{\prime}(x)=$ $h(x)$, leading to the expression

$$
\begin{equation*}
f^{ \pm}(x)=e^{\gamma x}\left(k^{ \pm}(\gamma ; h)+\int_{x}^{0} e^{-\gamma \xi} h(\xi) d \xi\right), \quad x \leq 0, \tag{5.7}
\end{equation*}
$$

where $k^{ \pm}(\gamma ; h) \in \mathbb{C}$ are integration constants depending on $\gamma$ and $h$, which will be specified below. For $x \geq 0$ the resolvent equation may be written as

$$
\begin{equation*}
\gamma f^{ \pm}(x)-\left(f^{ \pm}\right)^{\prime}(x)+U^{\prime}(x) f^{ \pm}(x)=\widetilde{h}^{ \pm}(x), \quad x \geq 0 \tag{5.8}
\end{equation*}
$$

where, for $x \geq 0$,

$$
\widetilde{h}^{ \pm}(x):=h(x) \pm U^{\prime}(x) f^{ \pm}(-x)=h(x) \pm U^{\prime}(x) e^{-\gamma x}\left[k^{ \pm}(\gamma ; h)+\int_{-x}^{0} e^{-\gamma \xi} h(\xi) d \xi\right]
$$

This first order ordinary differential equation admits the solution

$$
\begin{equation*}
f^{ \pm}(x)=e^{\gamma x+U(x)}\left(k^{ \pm}(\gamma ; h)-\int_{0}^{x} e^{-\gamma \xi-U(\xi)} \widetilde{h}^{ \pm}(\xi) d \xi\right), \quad x \geq 0 \tag{5.9}
\end{equation*}
$$

where the integration constant is chosen so that $f^{ \pm}$is continuous in $x$. In case $\gamma \notin \Sigma^{ \pm}$, let $k^{ \pm}(\gamma ; h)=m_{\gamma}^{ \pm}(h) / Z^{ \pm}(\gamma)$ with $m_{\gamma}^{ \pm}(h)$ as in (5.5). Analogously to the proof of Theorem 4.4, it may then be verified that $f^{ \pm}(x)$ admits the expression

$$
\begin{equation*}
f^{ \pm}(x)=e^{\gamma x+U(x)} \int_{x}^{\infty} e^{-\gamma \xi-U(\xi)} \widetilde{h}^{ \pm}(\xi) d \xi, \quad x \geq 0 \tag{5.10}
\end{equation*}
$$

Indeed, expression (5.10) is equal to (5.9) if and only if

$$
\begin{equation*}
k^{ \pm}(\gamma ; h)=\int_{0}^{\infty} e^{-\gamma \xi-U(\xi)} \widetilde{h}^{ \pm}(\xi) d \xi= \pm k^{ \pm}(\gamma ; h) \int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi+m_{\gamma}^{ \pm}(h) \tag{5.11}
\end{equation*}
$$

Solving for $k^{ \pm}(\gamma ; h)$ and using (4.1) yields the stated expression for $k^{ \pm}(\gamma ; h)$. We have, using (5.7) and Lemma A. 7 for $x \leq 0$, and (5.10) and Lemma A. 4 for $x \geq 0$, that $f^{ \pm} \in L^{2}(\nu)$. This fully determines the solution $f^{ \pm}$of the resolvent equation. If $\gamma \in \Sigma^{ \pm}$then (5.11) implies that $m_{\gamma}^{ \pm}(h)$ which is defined by (5.5) is identically equal to 0 ; a contradiction.

Now suppose $\gamma \in \sigma\left(L^{ \pm}\right)$. Since $\sigma\left(L^{ \pm}\right) \subset \sigma(L)$, it follows by Corollary 4.1 that $\gamma$ is in the point spectrum of $L^{ \pm}$. A solution to the equation $\left(\gamma-L^{ \pm}\right) f_{\gamma}^{ \pm}$is given by

$$
f_{\gamma}^{ \pm}(x)= \begin{cases}e^{\gamma x}, & x \leq 0, \\ e^{\gamma x+U(x)}\left(1 \mp \int_{0}^{x} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi\right), & x \geq 0,\end{cases}
$$

insisting upon continuity in $x=0$. Since we require $f_{\gamma}^{ \pm} \in L^{2}(\nu)$, its growth as $x \rightarrow \infty$ should be controlled, requiring that

$$
1 \mp \int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)} d \xi=0
$$

with the respective limits being well defined by an application of Lemma A.4. But this is equivalent to $\gamma \in \Sigma^{ \pm}$, and yields the alternative expression (5.6).

Remark 5.4 (Adjoint generator). Introduce the operator $J: L^{2}(\nu) \rightarrow L^{2}(\nu)$ by $J g(x)=g(-x)$. It may be verified that $\left(L^{ \pm}\right)^{\star}=J L^{ \pm} J$, a very similar situation as in the previous case where $L^{\star}=F L F$. Indeed by partial integration (taking for simplicity $f, g \in C_{c}^{\infty}(E)$ to be real valued),

$$
\begin{aligned}
\int_{\mathbb{R}} f L^{ \pm} g d \nu & =\int_{\mathbb{R}} f(x)\left[g^{\prime}(x)+\lambda^{+}(x)( \pm g(-x)-g(x))\right] e^{-U(x)} d x \\
& =\int_{\mathbb{R}}\left[-f^{\prime}(x) g(x)+U^{\prime}(x) f(x) g(x) \pm \lambda^{+}(x) f(x) g(-x)-\lambda^{+}(x) f(x) g(x)\right] e^{-U(x)} d x \\
& =\int_{\mathbb{R}}\left[-f^{\prime}(x) g(x) \pm \lambda^{+}(-x) f(-x) g(x)-\lambda^{-}(x) f(x) g(x)\right] e^{-U(x)} d x \\
& =\int_{\mathbb{R}}\left[J \partial_{x} J f(x) \pm \lambda^{+}(-x) f(-x)-\lambda^{+}(-x) f(x)\right] g(x) e^{-U(x)} d x \\
& =\int_{\mathbb{R}}\left(J L^{ \pm} J f\right) g d \nu,
\end{aligned}
$$

where we used partial integration in the second inequality, performed a change of variables $x \rightarrow-x$ for the integral over $\lambda^{+}(x) f(x) g(-x)$, and used $U^{\prime}(x)-\lambda^{+}(x)=-\lambda^{-}(x)=-\lambda^{+}(-x)$ in the third and fourth inequality.

The observation that $\left(L^{ \pm}\right)^{\star}=J L^{ \pm} J$ implies that the results of Lemma 4.2 concerning the spectrum of $L^{ \pm}$apply.
Corollary 5.5 (Spectral projection). Under the assumptions of Theorem 5.3 the spectral projection corresponding to $\gamma \in \sigma\left(L^{ \pm}\right)$is given by

$$
P_{\gamma}^{ \pm} h=\frac{1}{2 \pi \mathrm{i}} \int_{\Gamma} k^{ \pm}(\zeta, h) f_{\zeta}^{ \pm} d \zeta
$$

where $\Gamma$ is a Jordan contour in $\mathbb{C}$ enclosing the eigenvalue $\gamma$ and no other eigenvalues, where $k^{ \pm}$and $f_{\zeta}^{ \pm}$are as defined in Theorem 5.3.

In particular if $\gamma$ is a simple root of $Z^{ \pm}$, then $P_{\gamma}^{ \pm}$has rank one and is given by

$$
P_{\gamma}^{ \pm} h=\frac{m_{\gamma}^{ \pm}(h)}{\frac{d}{d \gamma} Z^{ \pm}(\gamma)} f_{\gamma}^{ \pm}=\frac{\left\langle h, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}}{\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}} f_{\gamma}^{ \pm},
$$

with $m_{\gamma}^{ \pm}$given in Theorem 5.3.
Proof. As in the proof of Proposition 4.5 the non-vanishing terms after integrating the resolvent along a simple closed contour $\Gamma$ enclosing the eigenvalue $\gamma$ are only those depending on $k^{ \pm}(\gamma ; h)$ since these have $Z^{ \pm}(\gamma)$ in the denominator and are thus non-holomorphic on the interior of $\Gamma$. This yields the stated expression. The expressions for the case of a simple root of $Z^{ \pm}$follow analogously to the proof of Corollary 4.6

Remark 5.6. As was the case with Remark 4.7, it may be verified by direct computation that, if $\gamma \in \sigma_{p}\left(L^{ \pm}\right)$, we have

$$
m_{\gamma}^{ \pm}(h)=\left\langle h, J \bar{f}_{\gamma}^{ \pm}\right\rangle \quad \text { and } \quad \frac{d}{d \gamma} Z^{ \pm}(\gamma)=\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle
$$

also when $\frac{d}{d \gamma} Z^{ \pm}(\gamma)=0$.

## 6. Bounded perturbations

So far the characterization of the spectrum relies upon the assumption that $\lambda_{\text {refr }}(x)=0$ for all $x$, in which case relatively explicit solutions to resolvent and eigenvalue equations can be obtained. In order to extend our analysis to the case $\lambda_{\text {refr }} \neq 0$, we investigate the effect of small perturbations of the zigzag generator on its spectrum.

Proposition 6.1. Let $(L, \mathcal{D}(L))$ denote the generator of the zigzag semigroup in $L^{2}(\mu)$ and let $B$ be a bounded operator on $L^{2}(\mu)$. Suppose the assumptions of Theorem 4.4 are satisfied. Suppose $\gamma$ is a root of $Z$, so that $\gamma \in \sigma_{p}(L)$.
(i) For $\epsilon>0$ sufficiently small there is a set of (repeated) eigenvalues $\gamma_{j}(\epsilon) \in \sigma_{p}(L+\epsilon B)=\sigma(L+\epsilon B), j=1, \ldots, m$, with total algebraic multiplicity equal to the algebraic multiplicity $m$ of $\gamma$, such that $\gamma_{j}(\epsilon) \rightarrow \gamma$ as $\epsilon \downarrow 0$. Furthermore the eigenvalues $\gamma_{j}(\epsilon)$ admit the asymptotic expansions

$$
\gamma_{j}(\epsilon)=\gamma+\epsilon \mu_{j}^{(1)}+o(\epsilon), \quad j=1, \ldots, m
$$

where $\left(\mu_{j}^{(1)}\right)$ are the repeated eigenvalues of the operator $P_{\gamma} B P_{\gamma}$ considered in the m-dimensional space $P_{\gamma} L^{2}(\mu)$, where $P_{\gamma}$ denotes the spectral projection corresponding to $\gamma$.
(ii) If moreover $\gamma$ is a simple root of $Z$, so that it is a simple eigenvalue of $L$, then for $\epsilon>0$ sufficiently small the operator $L+\epsilon B$ has a simple eigenvalue $\gamma(\epsilon)$ such that $\gamma(\epsilon) \rightarrow \gamma$ as $\epsilon \downarrow 0$, and $\gamma(\epsilon)$ admits the asymptotic expansion

$$
\begin{equation*}
\gamma(\epsilon)=\gamma+\epsilon\left\langle B f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)} /\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)}+o(\epsilon) \tag{6.1}
\end{equation*}
$$

where $f_{\gamma}$ is the eigenfunction associated to $\gamma$ as given by Theorem 4.4.
Proof. Since $L$ has compact resolvent, also $L+\epsilon B$ has compact resolvent by [10, Proposition III.1.12], so that $\sigma(L+$ $\epsilon B)=\sigma_{p}(L+\epsilon B)$. Using that the dependence of $L+\epsilon B$ on $\epsilon$ is holomorphic, by [17, Theorem VII.1.3] the resolvent $(\zeta-L-\epsilon B)^{-1}$ is jointly holomorphic in $(\zeta, \epsilon)$ for sufficiently small $\epsilon$. Then [17, Theorem VII.1.7] establishes that spectral projections depend holomorphically on $\varepsilon$, in the sense of [17, Equation (VII.1.4)]. It then follows that $\gamma$ is a stable eigenvalue in the terminology of [17, Section VIII.1.4]): condition (i) of the definition of stable eigenvalue is satisfied by the mentioned joint analyticity of $(\zeta-L-\epsilon B)^{-1}$ around $\gamma$ and condition (ii) is satisfied by the mentioned analyticity of the spectral projections. We may therefore apply [17, Theorem VIII.2.6] to conclude (i). Finally (ii) follows from (i) and the expression for $P_{\gamma}$ obtained in Corollary 4.6.

Remark 6.2 (Perturbations in the symmetric case). If Assumption A5 is satisfied, and if additionally the bounded perturbation $B: L^{2}(\mu)$ respects the direct sum $B: \mathcal{H}_{S} \oplus \mathcal{H}_{A} \rightarrow \mathcal{H}_{S} \oplus \mathcal{H}_{A}$, then the perturbed operator $L+\epsilon B$ decomposes through the transformation $T$, i.e.,

$$
T(L+\epsilon B) T^{\star}=\left(\begin{array}{cc}
L^{+}+\epsilon B^{+} & 0 \\
0 & L^{-}+\epsilon B^{-}
\end{array}\right)
$$

with $L^{ \pm}$as before and $B^{ \pm}: L^{2}(\nu) \rightarrow L^{2}(\nu)$ bounded linear mappingss. If moreover the conditions of Theorem 4.4 are satisfied and $\gamma$ is a simple eigenvalue of $L^{ \pm}$, then we obtain analogously to Proposition 6.1 an asymptotic expansion for the eigenvalues $\gamma(\epsilon)$ of the generators $L^{ \pm}+\epsilon B^{ \pm}$given by

$$
\begin{equation*}
\gamma(\epsilon)=\gamma+\epsilon\left\langle B^{ \pm} f_{\gamma}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)} /\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}+o(\epsilon), \tag{6.2}
\end{equation*}
$$

where $\gamma$ is a simple eigenvalue of $L^{ \pm}$and $f_{\gamma}^{ \pm}$is the associated eigenfunction as given by Theorem 5.3 and we recall that $J f(x)=f(-x)$.

### 6.1. Example (additive perturbation of the switching intensity)

Consider the bounded perturbation operator $B f=F f-f$, so that $L+\epsilon B$ admits the expression

$$
(L+\epsilon B) f(x, \theta)=\theta \partial_{x} f(x, \theta)+\left(\max \left(\theta U^{\prime}(x), 0\right)+\epsilon\right)(f(x,-\theta)-f(x, \theta)) .
$$

We see that the perturbation $\epsilon B$ effectively adds a constant $\epsilon$ to the switching intensity. In case the potential function $U$ is such that the assumptions of Theorem 4.4 are satisfied and furthermore all eigenvalues of $L$ are simple, then by Proposition 6.1(ii) the perturbed eigenvalues $\gamma(\epsilon)$ allow the asymptotic expansion

$$
\begin{aligned}
\gamma(\epsilon) & =\gamma+\epsilon\left\langle B f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)} /\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)}+o(\epsilon) \\
& =\gamma+\epsilon\left(\left\langle F f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)} /\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)}-1\right)+o(\epsilon) \\
& =\gamma+\epsilon\left(\left\langle f_{\gamma}, \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)} /\left\langle f_{\gamma}, F \bar{f}_{\gamma}\right\rangle_{L^{2}(\mu)}-1\right)+o(\epsilon),
\end{aligned}
$$

where we subsequently used (6.1), the expression for $B$, and finally the fact that $\langle F f, F g\rangle_{L^{2}(\mu)}=\langle f, g\rangle_{L^{2}(\mu)}$ for any $f, g \in L^{2}(\mu)$.

Furthermore the operator $B$ satisfies the condition stated in Remark 6.2 of preserving the decomposition $\mathcal{H}_{S} \oplus \mathcal{H}_{A}$. We compute

$$
T B T^{\star} f(x,+1)=f(-x,+1)-f(x,+1), \quad \text { and } \quad T B T^{\star} f(x,-1)=-f(-x,-1)-f(x,-1),
$$

so that in the notation of Remark 6.2, we have

$$
B^{ \pm} f= \pm J f-f .
$$

This means that if the additional assumption of symmetry (Assumption A5, $U(x)=U(-x)$ ) is satisfied, then we may determine from (6.2) the perturbed eigenvalues as

$$
\begin{aligned}
\gamma(\epsilon) & =\gamma+\epsilon\left\langle B^{ \pm} f_{\gamma}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)} /\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}+o(\epsilon) \\
& =\gamma+\epsilon\left(\left\langle \pm J f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)} /\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}-1\right)+o(\epsilon) \\
& =\gamma+\epsilon\left(\left\langle \pm f_{\gamma}^{ \pm}, \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)} /\left\langle f_{\gamma}^{ \pm}, J \bar{f}_{\gamma}^{ \pm}\right\rangle_{L^{2}(\nu)}-1\right)+o(\epsilon),
\end{aligned}
$$

where we should read ' + ' or ' - ' in the above expression depending on whether $\gamma \in \sigma\left(L^{+}\right)=\Sigma^{+}$or $\gamma \in \sigma\left(L^{-}\right)=\Sigma^{-}$.
Unfortunately the expression does not seem to simplify further and in order to compute the above expansion in particular situations we will have to resort to numerical computation of the integrals involved.

## 7. Examples

### 7.1. A family of medium-heavy to light-tailed symmetric distributions

Consider

$$
\begin{equation*}
U(x)=\frac{1}{\beta}\left[\left(1+x^{2}\right)^{\beta / 2}-1\right] \quad \text { where } \beta>1 . \tag{7.1}
\end{equation*}
$$

We compute

$$
\begin{aligned}
& U^{\prime}(x)=x\left(1+x^{2}\right)^{\beta / 2-1} \\
& U^{\prime \prime}(x)=\left(1+x^{2}\right)^{\beta / 2-2}\left(1+(\beta-1) x^{2}\right) .
\end{aligned}
$$

The case $\beta=2$ corresponds to the standard normal distribution. Note that for $\beta \leq 1$ already the most basic assumptions of our theory (e.g., Assumption A2) are not satisfied.

Lemma 7.1. Suppose $U$ is given by (7.1). Then Assumptions A1, A2, A3, $A 4$ and $A 5$ are satisfied.

Proof. Assumptions A3 and A5 are immediate. From the expression for $U^{\prime \prime}(x)$ it follows that $U^{\prime \prime}(x) \geq 0$ for all $x \in \mathbb{R}$ and $\beta>1$. We have

$$
\frac{U^{\prime \prime}(x)}{\left(U^{\prime}(x)\right)^{2}}=\left(1+x^{2}\right)^{-\beta / 2}\left(1+(\beta-1) x^{2}\right) x^{-2} \sim(\beta-1)|x|^{-\beta},
$$

establishing that Assumption A1 is satisfied. Furthermore $\left|U^{\prime}(x)\right| \sim|x|^{\beta-1} \rightarrow \infty$ as $|x| \rightarrow \infty$, establishing Assumption A2.

Asymptotically, $U^{\prime \prime}(x) \sim(\beta-1)|x|^{\beta-2}$ as $|x| \rightarrow \infty$, which establishes Assumption A4 for $\beta \geq 2$, using Lemma 2.3.
Finally we establish Assumption A4 for $1<\beta<2$. Let $M>1$. By symmetry it suffices to consider $y \geq x \geq M$ in Assumption A4. For all $x \geq M,\left(1+x^{2}\right) \leq\left(1+\frac{1}{M^{2}}\right) x^{2}$. It follows that

$$
\begin{equation*}
U^{\prime}(x)=x\left(1+x^{2}\right)^{\beta / 2-1} \geq\left(1+\frac{1}{M^{2}}\right)^{\beta / 2-1} x^{\beta-1}, \quad x \geq M . \tag{7.2}
\end{equation*}
$$

Define $h(\xi)=U(x+\xi)-U(x)-m \xi^{\beta}$ for $\xi \geq 0$ and fixed $x \geq M$. Then, using (7.2)

$$
\begin{aligned}
h^{\prime}(\xi) & =U^{\prime}(x+\xi)-m \beta \xi^{\beta-1} \geq\left(1+\frac{1}{M^{2}}\right)^{\beta / 2-1}(x+\xi)^{\beta-1}-m \beta \xi^{\beta-1} \\
& \geq\left[\left(1+\frac{1}{M^{2}}\right)^{\beta / 2-1}-m \beta\right] \xi^{\beta-1}=0,
\end{aligned}
$$

by taking $m=\frac{1}{\beta}\left(1+\frac{1}{M^{2}}\right)^{\beta / 2-1}$. We see that $h(\xi)$ is strictly non-decreasing as a function of $\xi$ for any $x \geq M$. Thus $h(y-x) \geq h(0)=0$ for all $y$, establishing that

$$
U(y) \geq U(x)+m|y-x|^{\beta}, \quad y \geq x \geq M
$$

as required.

### 7.2. Gaussian distribution

Suppose $U(x)=x^{2} /\left(2 \sigma^{2}\right)$. Without loss of generality we assume that $\sigma=1$ and consider $U(x)=x^{2} / 2$. This is an instance of the family discussed above, with $\beta=2$. In particular Assumptions A1 until A5 are satisfied. Let $\operatorname{erfc}(z)=$ $1-\operatorname{erf}(z)$ denote the complimentary error function, with the error function $\operatorname{erf}(z)$ given by

$$
\operatorname{erf}(z)=\frac{2}{\sqrt{\pi}} \int_{0}^{z} e^{-t^{2}} d t, \quad z \in \mathbb{C}
$$

Proposition 7.2. Suppose $U(x)=x^{2} / 2$ and $\lambda_{\text {refr }}(x)=0$ for all $x$. Then the spectrum of the zigzag semigroup is given by

$$
\sigma(L)=\sigma_{p}(L)=\Sigma^{+} \cup \Sigma^{-},
$$

where

$$
\Sigma^{+}=\{0\} \cup\{\gamma \in \mathbb{C}: \operatorname{erfc}(\sqrt{2} \gamma)=0\} \quad \text { and } \quad \Sigma^{-}=\left\{\gamma \in \mathbb{C}: \sqrt{\pi} \gamma e^{2 \gamma^{2}} \operatorname{erfc}(\sqrt{2} \gamma)=\sqrt{2}\right\} .
$$

All eigenvalues $\gamma \in \sigma(L)$ are algebraically simple.
Proof. As established in Theorem 5.3, the spectrum consists of eigenvalues only, and is given by $\sigma(L)=\Sigma^{+} \cup \Sigma^{-}$, where

$$
\Sigma^{ \pm}=\sigma\left(L^{ \pm}\right)=\left\{\gamma \in \mathbb{C}: \psi^{+}(\gamma)= \pm 1\right\} .
$$

By partial integration (Lemma 4.8)

$$
\begin{equation*}
\psi^{+}(\gamma)=1-\sqrt{2 \pi} \gamma e^{2 \gamma^{2}} \operatorname{erfc}(\sqrt{2} \gamma), \quad \gamma \in \mathbb{C} \tag{7.3}
\end{equation*}
$$

resulting in the characterization of $\sigma^{ \pm}$. From (4.11) and (4.12) we have

$$
\begin{equation*}
\psi^{+}(\gamma)=1+\gamma \frac{d}{d \gamma} \psi^{+}(\gamma)-4 \gamma^{2} \psi^{+}(\gamma) \tag{7.4}
\end{equation*}
$$

Note that $\psi^{-}(\gamma)=\psi^{+}(\gamma)$ by symmetry of $U$. Since $Z(\gamma)=1-\left(\psi^{+}(\gamma)\right)^{2}$, it follows that $\frac{d}{d \gamma} Z(\gamma)=$ $-2 \psi^{+}(\gamma) \frac{d}{d \gamma} \psi^{+}(\gamma)$. Thus if $\gamma$ is an eigenvalue with algebraic multiplicity larger than one, then $\psi^{+}(\gamma)= \pm 1$ and $\frac{d \psi^{+}}{d \gamma}=0$. Assume $\frac{d}{d \gamma} \psi^{+}(\gamma)=0$. Using (7.4), in case $\psi^{+}(\gamma)=1$, then necessarily $\gamma=0$; however it was established in the proof of Proposition 4.9 that $\frac{d}{d \gamma} Z(0) \neq 0$. In case $\psi^{+}(\gamma)=-1$, using again (7.4), we find that $\gamma= \pm \frac{1}{2} \sqrt{2} \mathrm{i}$, which is however not a root of $\psi^{+}$as defined in (7.3).

Asymptotic expansions for the elements of $\Sigma^{+}$can be found in [11]. However for a complete computation of the spectrum we have to rely on numerical computation.

### 7.3. Numerical results

By Theorems 4.4 and 5.3, the spectrum of $L$ and $L^{ \pm}$may be computed as the roots of $Z$ (see (4.2)) and $Z^{ \pm}$(see (5.2)), respectively. For ease of notation we only consider roots of $Z$ but the exposition applies equally to roots of $Z^{ \pm}$.

A suitable numerical routine for determining roots of $Z$ is provided by [8], to which we refer for details. For our purposes it is important to remark that the method of [8] relies on the ability to integrate $Z^{\prime}(\zeta) / Z(\zeta)$ efficiently along line segments in the complex plane. We have

$$
\frac{Z^{\prime}(\zeta)}{Z(\zeta)}=-\frac{\psi^{-}(\zeta) \frac{d}{d \zeta} \psi^{+}(\zeta)+\psi^{+}(\zeta) \frac{d}{d \zeta} \psi^{-}(\zeta)}{1-\psi^{+}(\zeta) \psi^{-}(\zeta)} .
$$

In the case of a Gaussian target distribution (Section 7.2), these functions may be expressed in terms of the error function. However for the general family considered in Section 7.1 we have to rely on numerical integration to determine the value of $Z^{\prime}(\zeta) / Z(\zeta)$.

Computer code (in Julia) for visualization of the spectrum and perturbations may be found at [2]. Graphical depictions of the spectrum, including the effect of perturbations by a positive switching rate, for the Gaussian case ( $\beta=2$ ) and two other cases are provided in Figure 1. In the Gaussian case the rightmost (pairs of) eigenvalues are given (up to 6 significant digits) by $0,-0.425665 \pm 1.02295 i,-0.957995 \pm 1.40818 \mathrm{i},-1.26616 \pm 1.66757 \mathrm{i},-1.53940 \pm-1.90293 \mathrm{i}$ and we find that the spectral gap is given by 0.425665 . An interesting observation is that a positive refreshment rate increases the spectral gap (since the spectrum moves towards the left). This phenomenon is also observed in [20]. Furthermore it seems that making the tail more heavy (i.e. decreasing $\beta$ ) results in an increase of the spectral gap.

## 8. Discussion

In this section we will explore the relation between the results obtained in this paper, and other results concerning exponential convergence to equilibrium. Assume $\mu(E)<\infty$ and recall the notation $\pi$ for the normalization of $\mu$, defined by (4.14).

### 8.1. Connection between spectral gap and geometric ergodicity

Suppose we have a spectral gap for the zigzag semigroup; e.g., assume the conditions of Theorem 4.14 are satisfied. We consider a time discretization of the zigzag semigroup $(P(t))_{t \geq 0}$ by defining, for a fixed $t_{0}>0$, the Markov kernel

$$
Q(\eta, A):=\int P\left(t_{0}\right) \mathbb{1}_{A}(x, \theta) \eta(d x, d \theta)
$$

where $\eta$ is any probability measure on $E=\mathbb{R} \times\{-1,+1\}$.
For any signed measure $\eta \ll \pi$ we define $\|\eta\|_{L^{2}(\pi)}^{2}=\int\left(\frac{d \eta}{d \pi}\right)^{2} d \pi$. We say that a Markov chain with transition kernel $Q$ on a general state space $(X, \mathcal{X})$ is $L^{2}(\pi)$-geometrically ergodic if there is a function $c(\eta)<\infty$ and a constant $\rho<1$ such that

$$
\left\|Q^{n}(\eta, \cdot)-\pi\right\|_{L^{2}(\pi)} \leq c(\eta) \rho^{n}
$$

for all Borel probability distributions $\eta$ on $E$ satisfying $\|\eta\|_{L^{2}(\pi)}<\infty$.


Fig. 1. The spectra of $L^{+}$and $L^{-}$for the distributions described in Section 7.1, along with the directions in which the spectrum is perturbed under a small refreshment rate $\lambda_{\text {refr }}=\epsilon>0$.

By considering the adjoint semigroup of the zigzag semigroup we then have the following result.
Corollary 8.1. Under the assumptions of Theorem 4.14, the Markov kernel $Q$ is $L^{2}(\pi)$-geometrically ergodic.
Proof. We have that $Q^{n}(\eta, A)=\int_{A} P^{\star}\left(n t_{0}\right) \frac{d \eta}{d \pi} d \pi$ where $P^{\star}(t)$ is the adjoint semigroup of $P(t)$. Since the result of Theorem 4.14 for $(P(t))_{t \geq 0}$ carries over to $\left(P^{\star}(t)\right)_{t \geq 0}$ (using Theorem 3.6), the stated claim follows.

An alternative notion of geometric ergodicity is the following. We say that a chain is $\pi$-almost everywhere geometrically ergodic if there is a function $W: E \rightarrow \infty$ and $\rho<1$ such that for $\pi$-almost all $z$,

$$
\left\|Q^{n}(z, \cdot)-\pi\right\|_{\mathrm{TV}} \leq W(z) \rho^{n} .
$$

Importantly the latter notion is phrased in terms of $\pi$-almost all initial positions. However, as a result of [23, Theorem 1] we have indeed that $Q$ is also $\pi$-almost everywhere geometrically ergodic.

### 8.2. Comparison to hypocoercivity analysis

In [1] conditions for hypocoercivity of piecewise deterministic process such as the zigzag process are established. We will rephrase the result of [1] to fit in the present context.

Theorem 8.2 ([1, Theorem 1]). Suppose the following assumptions hold:
(i) $U \in C^{3}(\mathbb{R}), \lambda_{\text {refr }}$ is continuous and $e^{-U} \in L^{1}(\mathbb{R})$.
(ii) There is a $p \geq 0$ such that $\sup _{x \in \mathbb{R}}\left|U^{(3)}(x)\right| /\left(1+|x|^{p}\right)<\infty$.
(iii) There is a constant $c_{1} \geq 0$ such that $U^{\prime \prime}(x) \geq-c_{1}$ for all $x \in \mathbb{R}$.
(iv) $\liminf \mathrm{f}_{|x| \rightarrow \infty}\left\{\left|U^{\prime}(x)\right|^{2} / 2-U^{\prime \prime}(x)\right\}>0$.
(v) The refreshment rate $\lambda_{\text {refr }}: \mathbb{R} \rightarrow(0, \infty)$ is bounded from below and from above as follows. There exists $\underline{\lambda}>0$ and $c_{\lambda} \geq 0$ such that for all $x \in \mathbb{R}$,

$$
0<\underline{\lambda} \leq \lambda_{\mathrm{refr}}(x) \leq \underline{\lambda}\left(1+c_{\lambda}\left|U^{\prime}(x)\right|\right) .
$$

Then there exist constants $C>0$ and $\alpha>0$ such that, for any $f \in L_{0}^{2}(\mu)$ and $t \geq 0$,

$$
\|P(t) f\|_{L_{0}^{2}(\mu)} \leq C e^{-\alpha t}\|f\|_{L_{0}^{2}(\mu)} .
$$

Proof. We may write the generator of the zigzag semigroup in the notation of [1, Equation (1)],

$$
L f(x, \theta)=\theta \partial_{x} f(x, \theta)+\lambda_{1}(x, \theta)\left(\mathcal{B}_{1}-I\right) f(x, \theta)+m_{2}^{1 / 2} \lambda_{\mathrm{ref}}(x) \mathcal{R}_{v} f(x, \theta)
$$

where

$$
\begin{aligned}
& \lambda_{1} f(x, \theta)=\left(\theta F_{1}(x)\right)_{+}, \quad \text { for } F_{1}(x)=U^{\prime}(x), \\
& \mathcal{B}_{1} f(x, \theta)=f(x,-\theta), m_{2}=1, \quad \lambda_{\text {ref }}(x)=2 \lambda_{\text {refr }}(x), \quad \mathcal{R}_{v} f(x, \theta)=\frac{1}{2}(f(x,-\theta)-f(x, \theta)) .
\end{aligned}
$$

(The factor 2 had to be redistributed between the refreshment rate and the refreshment kernel to be consistent with the definition of $\mathcal{R}_{v}$ in [1].)

We verify that the stated assumptions imply the assumptions A1 and H1-H6 of [1]. Some notations below are from [1], and may override notation from this paper; in particular $v$ for the uniform marginal stationary distribution on the velocities $\{-1,+1\}$.

By Lemma 3.2 and Theorem 3.6, Assumption A1 of [1] is satisfied.
By (ii), $U \in C_{\text {poly }}^{(3)}(\mathbb{R})$. Together with assumptions (iii) and (iv) this implies H 1 .
We have $F_{1}(x)=U^{\prime}(x)$, which by $(\mathrm{i})$ is in $C^{2}(\mathbb{R})$. It is then clear that H 2 is satisfied.
H3 is satisfied by taking $\varphi(s)=\max (0, s)$.
H 4 is satisfied by taking $V=\{-1,+1\}$ and $\nu=\frac{1}{2}\left(\delta_{-1}+\delta_{+1}\right)$.
H5 is satisfied for $\mathcal{R}_{v}:=\Pi_{v}-I$ where $\Pi_{v} f=v(f)$. In particular if $g \in L_{0}^{2}(\nu)$ we have, writing $g=\left(g_{1}, g_{2}\right)=$ $\left(g_{1},-g_{1}\right)$,

$$
\left\langle-\mathcal{R}_{v} g, g\right\rangle_{L^{2}(\nu)}=-\frac{1}{2} g_{1}^{2}-\frac{1}{2} g_{2}^{2}+\frac{1}{4}\left(g_{1}+g_{2}\right)^{2}=-\frac{1}{4}\left(g_{1}-g_{2}\right)^{2}=\|g\|_{L^{2}(v)} .
$$

Finally H6 is identical to (v).
Bounds for the constants $C$ and $\alpha$ can, in principle, be obtained in explicit form although this requires a tedious effort. The result above may be compared to the same result of Theorem 4.14, which depends on Assumptions A1, A2, A3 and A4. The above conditions are different from our conditions in several respects. Condition (iv) in the above theorem is stronger than our Assumption A1, for example. Furthermore a strictly positive refreshment rate is required. On the other hand we require that $\left|U^{\prime}(x)\right| \rightarrow \infty$ (Assumption A2), a zero refreshment rate and unimodal target distribution to obtain the explicit characterization of the eigenvalues.

### 8.3. Comparison to exponential ergodicity

In [5] the following result on exponential ergodicity of the zigzag process is established. A function $V: E \rightarrow \mathbb{R}$ is called norm-like if $V(x, \theta)>0$ for all $(x, \theta) \in E$ and $\lim _{|x| \rightarrow \infty} V(x, \theta)=\infty$ for $\theta= \pm 1$.

Theorem 8.3 ([5, Theorem 5]). Suppose $U$ is continuously differentiable, $\lambda_{\text {refr }}$ is continuous, and for $\lambda^{ \pm}(x)=\left( \pm U^{\prime}(x) \vee\right.$ $0)+\lambda_{\text {refr }}(x)$, we have that

$$
\liminf _{x \rightarrow \infty} \lambda^{+}(x)>\limsup _{x \rightarrow \infty} \lambda^{-}(x) \text { and } \liminf _{x \rightarrow-\infty} \lambda^{-}(x)>\limsup _{x \rightarrow-\infty} \lambda^{+}(x) .
$$

Then there are constants $c>0, \alpha>0$ and a continuous norm-like function $V \in L^{2}(\mu), V>0$ on $E$, such that for all $(x, \theta) \in E$,

$$
|P(t) f(x, \theta)-\pi(f)| \leq c(1+V(x, \theta)) e^{-\alpha t}, \quad t \geq 0
$$

for all measurable $f: E \rightarrow \mathbb{R}$ satisfying $|f(x, \theta)| \leq 1+V(x, \theta)$ on $E$.
Note that here the family of operators $P(t)$ represents the extension of the Markov semigroup to a suitable space of measurable functions. In the case with constant refreshment rate the condition of Theorem 8.3 reduces to the requirement that $\liminf _{x \rightarrow \infty} U^{\prime}(x)>0$ and $\lim \sup _{x \rightarrow-\infty} U^{\prime}(x)<0$. We see that the above result implies, under only mild conditions, exponentially fast convergence in $L^{2}(\pi)$ for functions $f$ satisfying $|f(x, \theta)| \leq 1+V(x, \theta)$. However, this class of functions does not contain the full space $L^{2}(\pi)$. Indeed, by the proof of [5, Theorem 5], the function $V$ is growing at an exponential rate as $|x| \rightarrow \infty$, putting a restriction on the growth of the functions $f(x, \theta)$.

## Appendix: Technical lemmas

## A.1. Weak derivatives

In this section we recall some basic results on weak derivatives as used in Section 3. A function $f: \mathbb{R} \rightarrow \mathbb{R}$ is said to have a weak derivative $f^{\prime}$ if, for all $\varphi \in C_{c}^{\infty}(\mathbb{R})$, we have

$$
\int_{\mathbb{R}} f \varphi^{\prime} d x=-\int_{\mathbb{R}} f^{\prime} \varphi d x
$$

If furthermore, for $p \geq 1,\left.f\right|_{U}$ and $\left.f^{\prime}\right|_{U} \in L^{p}(U)$ for any open $U \subset \mathbb{R}$ with compact closure, then we write $f \in W_{\text {loc }}^{1, p}(\mathbb{R})$.
Lemma A.1. Suppose $f \in W_{\mathrm{loc}}^{1, p}(\mathbb{R})$ and $g \in C^{1}(\mathbb{R})$. Then $f g \in W_{\mathrm{loc}}^{1, p}(\mathbb{R})$, and $(f g)^{\prime}=f^{\prime} g+f g^{\prime}$.
Proof. Let $\varphi \in C_{c}^{\infty}(\mathbb{R})$ and write $K=\operatorname{supp} \varphi$. By standard approximation arguments we may find $\tilde{g} \in C_{c}^{\infty}(R)$ such that $\sup _{x \in K}|\tilde{g}(x)-g(x)|+\left|\tilde{g}^{\prime}(x)-g^{\prime}(x)\right|<\varepsilon$. We then have by the definition of the weak derivative that

$$
\int_{\mathbb{R}} f \tilde{g} \varphi^{\prime} d x=\int_{\mathbb{R}} f(\tilde{g} \varphi)^{\prime} d x-\int_{\mathbb{R}} f \tilde{g}^{\prime} \varphi d x=-\int_{\mathbb{R}} f^{\prime} \tilde{g} \varphi d x-\int_{\mathbb{R}} f \tilde{g}^{\prime} \varphi d x .
$$

Since $\varepsilon>0$ is arbitrary, and $\varphi$ has compact support, this yields

$$
\int_{\mathbb{R}} f g \varphi^{\prime} d x=-\int_{\mathbb{R}}\left(f^{\prime} g+f g^{\prime}\right) \varphi d x
$$

so that $f g$ has weak derivative $f^{\prime} g+f g^{\prime}$. Now since $g$ and $g^{\prime}$ are bounded on any bounded set, and $f \in W_{\text {loc }}^{1, p}(\mathbb{R})$, it follows that $f g \in W_{\text {loc }}^{1, p}(\mathbb{R})$.

## A.2. Assumption A 4 revisited

The following lemma establishes that the constant $M$ in Assumption A4 can, without loss of generality, be assumed to be equal to zero.

Lemma A.2. Suppose Assumption A4 is satisfied for certain values $C \leq 0, M>0, m>0$ and $p>1$. Then it is also satisfied for $M=0$. That is, for $\widetilde{m}:=2^{1-p_{m}}$ and some $\widetilde{C} \leq C$ we have that

$$
\begin{equation*}
U(y) \geq U(x)+\widetilde{C}+\widetilde{m}|x-y|^{p} \quad \text { for all } x, y \text { for which } y \geq x \geq 0 \text { or } y \leq x \leq 0 . \tag{A.1}
\end{equation*}
$$

Proof. Suppose Assumption A4 holds. We prove (A.1) for $y \geq x \geq 0$; the proof for $y \leq x \leq 0$ is analogous. First suppose $0 \leq x \leq y \leq M$. Write $C_{1}:=\max _{0 \leq \xi \leq M} U(\xi)$ and $C_{2}:=\min _{0 \leq \xi \leq M} U(\xi)$. Define $\widetilde{C}_{1}=C_{2}-C_{1}-m M^{p} \leq 0$. Then

$$
U(y) \geq U(x)+\widetilde{C}_{1}+m|x-y|^{p}
$$

Next consider the situation $0 \leq x \leq M \leq y$. We have by Assumption A4 that

$$
U(y) \geq U(M)+C+m|M-y|^{p} .
$$

Next

$$
U(M) \geq U(x)+\widetilde{C}_{1}+m|M-x|^{p}
$$

by the first part of this proof. Finally we have by Jensen's inequality that

$$
|y-M|^{p}+|M-x|^{p} \geq 2^{1-p}|y-x|^{p} .
$$

It follows that

$$
U(y) \geq U(x)+\widetilde{C}_{1}+m|M-x|^{p}+C+m|M-y|^{p} \geq U(x)+\widetilde{C}_{1}+C+2^{1-p} m|y-x|^{p} .
$$

By assumption the required inequality also holds for $M \leq x \leq y$. The statement of the lemma follows by taking $\widetilde{C}:=$ $\min \left(C, \widetilde{C}_{1}, C+\widetilde{C}_{1}\right)=C+\widetilde{C}_{1}$.

## A.3. The operator $C$ and its core

As in Section 3.2, for $f \in W_{\text {loc }}^{1,2}(\mathbb{R})$ let $C f=f^{\prime}-U f^{\prime}$, and (with slight abuse of notation) restrict $C$ to the domain $\mathcal{D}(C)=\left\{f \in L^{2}(\nu) \cap W_{\text {loc }}^{1,2}(\mathbb{R}): C f \in L^{2}(\nu)\right\}$.

Lemma A.3. $(C, \mathcal{D}(C))$ is closed and $C_{c}^{\infty}(\mathbb{R})$ is a core for $(C, \mathcal{D}(C))$.
Proof. First we show that $(C, \mathcal{D}(C))$ is closed. Let $\left(f_{n}\right) \subset \mathcal{D}(C), g_{n}=C f_{n}$, and assume $f_{n} \rightarrow f$ and $g_{n} \rightarrow g$ in $L^{2}(\nu)$. Then $f_{n_{k}} \rightarrow f$ and $g_{n_{k}} \rightarrow g \nu$-almost everywhere for a suitable subsequence. Therefore

$$
f_{n_{k}}^{\prime}(x)=g_{n_{k}}(x)+U^{\prime}(x) f_{n_{k}}(x) \rightarrow g(x)+U^{\prime}(x) f(x), \quad v \text {-almost everywhere. }
$$

Now fix an open set $F \subset \mathbb{R}$ with compact closure. Since $U^{\prime}$ is bounded on $F$, it follows that the subsequence $\left(f_{n_{k}}^{\prime}\right)=$ $\left(g_{n_{k}}+U^{\prime} f_{n_{k}}\right)$ is Cauchy in $L^{2}(F)$. Thus $f \in W_{\mathrm{loc}}^{1,2}(\mathbb{R})$. Next

$$
\int_{\mathbb{R}}|C f-g|^{2} d \nu=\int_{\mathbb{R}} \lim _{k \rightarrow \infty}\left|C f_{n_{k}}-g_{n_{k}}\right|^{2} d \nu \stackrel{\text { Fatou }}{\leq} \lim _{k \rightarrow \infty} \int_{\mathbb{R}}\left|C f_{n_{k}}-g_{n_{k}}\right|^{2} d \nu=0 .
$$

Therefore $\|C f\|_{L^{2}(\nu)} \leq\|C f-g\|_{L^{2}(\nu)}+\|g\|_{L^{2}(\nu)}<\infty$ and $C f=g$.
Next we will show that $C_{c}^{\infty}(\mathbb{R})$ is a core for $(C, \mathcal{D}(C))$, in complete analogy to Lemma 3.2. Let $f \in \mathcal{D}(C)$. First take functions $\xi_{n}$ as in the proof of Lemma 3.2 and define $f_{n}(x):=f(x) \xi_{n}(x)$. Then $f_{n} \in \mathcal{D}(C), f_{n}$ has compact support and $f_{n} \rightarrow f$ in $L^{2}(\nu)$ by dominated convergence, and

$$
\left\|C f_{n}-C f\right\|_{L^{2}(\nu)} \leq\left\|(C f)\left(\xi_{n}-1\right)\right\|_{L^{2}(\nu)}+\left\|f \xi_{n}^{\prime}\right\|_{L^{2}(\nu)} \rightarrow 0 \quad \text { as } n \rightarrow \infty,
$$

the first term by dominated convergence and the second term since $\sup _{x \in \mathbb{R}}\left|\xi_{n}^{\prime}(x)\right| \leq c / n$ for some positive $c>0$ independent of $n$. This shows that compactly supported functions are dense in $\mathcal{D}(C)$ with respect to the graph norm.

Next take $\varphi_{n} \in C_{c}^{\infty}(\mathbb{R})$, again as in the proof of Lemma 3.2. Let $f \in \mathcal{D}(C)$ with compact support, and define $f_{n}(x)=$ $\left(f \star \varphi_{n}\right)(x)$. Then every $f_{n} \in C_{c}^{\infty}(\mathbb{R})$ is supported in $\operatorname{supp}(f)+B(1)$ and $f_{n} \rightarrow f$ in $L^{2}(\nu)$. Since $\partial_{x} f_{n}=\left(\partial_{x} f\right) \star \varphi_{n}$ and $U^{\prime}$ is bounded uniformly in $n$ on the support of $\left(f_{n}\right)$, it follows that $\left\|C f-C f_{n}\right\|_{L^{2}(\nu)} \rightarrow 0$.

## A.4. Estimates in the proof of Theorem 4.4

In the following lemma we verify that the constants that appear as integrals in the proof of Theorem 4.4 are well defined.
Lemma A.4. Suppose Assumption A2 holds. Let $h \in L^{2}(v)$. For any $\gamma \in \mathbb{C}$, we have that the following mappings are in $L^{1}(\mathbb{R})$.

$$
\begin{align*}
& \xi \mapsto U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)}  \tag{A.2}\\
& \xi \mapsto e^{-2 \gamma \xi-U(\xi)}  \tag{A.3}\\
& \xi \mapsto e^{-\gamma \xi-U(\xi)} h(\xi)  \tag{A.4}\\
& \xi \mapsto U^{\prime}(\xi) e^{-\gamma \xi} \int_{0}^{\xi} e^{\gamma \eta} h(\eta) d \eta \tag{A.5}
\end{align*}
$$

Proof. Write $\alpha=\operatorname{Re} \gamma$. It suffices to check integrability over $[0, \infty)$, as integrability over $(-\infty, 0]$ is analogous. Next, by continuity and where necessary local integrability of $h$ it suffices to check integrability over $\left[x_{0}, \infty\right)$ for some $x_{0}>0$. Let $x_{0}$ be sufficiently large, such that, for some constant $c>0$ such that $c+2 \alpha>0$, we have that $U^{\prime}(x) \geq c$ for $x \geq x_{0}$. We find

$$
\begin{aligned}
\int_{x_{0}}^{\infty}\left|U^{\prime}(\xi) e^{-2 \gamma \xi-U(\xi)}\right| d \xi & =\int_{x_{0}}^{\infty} U^{\prime}(\xi) e^{-2 \alpha \xi-U(\xi)} d \xi \\
& =e^{-U\left(x_{0}\right)}-2 \alpha \int_{0}^{\infty} e^{-U(\xi)-2 \alpha \xi} d \xi
\end{aligned}
$$

If $\alpha \geq 0$, we have already succeeded in establishing (A.2). Moreover, for any value of $\alpha \in \mathbb{R}$,

$$
\int_{x_{0}}^{\infty} e^{-U(\xi)-2 \alpha \xi} d \xi \leq \int_{x_{0}}^{\infty} e^{-U\left(x_{0}\right)-c\left(\xi-x_{0}\right)-2 \alpha \xi} d \xi<\infty
$$

establishing (A.2), (A.3). This yields, for (A.4), that

$$
\int_{0}^{\infty} e^{-\alpha \xi-U(\xi)}|h(\xi)| d \xi \leq\|h\|_{L^{2}(\nu)} \int_{0}^{\infty} e^{-2 \alpha \xi-U(\xi)} d \xi<\infty
$$

Finally, for (A.5),

$$
\begin{aligned}
& \int_{0}^{\infty}\left|U^{\prime}(\xi)\right| e^{-2 \alpha \xi-U(\xi)} \int_{0}^{\xi} e^{\alpha \eta}|h(\eta)| d \eta d \xi \\
& \quad=\int_{0}^{\infty} \int_{\eta}^{\infty}\left|U^{\prime}(\xi)\right| e^{-2 \alpha \xi-U(\xi)} d \xi e^{\alpha \eta+U(\eta)}|h(\eta)| e^{-U(\eta)} d \eta
\end{aligned}
$$

so it suffices (for this term) to check that

$$
\eta \mapsto \int_{\eta}^{\infty}\left|U^{\prime}(\xi)\right| e^{-2 \alpha \xi-U(\xi)} d \xi e^{\alpha \eta+U(\eta)} \in L^{2}(\nu)
$$

We have for $\eta \geq x_{0}$, by partial integration,

$$
\int_{\eta}^{\infty}\left|U^{\prime}(\xi)\right| e^{-2 \alpha \xi-U(\xi)} d \xi e^{\alpha \eta+U(\eta)}=e^{-\alpha \eta}-2 \alpha \int_{\eta}^{\infty} e^{-2 \alpha \xi-U(\xi)} d \xi e^{\alpha \eta+U(\eta)}
$$

The first term has already been establised above to belong to $L^{2}(\nu)$, the second term can be estimated using $U(\xi)-$ $U(\eta) \geq c(\xi-\eta)$ as

$$
\int_{\eta}^{\infty} e^{-2 \alpha \xi-U(\xi)} d \xi e^{\alpha \eta+U(\eta)} \leq \int_{\eta}^{\infty} e^{-2 \alpha \xi-c(\xi-\eta)} d \xi e^{\alpha \eta}=\frac{e^{-\alpha \eta}}{2 \alpha+c}<\infty .
$$

## A.5. Lemmas for the proof of Theorem 4.11

The following two lemmas are instrumental in establishing a spectral gap, i.e., Theorem 4.11.
Lemma A.5. Suppose the assumptions of Theorem 4.4 are satisfied. Then $\mathrm{i} \beta \in \sigma(L)$ for $\beta \in \mathbb{R}$ if and only if $\beta=0$. In particular $\sigma(L)$ does not have an accumulation point at $\mathrm{i} \beta$ for $\beta \in \mathbb{R}$.

Proof. First note that by Assumption A3, $U^{\prime}(0)=0$ and $x \mapsto U^{\prime}(x) e^{-U(x)}$ is a probability density function on $[0, \infty)$. We have

$$
\psi^{+}(\mathrm{i} \beta)=\int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \mathrm{i} \beta \xi-U(\xi)} d \xi=\int_{0}^{\infty} U^{\prime}(\xi)[\cos (2 \beta \xi)-\mathrm{i} \sin (2 \beta \xi)] e^{-U(\xi)} d \xi,
$$

so that, by Jensen's inequality,

$$
\begin{aligned}
\left|\psi^{+}(\mathrm{i} \beta)\right|^{2} & =\left(\int_{0}^{\infty} U^{\prime}(\xi) \cos (2 \beta \xi) e^{-U(\xi)} d \xi\right)^{2}+\left(\int_{0}^{\infty} U^{\prime}(\xi) \sin (2 \beta \xi) e^{-U(\xi)} d \xi\right)^{2} \\
& \leq \int_{0}^{\infty} U^{\prime}(\xi)\left[\cos ^{2}(2 \beta \xi)+\sin ^{2}(2 \beta \xi)\right] e^{-U(\xi)} d \xi=\int_{0}^{\infty} U^{\prime}(\xi) e^{-U(\xi)} d \xi=1
\end{aligned}
$$

with equality if and only if the integrands are constant, i.e., if and only if $\beta=0$. Similarly $\left|\psi^{-}(\mathrm{i} \beta)\right| \leq 1$ with equality if and only if $\beta=0$. It follows that $Z(\mathrm{i} \beta)=1-\psi^{+}(\mathrm{i} \beta) \psi^{-}(\mathrm{i} \beta)=0$ if and only if $\beta=0$.

Now suppose there is $\beta \neq 0$ and a sequence $\gamma_{n} \in \sigma(L)$ such that $\lim _{n \rightarrow \infty} \gamma_{n}=\mathrm{i} \beta$. By continuity of $Z$ (Lemma 4.8), it follows that $Z(\mathrm{i} \beta)=0$, a contradiction. Furthermore, since 0 an eigenvalue, and the spectrum consists of isolated eigenvalues only, 0 cannot be an accumulation point either.

Lemma A.6. Suppose the assumptions of Theorem 4.4 are satisfied. For every closed and bounded interval $[a, b] \subset \mathbb{R}$ there is a constant $C>0$ such that

$$
\left|\psi^{ \pm}(\alpha+\mathrm{i} \beta)\right| \leq \frac{C}{|\beta|} \quad \text { for all } \alpha \in[a, b] \text { and } \beta \neq 0 .
$$

In particular, $\alpha+\mathrm{i} \beta \in \rho(L)$ for $\alpha \in[a, b],|\beta|>C$.
Proof. By partial integration, and using that by Assumption A1, $\left|U^{\prime \prime}(\xi)\right| \leq c_{1}+c_{2}\left|U^{\prime}(\xi)\right|^{2}$, for some constants $c_{1}, c_{2}>0$,

$$
\begin{aligned}
\left|\psi^{+}(\alpha+\mathrm{i} \beta)\right| & =\left|\int_{0}^{\infty} U^{\prime}(\xi) e^{-2 \alpha \xi-2 \mathrm{i} \beta \xi-U(\xi)} d \xi\right| \\
& =\left|\frac{1}{2 \mathrm{i} \beta} \int_{0}^{\infty}\left[U^{\prime \prime}(\xi)-2 \alpha U^{\prime}(\xi)-\left(U^{\prime}(\xi)\right)^{2}\right] e^{-2 \alpha \xi-2 \mathrm{i} \beta \xi-U(\xi)} d \xi\right| \\
& \leq \frac{1}{2|\beta|} \int_{0}^{\infty}\left|U^{\prime \prime}(\xi)-2 \alpha U^{\prime}(\xi)-\left(U^{\prime}(\xi)\right)^{2}\right| e^{-2 \alpha \xi-U(\xi)} d \xi \\
& \leq \frac{1}{2|\beta|} \int_{0}^{\infty}\left[c_{1}+2 \alpha^{2}+\left(c_{2}+3 / 2\right)\left(U^{\prime}(\xi)\right)^{2}\right] e^{-2 \alpha \xi-U(\xi)} d \xi,
\end{aligned}
$$

and the integral converges by an application of Lemma 3.7. It follows that, for $\alpha \in[a, b]$,

$$
\left|\psi^{+}(\alpha+\mathrm{i} \beta)\right| \leq \frac{1}{2|\beta|} \int_{0}^{\infty}\left[c_{1}+2 \max (|a|,|b|)^{2}+\left(c_{2}+3 / 2\right)\left(U^{\prime}(\xi)\right)^{2}\right] e^{-2 a \xi-U(\xi)} d \xi=: C^{+} /|\beta| .
$$

The result for $\psi^{-}$follows analogously, yielding a constant $C^{-}$, and we may take $C$ as the maximum of the two constants obtained. Thus for $|\beta|>C$, we have that

$$
\left|\psi^{+}(\alpha+\mathrm{i} \beta) \psi^{-}(\alpha+\mathrm{i} \beta)\right| \leq C^{2} / \beta^{2}<1
$$

so that $Z(\alpha+\mathrm{i} \beta) \neq 0$ and therefore $\alpha+\mathrm{i} \beta \in \rho(L)$.

## A.6. A convolution estimate

In the proof of Proposition 4.12 we will make repeated use of the following lemma.
Lemma A.7. Suppose Assumption $A 4$ is satisfied. Let $\phi: \mathbb{R} \rightarrow \mathbb{C}$ be Lebesgue measurable and $\psi \in L^{2}(v)$. Furthermore assume that $\phi(x)=\psi(x)=0$ for almost every $x<0$. Let

$$
(\phi \star \psi)(x):=\int_{0}^{x} \phi(x-y) \psi(y) d y, \quad \text { and } \quad \zeta(x):=e^{U(x)} \int_{x}^{\infty} \phi(y-x) \psi(y) e^{-U(y)} d y, \quad x \geq 0
$$

Then

$$
\|\phi \star \psi\|_{L^{2}(\nu)} \leq\left(\int_{0}^{\infty}|\phi(x)| e^{-C / 2-(m / 2)|x|^{p}} d x\right)\|\psi\|_{L^{2}(\nu)}
$$

and

$$
\|\zeta\|_{L^{2}(\nu)} \leq\left(\int_{0}^{\infty}|\phi(x)| e^{-C / 2-(m / 2)|x|^{p}} d x\right)\|\psi\|_{L^{2}(\nu)}
$$

Proof. By Lemma A. 2 we may assume without loss of generality that Assumption A4 holds with $M=0$, i.e., there are constants $C \leq 0, m>0$ and $p>1$ such that

$$
U(y) \geq U(x)+C+m|x-y|^{p} \quad \text { for all } x, y \text { for which } 0 \leq x \leq y \text { or } y \leq x \leq 0
$$

Write $\hat{\psi}(x)=e^{-U(x) / 2} \psi(x)$, so that $\|\hat{\psi}\|_{L^{2}(\mathbb{R})}=\|\psi\|_{L^{2}(\nu)}$. By Assumption A4,

$$
\begin{aligned}
e^{-U(x) / 2} \int_{0}^{x} \phi(x-y) \psi(y) d y & =e^{-U(x) / 2} \int_{0}^{x} \phi(x-y) \hat{\psi}(y) e^{U(y) / 2} d y \\
& \leq \int_{0}^{x}|\phi(x-y)| e^{-C / 2-(m / 2)|y-x|^{p}}|\hat{\psi}(y)| d y \\
& =(|\hat{\phi}| \star|\hat{\psi}|)(x)
\end{aligned}
$$

where

$$
\hat{\phi}(x)=\phi(x) e^{-C / 2-(m / 2)|x|^{p}}
$$

The first result of the lemma follows now from the convolutional inequality $\|f \star g\|_{L^{2}(\mathbb{R})} \leq\|f\|_{L^{1}(\mathbb{R})}\|g\|_{L^{2}(\mathbb{R})}$. Next

$$
\begin{aligned}
e^{-U(x) / 2} \zeta(x) & =e^{U(x) / 2} \int_{x}^{\infty} \phi(y-x) \hat{\psi}(y) e^{-U(y) / 2} d y \\
& \leq \int_{x}^{\infty}|\phi(y-x)| e^{-C / 2-(m / 2)|x-y|^{p} / 2}|\hat{\psi}(y)| d y \\
& =(|\tilde{\phi}| \star|\hat{\psi}|)(x)
\end{aligned}
$$

where

$$
\tilde{\phi}(x)=\phi(-x) e^{-C / 2-(m / 2)|x|^{p}}
$$

which yields the second statement after applying the convolutional inequality again.

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## References

[1] C. Andrieu, A. Durmus, N. Nüsken and J. Roussel Hypercoercivity of piecewise deterministic Markov process-Monte Carlo. arXiv preprint, 2018. Available at arXiv:1808.08592.
[2] J. Bierkens. Numerical Computation of Zigzag Semigroup Spectrum. GitHub Repository https://github.com/jbierkens/spectral, 2019.
[3] J. Bierkens and A. Duncan. Limit theorems for the zig-zag process. Adv. in Appl. Probab. 49 (3) (2017) 791-825. MR3694318 https://doi.org/10. 1017/apr. 2017.22
[4] J. Bierkens, P. Fearnhead and G. O. Roberts. The zig-zag process and super-efficient sampling for Bayesian analysis of big data. Ann. Statist. 47 (3) (2019) 1288-1320. MR3911113 https://doi.org/10.1214/18-AOS1715
[5] J. Bierkens and G. Roberts. A piecewise deterministic scaling limit of lifted Metropolis-Hastings in the Curie-Weiss model. Ann. Appl. Probab. 27 (2) (2017) 846-882. MR3655855 https://doi.org/10.1214/16-AAP1217
[6] A. Bouchard-Côté, S. J. Vollmer and A. Doucet. The bouncy particle sampler: A non-reversible rejection-free Markov chain Monte Carlo method. J. Amer. Statist. Assoc. (2017). MR3832232 https://doi.org/10.1080/01621459.2017.1294075
[7] E. B. Davies. Linear Operators and Their Spectra. Cambridge Studies in Advanced Mathematics 106. Cambridge University Press, Cambridge, 2007. MR2359869 https://doi.org/10.1017/CBO9780511618864
[8] M. Dellnitz, O. Schütze and Q. Zheng. Locating all the zeros of an analytic function in one complex variable. J. Comput. Appl. Math. 138 (2) (2002) 325-333. MR1876239 https://doi.org/10.1016/S0377-0427(01)00371-5
[9] J. Dolbeault, C. Mouhot and C. Schmeiser. Hypocoercivity for linear kinetic equations conserving mass. Trans. Amer. Math. Soc. 367 (6) (2015) 3807-3828. MR3324910 https://doi.org/10.1090/S0002-9947-2015-06012-7
[10] K.-J. Engel and R. Nagel. One-Parameter Semigroups for Linear Evolution Equations. Graduate Texts in Mathematics 194. Springer-Verlag, New York, 2000. MR1721989
[11] H. E. Fettis, J. C. Caslin and K. R. Cramer. Complex zeros of the error function and of the complementary error function. Math. Comp. 27 (122) (1973) 401-407. MR0326991 https://doi.org/10.2307/2005630
[12] J. Fontbona, H. Guérin and F. Malrieu. Quantitative estimates for the long-time behavior of an ergodic variant of the telegraph process. Adv. in Appl. Probab. 44 (4) (2012) 977-994. MR3052846
[13] J. Fontbona, H. Guérin and F. Malrieu. Long time behavior of telegraph processes under convex potentials. Stochastic Process. Appl. 126 (10) (2015) 1-26. MR3542627 https://doi.org/10.1016/j.spa.2016.04.002
[14] L. Gearhart. Spectral theory for contraction semigroups on Hilbert space. Trans. Amer. Math. Soc. 236 (1978) 385-394. MR0461206 https://doi.org/10.2307/1997792
[15] M. A. Kaashoek and S. M. Verduyn Lunel. Characteristic matrices and spectral properties of evolutionary systems. Trans. Amer. Math. Soc. 334 (2) (1992) 479-517. MR1155350 https://doi.org/10.2307/2154470
[16] M. Kac. On some connections between probability theory and differential and integral equations. In Proceedings of the Second Berkeley Symposium on Mathematical Statstics and Probability 189, 1951. MR0045333
[17] T. Kato. Perturbation Theory for Linear Operators. Classics in Mathematics. Springer-Verlag, Berlin, 1995. MR 1335452
[18] L. Lorenzi and M. Bertoldi Analytical methods for Markov semigroups. In Pure and Applied Mathematics (Boca Raton) 283. Chapman \& Hall/CRC, Boca Raton, FL, 2007. MR2313847
[19] J. Lu and L. Wang On explicit $L^{2}$-convergence rate estimate for piecewise deterministic Markov processes. arXiv preprint, 2020. Available at arXiv:2007.14927.
[20] L. Miclo and P. Monmarché. Étude Spectrale Minutieuse de Processus Moins Indécis Que les Autres 459-481. Springer International Publishing, Heidelberg, 2013. MR3185926 https://doi.org/10.1007/978-3-319-00321-4_18
[21] P. Monmarché. Piecewise deterministic simulated annealing. ALEA 13 (1) (2016) 357-398. MR3487077
[22] E. A. J. F. Peters and G. De With. Rejection-free Monte Carlo sampling for general potentials. Phys. Rev. E 85 (2) (2012) 1-5.
[23] G. O. Roberts and R. L. Tweedie. Geometric $L^{2}$ and $L^{1}$ convergence are equivalent for reversible Markov chains. J. Appl. Probab. 38 (A) (2001) 37-41. MR1915532 https://doi.org/10.1239/jap/1085496589
[24] J. S. Rosenthal. Asymptotic variance and convergence rates of nearly-periodic Markov chain Monte Carlo algorithms. J. Amer. Statist. Assoc. 98 (461) (2003) 169-177. MR $1965683 \mathrm{https}: / /$ doi.org/10.1198/016214503388619193
[25] M. Vialaret and F. Maire. On the convergence time of some non-reversible Markov chain Monte Carlo methods. Methodol. Comput. Appl. Probab. (2020) 1-39. MR4129137 https://doi.org/10.1007/s1 1009-019-09766-w

