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Pin groups in general relativity

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There are eight possible Pin groups that can be used to describe the transformation behavior of fermions under parity and time reversal. We show that only *two* of these are compatible with general relativity, in the sense that the configuration space of fermions coupled to gravity transforms appropriately under the space-time diffeomorphism group.

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I. INTRODUCTION

For bosons, the space-time transformation behavior is governed by the Lorentz group $O(3, 1)$, which comprises four connected components. Rotations and boosts are contained in the connected component of unity, the proper orthochronous Lorentz group $SO^\uparrow(3, 1)$. Parity (P) and time reversal (T) are encoded in the other three connected components of the Lorentz group, the translates of $SO^\uparrow(3, 1)$ by P , T and PT .

For fermions, the space-time transformation behavior is governed by a double cover of $O(3, 1)$. Rotations and boosts are described by the unique simply connected double cover of $SO^\uparrow(3, 1)$, the spin group $\text{Spin}^\uparrow(3, 1)$. However, in order to account for parity and time reversal, one needs to extend this cover from $SO^\uparrow(3, 1)$ to the full Lorentz group $O(3, 1)$.

This extension is by no means unique. There are no less than *eight* distinct double covers of $O(3, 1)$ that agree with $\text{Spin}^\uparrow(3, 1)$ over $SO^\uparrow(3, 1)$. They are the Pin groups Pin^{abc} , characterized by the property that the elements Λ_P and Λ_T covering P and T satisfy $\Lambda_P^2 = -a$, $\Lambda_T^2 = b$ and $(\Lambda_P\Lambda_T)^2 = -c$, where a , b and c are either 1 or -1 (cf. [1,2]).

In this paper, we show that the consistent description of fermions in the presence of general relativity (GR) imposes severe restrictions on the choice of Pin group. In fact, we find that only *two* of the eight Pin groups are admissible: the group $\text{Pin}^+ = \text{Pin}^{+++}$ and the group $\text{Pin}^- = \text{Pin}^{---}$. The source of these restrictions is the double cover of the frame bundle, which, in the context of GR, is needed in order to obtain an infinitesimal action of the space-time diffeomorphism group on the configuration space of fermions coupled to gravity.

We derive these restrictions in the “universal spinor bundle approach” for fermions coupled to gravity, as developed in [3–5] for the Riemannian and in [6–9] for the Lorentzian case. However, our results remain valid in other formulations that are covariant under infinitesimal space-time diffeomorphisms, such as the “global” approach of [2,10–12]. To underline this point, we highlight the role of the space-time diffeomorphism group in restricting the admissible Pin groups.

Selecting the correct Pin groups is important from a fundamental point of view—it determines the transformation behavior of fermionic fields under reflections—but also because the Pin group can affect observable quantities such as currents [13–15]. Due to their transparent definition in terms of Clifford algebras, the “Cliffordian” Pin groups $\text{Pin}(3, 1) = \text{Pin}^{+++}$ and $\text{Pin}(1, 3) = \text{Pin}^{---}$ have attracted much attention [13,16–19]. Remarkably, the two Pin groups Pin^+ and Pin^- that are compatible with GR are *not* the widely used Cliffordian Pin groups $\text{Pin}(3, 1)$ and $\text{Pin}(1, 3)$.

II. THE LORENTZIAN METRIC

In order to establish notation, we briefly recall the *frame* or *vierbein formalism* for a Lorentzian metric g on a four-dimensional space-time manifold M .

A *frame* e_x based at x is a basis $e_a^\mu \partial_\mu$ of the tangent space $T_x M$, with basis vectors labeled by $a = 0, 1, 2, 3$. The space $F(M)$ of all frames (with arbitrary x) is called the *frame bundle*, and we denote by $F_x(M)$ the set of frames with base point x . Note that the group $\text{Gl}(4, \mathbb{R})$ of invertible 4×4 matrices A_b^a acts from the right on $F_x(M)$, sending e_x to the frame $e'_x = e_x A$ with $e_a'^\mu = e_b^\mu A_b^a$. This action is free and transitive; any two frames e_x and e'_x over the same point x are related by $e'_x = e_x A$ for a unique matrix A_b^a .

For a given Lorentzian metric g , the *orthonormal frame bundle* $O^g(M) \subset F(M)$ is the space of all orthonormal frames e_a^μ , satisfying $g_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$. Since two orthonormal frames e_x and e'_x over the same point x differ by a Lorentz transformation Λ , $e'_x = e_x \Lambda$, the Lorentz group $O(3, 1)$ acts

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freely and transitively on the set $O_x^g(M) \subset F_x(M)$ of orthonormal frames based at x .

Specifying a metric g at x is equivalent to specifying the set $O_x^g(M)$ of orthonormal frames. Since $O_x^g(M) \subset F_x(M)$ is an orbit under the action of the Lorentz group $O(3, 1)$ on $F_x(M)$, specifying the metric at x is equivalent to picking a point in the *orbit space* $\mathcal{R}_x(M) = F_x(M)/O(3, 1)$. This is the set of equivalence classes $[e_x]$ of frames at x , where two frames e_x and e'_x are deemed equivalent if they differ by a Lorentz transformation Λ , $e'_x = e_x \Lambda$. We denote the bundle of all equivalence classes $[e_x]$ (with arbitrary x) by $\mathcal{R}(M)$.

To describe fermions in the presence of GR, it will be convenient to view a metric g on M as a *section* of $\mathcal{R}(M)$; a smooth map $g: M \rightarrow \mathcal{R}(M)$ that takes a point x to an equivalence class $[e_x]$ of frames at x . The configuration space [20] of general relativity can thus be seen as the space $\Gamma(\mathcal{R}(M))$ of sections of the bundle $\mathcal{R}(M)$.

III. FERMIONIC FIELDS IN A FIXED BACKGROUND

We start by describing fermionic fields on M in the presence of a fixed background metric g . In order to do this, a number of choices have to be made, especially if we wish to keep track of the transformation behavior of spinors under parity and time reversal.

The local transformation behavior is fixed by choosing one out of the eight possible Pin groups Pin^{abc} , together with a (not necessarily \mathbb{C} -linear) representation V that extends the spinor representation of $\text{Spin}^\uparrow(3, 1) \subset \text{Pin}^{abc}$. For example, V consists of n copies of \mathbb{C}^4 in the case of n Dirac fermions, and it consists of m copies of \mathbb{C}^2 in the case of m Majorana fermions [21].

Once a Pin group has been selected, the second choice one has to make is a choice of *Pin structure*. A Pin structure is a twofold cover $u: Q^g \rightarrow O^g(M)$ of the orthonormal frame bundle, equipped with a Pin^{abc} -action that is compatible with the action of the Lorentz group on $O^g(M)$. The compatibility entails that if $\tilde{\Lambda} \in \text{Pin}^{abc}$ covers $\Lambda \in O(3, 1)$, then $u(q_x \tilde{\Lambda}) = u(q_x) \Lambda$ for all pin frames q_x in Q^g . A pin frame q_x is based at the same point as its image, the frame $u(q_x)$. We denote by Q_x^g the set of pin frames based at x .

For a given manifold M and a given Pin group Pin^{abc} , a Pin structure may or may not exist, and if it does, it need not be unique. The obstruction theory for this problem has been completely solved for the Cliffordian Pin groups in [22], and for the general case in [1].

Once a Pin structure Q^g has been chosen, one can construct the associated bundle $S^g = (Q^g \times V)/\text{Pin}^{abc}$ of *spinors*. A spinor $\psi_x = [q_x, \vec{v}]$ at x is thus an equivalence class of a pin frame $q_x \in Q_x^g$ and a vector $\vec{v} \in V$, where $(q_x \tilde{\Lambda}, \vec{v})$ is identified with $(q_x, \tilde{\Lambda} \vec{v})$ for every element $\tilde{\Lambda}$ of the Pin group Pin^{abc} .

For a given background metric g , the fermionic fields are then described by sections of the spinor bundle S^g , that is,

by smooth maps $\psi: M \rightarrow S^g$ that assign to each space-time point x a spinor ψ_x based at x . The configuration space for the fermionic fields at a fixed metric g is thus the space $\Gamma(S^g)$ of sections of the spinor bundle S^g .

IV. FERMIONIC FIELDS COUPLED TO GR

We now wish to describe the configuration space for fermionic fields coupled to gravity. This is *not* simply the product of the configuration space of general relativity and that of a fermionic field; the main difficulty here is that the very space S^g where the spinor field ψ takes values depends on the metric g . A solution to this problem was proposed in [3,4] for the Riemannian case, and in [6–9] for metrics of Lorentzian signature. In order to handle reflections, we need to adapt this procedure as follows.

First, we choose a twofold cover of $\text{Gl}(4, \mathbb{R})$ that agrees with the universal cover $\tilde{\text{Gl}}_+(4, \mathbb{R})$ over $\text{Gl}_+(4, \mathbb{R})$. In Sec. V we show that there are only *two* such covers, which, for want of a better name, we will call Gin^+ and Gin^- . Having made our choice of Gin^\pm , we choose what one may call a *Gin structure*; a twofold cover $u: \hat{Q} \rightarrow F(M)$ with a Gin^\pm -action that is compatible with the $\text{Gl}(4, \mathbb{R})$ -action on $F(M)$. Corresponding to every (not necessarily orthogonal) frame e_x , there are thus two gin frames \hat{q}_x and \hat{q}'_x . If $\tilde{A} \in \text{Gin}^\pm$ covers $A \in \text{Gl}(4, \mathbb{R})$, then the two gin frames corresponding to $e_x A$ are $\hat{q}_x \tilde{A}$ and $\hat{q}'_x \tilde{A}$.

We denote by Pin^\pm the twofold cover of $O(3, 1)$ inside Gin^\pm . Choosing a Gin structure \hat{Q} for the group Gin^\pm is equivalent to choosing a Pin structure Q^g for the group Pin^\pm . Indeed, for every Gin^\pm structure \hat{Q} , the preimage $Q^g \subset \hat{Q}$ of $O^g(M) \subset F(M)$ under the map $u: \hat{Q} \rightarrow F(M)$ is a Pin^\pm -structure, since the restriction $u^g: Q^g \rightarrow O^g(M)$ of u to Q^g intertwines the Pin^\pm -action on Q^g with the action of the Lorentz group $O(3, 1)$ on $O^g(M)$. Conversely, every Pin^\pm -structure $u: Q^g \rightarrow O^g(M)$ gives rise to the associated Gin^\pm -structure $\hat{Q} = (Q^g \times \text{Gin}^\pm)/\text{Pin}^\pm$. This is the space of equivalence classes $[q_x, \tilde{A}]$, where $(q_x \tilde{\Lambda}, \tilde{A})$ is identified with $(q_x, \tilde{\Lambda} \tilde{A})$ for every $\tilde{\Lambda}$ in Pin^\pm . The obstruction theory for Gin^\pm -structures therefore reduces to the obstruction theory for Pin^\pm -structures, which has been worked out in [1].

Using the Gin structure \hat{Q} , one constructs the *universal spinor bundle* $\Sigma = (\hat{Q} \times V)/\text{Pin}^\pm$ in analogy with [9]. A universal spinor $\Psi_x = [\hat{q}_x, \vec{v}]$ at x is an equivalence class of a gin frame $\hat{q}_x \in \hat{Q}_x$ and a vector $\vec{v} \in V$, where $(\hat{q}_x \tilde{\Lambda}, \vec{v})$ is identified with $(\hat{q}_x, \tilde{\Lambda} \vec{v})$ for every $\tilde{\Lambda}$ in Pin^\pm . Note that a universal spinor Ψ_x in $\Sigma = (\hat{Q} \times V)/\text{Pin}^\pm$ defines a metric $g_{\mu\nu}$ at x , together with a spinor ψ_x in the spinor bundle $S^g = (Q^g \times V)/\text{Pin}^\pm$ that corresponds with the metric $g_{\mu\nu}$ induced by Ψ_x .

Indeed, since the covering map $u: \hat{Q} \rightarrow F(M)$ intertwines the Pin^\pm -action on \hat{Q} with the $\text{Gl}(4, \mathbb{R})$ -action on $F(M)$, it identifies the quotient of \hat{Q} by Pin^\pm with the

quotient of $F(M)$ by $O(3,1)$, which is the orbit space $\mathcal{R}(M)$. From a universal spinor $\Psi_x = [\hat{q}_x, \vec{v}]$ at x , we thus obtain an equivalence class $[u(\hat{q}_x)]$ in $\mathcal{R}_x(M)$, and hence a metric $g_{\mu\nu}$ at the point x .

To obtain not only the metric $g_{\mu\nu}$ but also the spinor ψ_x , recall that the Pin structure Q^g corresponding to $g_{\mu\nu}$ is the preimage of $O^g(M)$ under the double cover $u: \hat{Q} \rightarrow F(M)$. Since $Q^g \subset \hat{Q}$ contains the gin frame \hat{q}_x , the equivalence class $\Psi_x = [\hat{q}_x, \vec{v}]$ in $\Sigma = (\hat{Q} \times V)/\text{Pin}^\pm$ yields an equivalence class $\psi_x = [q_x, \vec{v}]$ in the spinor bundle $S^g = (Q^g \times V)/\text{Pin}^\pm$ by setting $q_x = \hat{q}_x$. Here, S^g is the spinor bundle derived from the metric $g_{\mu\nu}$ that is induced by Ψ .

We conclude that *both* the metric g and the fermionic field ψ are described by a *single* section $\Psi: M \rightarrow \Sigma$, a smooth map assigning to each point x of space-time a universal spinor Ψ_x based at x . The configuration space of fermionic fields coupled to gravity is thus the space $\Gamma(\Sigma)$ of sections of the universal spinor bundle Σ .

V. COVERING GROUPS

Out of the eight Pin groups covering $O(3,1)$, the only two that are compatible with this formalism are the twofold cover Pin^+ of $O(3,1)$ inside Gin^+ , and the twofold cover Pin^- of $O(3,1)$ inside Gin^- . We show that their coefficients in the sense of Sec. I are $(a, b, c) = (+, +, -)$ and $(a, b, c) = (-, -, -)$.

First we show that there are only two double covers of $\text{Gl}(4, \mathbb{R})$ that reduce to the universal cover over $\text{Gl}_+(4, \mathbb{R})$. Assume that G is such a cover. If Λ_T is an element of G that covers the time reversal operator $T \in \text{Gl}(4, \mathbb{R})$, then the automorphism $\text{Ad}_{\Lambda_T}(\tilde{A}) := \Lambda_T \tilde{A} \Lambda_T^{-1}$ of $\tilde{\text{Gl}}_+(4, \mathbb{R})$ covers the automorphism $\text{Ad}_T(A) := T A T^{-1}$ of $\text{Gl}_+(4, \mathbb{R})$. By the universal covering property, $\Lambda_T \tilde{A} \Lambda_T^{-1}$ is uniquely determined by \tilde{A} , and it depends neither on the choice of G , nor on the choice of Λ_T inside G . Since every element of G can be written as either \tilde{A} or $\tilde{B} \Lambda_T$, there are four types of products, namely those of the form $\tilde{A} \tilde{A}'$, $\tilde{A}(\tilde{B} \Lambda_T)$, $(\tilde{B} \Lambda_T) \tilde{A}$ and $(\tilde{B} \Lambda_T)(\tilde{B}' \Lambda_T)$, where $\tilde{A}, \tilde{A}', \tilde{B}, \tilde{B}'$ are in $\tilde{\text{Gl}}_+(4, \mathbb{R})$. Products of the first 2 types are determined by the group structure on $\tilde{\text{Gl}}_+(4, \mathbb{R})$. This is true for the third type as well, since $(\tilde{B} \Lambda_T) \tilde{A} = \tilde{B}(\Lambda_T \tilde{A} \Lambda_T^{-1}) \Lambda_T$, and $\Lambda_T \tilde{A} \Lambda_T^{-1}$ is independent of G . As $(\tilde{B} \Lambda_T)(\tilde{B}' \Lambda_T) = (\tilde{B}(\Lambda_T \tilde{B}' \Lambda_T^{-1})) \Lambda_T^2$, the only choice in the product structure on G lies in the sign of $\Lambda_T^2 = \pm 1$, yielding the two groups Gin^\pm . The twofold cover Pin^+ of $O(3,1)$ inside Gin^+ thus has $b = +1$, whereas the twofold cover Pin^- inside Gin^- has $b = -1$.

To establish that both Pin^+ and Pin^- satisfy $c = -1$, note that although the central element $PT = \text{diag}(-1, -1, -1, -1)$ does not lie in the connected component of unity for the Lorentz group $O(3,1)$, it does lie in the connected subgroup $\text{SO}(4)$ of $\text{Gl}_+(4, \mathbb{R})$. As the inverse image of $\text{SO}(4)$ under the universal cover $\tilde{\text{Gl}}_+(4, \mathbb{R}) \rightarrow \text{Gl}_+(4, \mathbb{R})$

is its universal cover $\text{Spin}^\uparrow(4)$, the square of $\Lambda_P \Lambda_T$ inside $\tilde{\text{Gl}}_+(4, \mathbb{R})$ equals its square in $\text{Spin}^\uparrow(4)$. Here, the elements $\pm i\gamma_5 = \mp \gamma_0 \gamma_1 \gamma_2 \gamma_3$ that cover PT square to $+1$, as one easily derives using the Clifford relations $\{\gamma_\mu, \gamma_\nu\} = 2\delta_{\mu\nu}$ for the Euclidean gamma matrices γ_μ . It follows that $(\Lambda_P \Lambda_T)^2 = 1$, and hence $c = -1$.

It remains to show that $a = b$. For this, note that the restriction of the automorphism Ad_{Λ_T} of $\tilde{\text{Gl}}_+(4, \mathbb{R})$ to the simply connected subgroup $\text{Spin}^\uparrow(4) \subset \tilde{\text{Gl}}_+(4, \mathbb{R})$ is uniquely determined by its induced Lie algebra automorphism. On $\text{Spin}^\uparrow(4)$, we thus have $\text{Ad}_{\Lambda_T}(u) = \gamma_0 u \gamma_0^{-1}$. As $\gamma_0(i\gamma_5)\gamma_0^{-1} = -i\gamma_5$, we find that $\Lambda_T \Lambda_P = \text{Ad}_{\Lambda_T}(\Lambda_P \Lambda_T) = -\Lambda_P \Lambda_T$. As we already established that $(\Lambda_P \Lambda_T)^2 = 1$, it follows that $\Lambda_P^2 \Lambda_T^2 = -1$, and hence that $a = -\Lambda_P^2$ is equal to $b = \Lambda_T^2$. We thus conclude that $(a, b, c) = (+, +, -)$ for Pin^+ , and $(a, b, c) = (-, -, -)$ for Pin^- .

The groups Pin^+ and Pin^- are therefore *not* isomorphic to the Cliffordian Pin groups $\text{Pin}(3,1)$ and $\text{Pin}(1,3)$. These are generated by the Clifford elements $v^\mu \tilde{\gamma}_\mu$ with $\eta_{\mu\nu} v^\mu v^\nu = \pm 1$, where the Lorentzian gamma matrices $\tilde{\gamma}_\mu$ satisfy $\{\tilde{\gamma}_\mu, \tilde{\gamma}_\nu\} = 2\eta_{\mu\nu}$ for $\text{Pin}(3,1)$, and $\{\tilde{\gamma}_\mu, \tilde{\gamma}_\nu\} = -2\eta_{\mu\nu}$ for $\text{Pin}(1,3)$. Since the group elements covering P and T are $\Lambda_P = \tilde{\gamma}_1 \tilde{\gamma}_2 \tilde{\gamma}_3$ and $\Lambda_T = \tilde{\gamma}_0$, one readily verifies that $(a, b, c) = (+, -, +)$ for $\text{Pin}(3,1)$, and that $(a, b, c) = (-, +, +)$ for $\text{Pin}(1,3)$ (cf. [1,2]).

In particular, we conclude that the two Pin groups Pin^\pm compatible with GR are *not* the widely used Cliffordian Pin groups $\text{Pin}(3,1)$ and $\text{Pin}(1,3)$.

VI. TRANSFORMATION UNDER DIFFEOMORPHISMS

In the above derivation of the two admissible Pin groups, a crucial role is played by the continuous covering map $u: \hat{Q} \rightarrow F(M)$. This map has physical significance, since it induces an infinitesimal action of the space-time diffeomorphism group $\text{Diff}(M)$ on the configuration space of fermions coupled to gravity (cf. [3,8]). This allows one to formulate a theory which is (up to sign) covariant under general coordinate transformations (cf. [3,11]), and to construct a stress-energy-momentum tensor via Noether's theorem (cf. [23,24], and cf. [6], Sec. 6) for an approach using variation of the metric).

To construct the infinitesimal action, note that $\text{Diff}(M)$ acts by automorphisms on the frame bundle $F(M)$, a diffeomorphism ϕ maps $e_x \in F_x(M)$ to $D\phi(e_x) := \partial_\mu \phi^\mu e_a^a$ in $F_{\phi(x)}(M)$. A one-parameter group ϕ_ϵ of diffeomorphisms thus yields a one-parameter group $D\phi_\epsilon$ of automorphisms of $F(M)$. Since $u: \hat{Q} \rightarrow F(M)$ is a double cover, this lifts to a unique one-parameter group $D\hat{\phi}_\epsilon$ of automorphisms of \hat{Q} . On the universal spinor bundle $\Sigma = (\hat{Q} \times V)/\text{Pin}^\pm$, we define the lift by $D\hat{\phi}_\epsilon[\hat{q}_x, \vec{v}] = [D\hat{\phi}_\epsilon(\hat{q}_x), \vec{v}]$. For the infinitesimal variation of the universal spinor field $\Psi: M \rightarrow \Sigma$ along ϕ_ϵ , this yields $\delta\Psi_x = \frac{d}{d\epsilon}|_0 D\hat{\phi}_\epsilon(\Psi_{\phi_\epsilon^{-1}(x)})$.

VII. THE ROLE OF DIFFEOMORPHISMS IN RESTRICTING THE PIN GROUPS

We stress that the above restrictions on the Pin groups are not needed to *construct* the configuration space for fermions coupled to gravity, but to ensure that it *transforms* appropriately under space-time diffeomorphisms.

Indeed, to construct the configuration space, one could simply choose any principal Pin^{abc} -bundle $P \rightarrow \mathcal{R}(M)$ (for example the trivial one), and construct the universal spinor bundle $\Sigma = (P \times V)/\text{Pin}^{abc}$ as in Sec. IV. Its sections $\Psi \in \Gamma(\Sigma)$ can be interpreted as a fermionic field ψ together with a metric g , so $\Gamma(\Sigma)$ may serve as a configuration space. This requires no restrictions on the Pin groups, nor on the topology of M .

However, this simple construction leaves the space-time transformation behavior undetermined. We show that the restrictions on the Pin groups are recovered by imposing appropriate transformation behavior on $\Gamma(\Sigma)$. Compatibility with the Lorentz group leads to the familiar restrictions on the topology of M , compatibility with infinitesimal diffeomorphisms leads to Pin groups with $c = -1$, and compatibility with a double cover of the diffeomorphism group requires Pin groups with $a = b$ as well as $c = -1$.

A. Lorentz transformations

The *pullback* of a bundle $E \rightarrow Y$ along a map $f: X \rightarrow Y$ is the bundle $f^*E \rightarrow X$ with $(f^*E)_x := E_{f(x)}$. Starting from the principal bundle $P \rightarrow \mathcal{R}(M)$, one thus obtains for every metric $g: M \rightarrow \mathcal{R}(M)$ a principal Pin^{abc} -bundle $g^*P \rightarrow M$. Its fibre g^*P_x at x is the fibre $P_{g(x)}$ of P at $g(x) \in \mathcal{R}(M)$. The bundle g^*P is not quite a Pin structure, since the action of Pin^{abc} on g^*P is as yet unrelated to the action of $\text{O}(3, 1)$ on $O^g(M)$. To define the transformation behavior of Ψ under infinitesimal isometries, we need to choose a Pin structure on each of the bundles g^*P . That is, for any possible metric $g \in \Gamma(\mathcal{R}(M))$, we need to choose a double cover $u^g: g^*P \rightarrow O^g(M)$ that intertwines the action of Pin^{abc} on g^*P with the action of $\text{O}(3, 1)$ on $O^g(M)$. This is where the restrictions on the topology of M arise: if the conditions in [1] are met, then it is possible to endow every *single* bundle $g^*P \rightarrow M$ with a double covering map $u^g: g^*P \rightarrow O^g(M)$, making it into a Pin structure.

B. Infinitesimal diffeomorphisms

The problem is that, in general, these covering maps u^g do not depend continuously on the metric g . If we require this to be the case, then we recover the infinitesimal action of the diffeomorphism group on the configuration space, as well as the restriction $c = -1$ on the Pin groups. This already excludes the ‘‘Cliffordian’’ Pin groups $\text{Pin}(3, 1)$ and $\text{Pin}(1, 3)$.

If we pull back $P \rightarrow \mathcal{R}(M)$ along the evaluation map $\text{ev}: M \times \Gamma(\mathcal{R}(M)) \rightarrow \mathcal{R}(M)$, defined as $\text{ev}(x, g) := g(x)$, we obtain the principal Pin^{abc} -bundle $\text{ev}^*P \rightarrow M \times \Gamma(\mathcal{R}(M))$.

It consists of all pairs $(p, g) \in P \times \Gamma(\mathcal{R}(M))$ where p lies in g^*P . The maps u^g for the different metrics $g \in \Gamma(\mathcal{R}(M))$ then combine to a single map $u: \text{ev}^*P \rightarrow F(M)$, defined by $u(p, g) := u^g(p)$. We say that u^g *depends continuously on g* if the map $u: \text{ev}^*P \rightarrow F(M)$ is continuous.

If u^g depends continuously on g , then we obtain an infinitesimal action of $\text{Diff}(M)$ on the configuration space $\Gamma(\Sigma)$ of fermions coupled to gravity. Since the (left) action of $\text{Diff}(M)$ on $F(M)$ commutes with the (right) action of $\text{Gl}(4, \mathbb{R})$, we have an action of $\text{Diff}(M)$ on $\mathcal{R}(M)$, yielding the usual space-time transformation behavior $g_x \mapsto D\phi g_{\phi^{-1}(x)}$ on the space $\Gamma(\mathcal{R}(M))$ of metrics. To obtain the transformation behavior of spinors coupled to gravity, note that since $u: \text{ev}^*P \rightarrow F(M)$ is continuous, it induces a double cover from ev^*P to $\text{ev}^*F(M)$, the space of all pairs $(e_x, g) \in F(M) \times \Gamma(\mathcal{R}(M))$ with $e_x \in O^g(M)$. Since $\text{Diff}(M)$ acts on $\text{ev}^*F(M)$, it has an *infinitesimal* action on the double cover ev^*P . This yields an infinitesimal action on $\text{ev}^*\Sigma \rightarrow M \times \Gamma(\mathcal{R}(M))$, the space of all pairs $([\hat{q}_x, v], g) \in \Sigma \times \Gamma(\mathcal{R}(M))$ where \hat{q}_x is in g^*P . This yields an infinitesimal action on $\Gamma(\Sigma)$, since a section $\Psi \in \Gamma(\Sigma)$ can be viewed as a map from M to $\text{ev}^*\Sigma$, sending $x \in M$ to the pair (Ψ_x, g) , where g is the metric obtained from the section Ψ .

To recover the restriction $c = -1$, consider the case $M = \mathbb{R}^4$. Since $\text{Gl}(4, \mathbb{R})$ is a subgroup of $\text{Diff}(\mathbb{R}^4)$, it acts from the left on $F(\mathbb{R}^4)$, and hence on $\mathcal{R}(\mathbb{R}^4)$. Since only the Lorentz group $\text{O}(3, 1)$ leaves the Minkowski metric η invariant, we obtain an injective, continuous map $\sigma: \text{Gl}(4, \mathbb{R})/\text{O}(3, 1) \rightarrow \mathbb{R}^4 \times \Gamma(\mathcal{R}(\mathbb{R}^4))$ by $\sigma([A]) := (0, A\eta)$. The pullback bundle $\sigma^*\text{ev}^*P$ is a principal Pin^{abc} -bundle over $\text{Gl}(4, \mathbb{R})/\text{O}(3, 1)$. Note that $\text{ev} \circ \sigma$ is a diffeomorphism from $\text{Gl}(4, \mathbb{R})/\text{O}(3, 1)$ to $\mathcal{R}_0(\mathbb{R}^4)$, the space of all Lorentzian metrics on the tangent space $T_0\mathbb{R}^4$ at the origin, so $\sigma^*\text{ev}^*P = (\text{ev} \circ \sigma)^*P$ can be identified with the restriction P_0 of P to $\mathcal{R}_0(\mathbb{R}^4)$. Since the image of the pullback map $\sigma^*u: \sigma^*\text{ev}^*P \rightarrow F(\mathbb{R}^4)$ is the set $F_0(\mathbb{R}^4) \simeq \text{Gl}(4, \mathbb{R})$ of frames at the origin, we obtain a continuous double cover $P_0 \rightarrow \text{Gl}(4, \mathbb{R})$. As this double cover intertwines the (right) Pin^{abc} -action on P_0 with the (right) $\text{O}(3, 1)$ -action on $F_0(\mathbb{R}^4) \simeq \text{Gl}(4, \mathbb{R})$, the preimage P_0^+ of $\text{Gl}(4, \mathbb{R})^+$ is the universal covering group $\tilde{\text{Gl}}(4, \mathbb{R})^+$, and the orientation-preserving subgroup of Pin^{abc} coincides with the subgroup of $\tilde{\text{Gl}}(4, \mathbb{R})^+$ that covers $\text{SO}(3, 1)$. Since $(\Lambda_P \Lambda_T)^2 = 1$ in $\tilde{\text{Gl}}(4, \mathbb{R})^+$, we recover the restriction $c = -1$ of Sec. V.

C. Double cover of the diffeomorphism group

In the above line of reasoning, the group structure on P_0^+ stems from its identification with the universal cover of the connected Lie group $\text{Gl}(4, \mathbb{R})^+$. Since we lack a group structure on the disconnected space P_0 , we cannot directly infer that $a = b$. This does, however, follow from the slightly stronger assumption that the $\text{Diff}(M)$ -action on $\text{ev}^*F(M)$ lifts to an action by automorphisms of a double

cover $\widehat{\text{Diff}}(M)$ on ev^*P . This yields an action of $\widehat{\text{Diff}}(M)$ on $\text{ev}^*\Sigma$, and by identifying $\Psi \in \Gamma(\Sigma)$ with a map from M to $\text{ev}^*\Sigma$ as before, one obtains an action of $\widehat{\text{Diff}}(M)$ on $\Gamma(\Sigma)$. Explicitly, $\phi \in \text{Diff}(M)$ acts on $\text{ev}^*F(M)$ by taking (e_x, g) to $(D\phi(e_x), D\phi \circ g \circ \phi^{-1})$. If this lifts to an automorphism $D\hat{\phi}$ of $\text{ev}^*\Sigma$, then $D\hat{\phi}$ maps $\Psi \in \Gamma(\Sigma)$ to the unique $\Psi' \in \Gamma(\Sigma)$ with $(\Psi'_x, D\phi \circ g \circ \phi^{-1}) = D\hat{\phi}(\Psi_{\phi^{-1}(x)}, g)$.

To see that this yields the restriction $a = b$, consider the case $M = \mathbb{R}^4$. Then $\text{Gl}(4, \mathbb{R})$ is a subgroup of $\text{Diff}(\mathbb{R}^4)$, and its preimage in $\widehat{\text{Diff}}(\mathbb{R}^4)$ is one of the two Gin groups Gin^\pm . The (left) action of Gin^\pm by automorphisms on ev^*P covers the (left) action of $\text{Gl}(4, \mathbb{R})$ by automorphisms on $\text{ev}^*F(M)$, so in particular, the (left) action of Gin^\pm on $\sigma^*\text{ev}^*P = P_0$ covers the (left) action of $\text{Gl}(4, \mathbb{R})$ on $\sigma^*\text{ev}^*F(\mathbb{R}^4) = F_0(\mathbb{R}^4)$. This intertwines the (right) action of Pin^{abc} on P_0 with the (right) action of $\text{O}(3, 1)$ on $F_0(\mathbb{R}^4)$. Since all these actions are free, we can identify Pin^{abc} with a subgroup of Gin^\pm that covers the Lorentz group $\text{O}(3, 1)$. Following the line of reasoning in Sec. V, we thus find $a = b$ as well as $c = -1$.

We conclude that although an infinitesimal action of the space-time diffeomorphism group on the configuration space of fermions coupled to gravity requires $c = -1$, an action of a double cover of the diffeomorphism group can only be achieved if the Pin group additionally satisfies the relation $a = b$.

VIII. DISCUSSION

The conclusion that only two of the eight Pin groups are compatible with general relativity, appears to be quite robust. It is based on the elementary observation that the twofold spin cover of the orthonormal frame bundle $O^g(M)$ is compatible with a twofold cover of the *full* frame bundle $F(M)$. Although we derived this from the setting outlined in Sec. IV (going back to [3–5] in the Riemannian and [6–9] in the Lorentzian case), the use of double covers of the full frame bundle—and hence our conclusion that only two Pin groups are admissible—is common to many other approaches, such as the more “global” formalism developed in [2, 10–12]. In fact, the restrictions on the Pin groups are closely linked to the transformation behavior of fermions coupled to gravity under space-time diffeomorphisms.

Since *any* principal bundle with an infinitesimal action of the space-time diffeomorphism group is associated to a discrete cover of a (higher order) frame bundle [7, 8], we expect that our restrictions on the Pin group are not an artefact of the particular description that we have adopted.

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- [20] In first order formalisms such as the Palatini approach, one considers the bigger configuration space of metrics g together with a connection ∇ . This amounts to replacing $\mathcal{R}(M)$ by $J^1F(M)/O(3, 1)$.
- [21] The requirement that V extends to an \mathbb{R} -linear Pin^{abc} -representation may place restrictions on a , b and c .
- For instance, the $\text{Spin}^\uparrow(3, 1)$ -representation $V = \mathbb{C}^2$ (a single Majorana fermion) extends to Pin^{abc} if and only if $a = 1$ and $b = -1$. To be consistent with the topological restrictions derived in Sec. V, we therefore need $m \geq 2$.
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